Global solutions to the stochastic reaction-diffusion equation with superlinear forcing and superlinear multiplicative noise

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Bounded spatial domain $x \in D \subset \mathbb{R}^d$.

$$\begin{cases} \frac{\partial u}{\partial t}(t,x) = \mathcal{A}u(t,x) + f(u(t,x)) + \sigma(u(t,x))\dot{w}(t,x) \quad t > 0, \quad x \in D, \\ u(t,x) = 0, \quad x \in \partial D \end{cases}$$

• Second-order elliptic differential operator $\mathcal{A} = \sum_{i=1}^{d} \sum_{j=1}^{d} \frac{\partial}{\partial x_i} \left(a_{ij}(x) \frac{\partial}{\partial x_j} \right).$ Bounded spatial domain $x \in D \subset \mathbb{R}^d$.

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- Multiplicative noise term $\sigma(u(t,x))\dot{w}$.

• The solution to the non-stochastic heat equation $\frac{\partial u}{\partial t} = \frac{1}{2}\Delta u$, $u(0,x) = u_0(x)$ can be written as a convolution with a heat kernel

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- If the domain is the whole space then $K(t, x, y) = (2\pi t)^{-\frac{d}{2}} e^{-\frac{|x-y|^2}{2t}}$.
- Duhamel's principle gives the solution to the semilinear heat equation $\frac{\partial u}{\partial t} = \frac{1}{2}\Delta u + f(u)$

$$u(t,x) = \int_D K(t,x,y)u_0(y)dy + \int_0^t \int_D K(t-s,x,y)f(u(s,y))dyds.$$

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• Formally, solution should be given by

$$u(t) = S(t)u_0 + \int_0^t S(t-s)f(u(s))ds + \int_0^t S(t-s)\sigma(u(s))dw(s)$$

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• Mueller and collaborators (1991, with Sowers 1993, 1998, 2000) showed that in the space-time white noise case with $\mathcal{A} = \frac{\partial^2}{\partial x^2}$, $f \equiv 0$ and $\sigma(u) = |u|^{\gamma}$, solutions never explode if $\gamma < \frac{3}{2}$ and explode with positive probability if $\gamma > \frac{3}{2}$.

Salins (BU)

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- Dissipative forces push solutions away from $\pm \infty$, counteracting expansion due to noise.
- I outline some recent results about sufficient conditions in the accretive and dissipative settings that guarantee that mild solutions to the SRDE do not explode.

• Eigenvalues of \mathcal{A} , $\{e_k(x)\}$ complete orthonormal basis of $L^2(D)$.

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 Ae_k = −α_ke_k, 0 ≤ α_k ≤ α_{k+1}.
- Formal definition of the noise: Sequence $\lambda_j \ge 0$

 $\dot{w}(t,x) = \sum_{j=1}^{\infty} \lambda_j e_j(x) d\beta_j(t), \quad \beta_j(t) \text{ i.i.d. one-dimensional B.M.}$

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• Condition (Cerrai 2003) - There exist $\theta>0,\,\rho\in[2,+\infty]$

$$\begin{split} \sum_{k=1}^{\infty} \alpha_k^{-\theta} |e_k|_{L^{\infty}}^2 < +\infty, \quad \sum_{j=1}^{\infty} \lambda_j^{\rho} |e_j|_{L^{\infty}}^2 < +\infty \text{ or } \sup_j \lambda_j < +\infty \left(\rho = \infty\