Global solutions to the stochastic reaction-diffusion equation with superlinear forcing and superlinear multiplicative noise

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## Stochastic reaction-diffusion equation

Bounded spatial domain $x \in D \subset \mathbb{R}^{d}$.

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\left\{\begin{array}{l}
\frac{\partial u}{\partial t}(t, x)=\mathcal{A} u(t, x)+f(u(t, x))+\sigma(u(t, x)) \dot{w}(t, x) \quad t>0, \quad x \in D \\
u(t, x)=0, \quad x \in \partial D
\end{array}\right.
$$

- Second-order elliptic differential operator

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\mathcal{A}=\sum_{i=1}^{d} \sum_{j=1}^{d} \frac{\partial}{\partial x_{i}}\left(a_{i j}(x) \frac{\partial}{\partial x_{j}}\right) .
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- Multiplicative noise term $\sigma(u(t, x)) \dot{w}$.


## Introduction to stochastic PDE

- The solution to the non-stochastic heat equation $\frac{\partial u}{\partial t}=\frac{1}{2} \Delta u$, $u(0, x)=u_{0}(x)$ can be written as a convolution with a heat kernel

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- Duhamel's principle gives the solution to the semilinear heat equation $\frac{\partial u}{\partial t}=\frac{1}{2} \Delta u+f(u)$

$$
u(t, x)=\int_{D} K(t, x, y) u_{0}(y) d y+\int_{0}^{t} \int_{D} K(t-s, x, y) f(u(s, y)) d y d s
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- Formally, solution should be given by

$$
u(t)=S(t) u_{0}+\int_{0}^{t} S(t-s) f(u(s)) d s+\int_{0}^{t} S(t-s) \sigma(u(s)) d w(s)
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- Mueller and collaborators (1991, with Sowers 1993, 1998, 2000) showed that in the space-time white noise case with $\mathcal{A}=\frac{\partial^{2}}{\partial x^{2}}$, $f \equiv 0$ and $\sigma(u)=|u|^{\gamma}$, solutions never explode if $\gamma<\frac{3}{2}$ and explode with positive probability if $\gamma>\frac{3}{2}$.


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- Dissipative forces push solutions away from $\pm \infty$, counteracting expansion due to noise.
- I outline some recent results about sufficient conditions in the accretive and dissipative settings that guarantee that mild solutions to the SRDE do not explode.


## Technical Assumptions on $\mathcal{A}$ and $\dot{w}$

- Eigenvalues of $\mathcal{A},\left\{e_{k}(x)\right\}$ complete orthonormal basis of $L^{2}(D)$.

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- Formal definition of the noise: Sequence $\lambda_{j} \geq 0$

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\dot{w}(t, x)=\sum_{j=1}^{\infty} \lambda_{j} e_{j}(x) d \beta_{j}(t), \quad \beta_{j}(t) \text { i.i.d. one-dimensional B.M. }
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- Condition (Cerrai 2003) - There exist $\theta>0, \rho \in[2,+\infty]$

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\begin{aligned}
\sum_{k=1}^{\infty} \alpha_{k}^{-\theta}\left|e_{k}\right|_{L^{\infty}}^{2} & <+\infty, \sum_{j=1}^{\infty} \lambda_{j}^{\rho}\left|e_{j}\right|_{L^{\infty}}^{2}<+\infty \text { or } \sup _{j} \lambda_{j}<+\infty(\rho=\infty) \\
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- Space Hölder continuity $\approx(1-\eta)$. Time Hölder continuity $\approx \frac{1-\eta}{2}$.


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- If $d=1$ we can take $\lambda_{j} \equiv 1$ (space-time white noise), $\rho=\infty$ $\theta \in(1 / 2,1), \eta=\theta$ arbitrarily close to $1 / 2$.
- If $d>1$, space-time white noise not allowed. On rectangular domains we need $\theta>\frac{d}{2}, \rho<\frac{2 d}{d-2}$.


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- Generalization to higher spatial dimensions and other more general conditions in dissertation by Bezdek.


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- Similar conditions identified by Liu and Röckner (2010) for monotone SPDEs

$$
2_{V^{*}}\langle A(t, v), v\rangle_{V}+\|B(t, v)\|_{2}^{2}+\theta\|v\|_{V}^{\alpha} \leq C+K\|v\|_{H}^{2} .
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- Note: Trace-class noise means $\eta=0$ and this coincides with the Ito formula condition $\gamma<\frac{\beta+1}{2}$.
- When $\eta$ not trace-class, Ito formula arguments are not available.
- As long as $\beta>1, \gamma$ can grow superlinearly.


## Idea of the proof

- Set up a sequence of stopping times

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- There exist $C>0, q>1$ such that for any $k \in \mathbb{N}, \varepsilon>0$,

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- This is enough to prove that $\sum_{k}\left(\tau_{k+1}-\tau_{k}\right)=+\infty$. Cannot explode in finite time.


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u\left(t+\tau_{k}\right)= & S(t) u\left(\tau_{k}\right)+\int_{\tau_{k}}^{\tau_{k}+t} S\left(\tau_{k}+t-s\right) f(u(s)) d s \\
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- If $Z_{k}(t):=\int_{\tau_{k}}^{\tau_{k}+t} S\left(\tau_{k}+t-s\right) \sigma(u(s)) d W(s)$ satisfies

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- When $2 \alpha+\eta+\frac{2(\gamma-1)}{\beta-1}<1$, this inner integral is bounded as $t \downarrow 0$. (Beta function)


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- $\mathbb{P}\left(\left|u\left(\tau_{k+1}\right)\right|_{L^{\infty}}=3\left|u\left(\tau_{k}\right)\right|_{L^{\infty}}, \tau_{k+1}-\tau_{k}<\left.\varepsilon| | u\left(\tau_{k}\right)\right|_{L^{\infty}}=c_{0} 3^{n}\right) \leq$ $\mathbb{P}\left(\sup _{t \in\left[0, \varepsilon \wedge\left(\tau_{k+1}-\tau_{k}\right)\right]}\left|Z_{k}(t)\right|_{L^{\infty}} \geq\left.\frac{1}{9}\left|u\left(\tau_{k}\right)\right|_{L^{\infty}}| | u\left(\tau_{k}\right)\right|_{L^{\infty}}=c_{0} 3^{n}\right)$.


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## Example

Theorem (S. 2022 - same theorem)
Assume $f(u) \operatorname{sign}(u) \leq-\mu|u|^{\beta}$ for $|u|>c_{0},|\sigma(u)| \leq C\left(1+|u|^{\gamma}\right)$, with $\gamma<1+\frac{(1-\eta)(\beta-1)}{2}$. Then mild solutions can never explode.

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- If $\beta \leq 3$, the Mueller results give a stronger result, so $\gamma$ can be any value $\gamma<\frac{3}{2}$.
- The semigroup causes $t^{-\frac{1}{2}}$ decay but the nonlinear term causes $t^{-\frac{1}{\beta-1}}$ decay.


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## Osgood condition for SRDE $\frac{\partial u}{\partial t}=\mathcal{A} u+f(u)+\sigma(u) \dot{W}$

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- Bonder and Groisman (2009) showed that in the case of additive noise $\sigma(u) \equiv 1>0$, if $\int_{c}^{\infty} \frac{1}{f(x)} d x<\infty$ for some $c>0$, then the SRDE explodes in finite time with probability one.


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- Dalang, Khoshnevisan, and Zhang (2019) showed that $\sigma$ can be superlinear too. Studied the space-time white noise on a bounded one-dimensional spatial domain
- Assume $|f(u)| \leq C(1+|u| \log |u|), \sigma \in o\left(|u|(\log |u|)^{\frac{1}{4}}\right)$.
- Then solutions never explode.


## Result for accretive $f$

## Theorem (S. 2022)

Assume $f$ and $\sigma$ are locally Lipschitz continuous functions. Assume that there exists a positive, increasing function $h:[0,+\infty) \rightarrow[0,+\infty)$ such that

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\begin{gathered}
\int_{0}^{\infty} \frac{1}{h(u)} d u=\infty \text { and } \\
|f(u)| \leq h(|u|)
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and there exists $\gamma \in\left(0, \frac{1-\eta}{2}\right)$ such that

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If $h(u)=u \log (u) \log \log (u)$ then $\sigma(u) \leq u(\log u \log \log u)^{\gamma}$.

## Local mild solution and explosion

- Define the cutoff versions of $f$ and $\sigma$

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f_{n}(x)= \begin{cases}f(x) & \text { if } x \in\left[-3^{n}, 3^{n}\right] \\ f\left(3^{n}\right) & \text { if } x>3^{n} \\ f\left(-3^{n}\right) & \text { if } x<-3^{n}\end{cases}
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- Solution EXPLODES in finite time if $\sup _{n} \tau_{n}<+\infty$ and solution is global in time if $\sup _{n} \tau_{n}=+\infty$.


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- One interpretation is that time Hölder continuity is slightly worse than $\frac{1-\eta}{2}$ - same number that describes how superlinear $\sigma$ can be.


## Idea of the proof

- Consider the ODE $\frac{d v}{d t}=f(v(t))$ where $|f(v)| \leq h(|v|)$ and $\int_{0}^{\infty} \frac{1}{h(v)} d v=+\infty$ and $h$ increasing.


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- This is enough to prove that solutions cannot explode in finite time.

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\sum_{n} \frac{2^{n}}{h\left(2^{n+1}\right)}=+\infty \text { iff } \int_{1}^{\infty} \frac{2^{x}}{h\left(2^{x}\right)} d x=+\infty \text { iff } \int_{0}^{\infty} \frac{1}{h(x)} d x=+\infty
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- So the times required to double have the property that $\sum\left(T_{n+1}-T_{n}\right)=+\infty$. Cannot explode in finite time.


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- By the moment bounds and $\sigma(u(s)) \leq 3^{(n+1)(1-\gamma)}\left(h\left(3^{n+1}\right)^{\gamma}\right.$,

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- By the Borel-Cantelli Lemma

$$
\sum_{n} \mathbb{P}\left(\tau_{n+1}-\tau_{n} \leq a_{n+1}\right) \leq \sum_{n} C n^{-q} \text { for some } q>1
$$

- With probability one $\tau_{n+1}-\tau_{n}>a_{n}$ for all large $n$.
- Therefore $\mathbb{P}\left(\sup _{n} \tau_{n}=+\infty\right)=1$.


## Thank you

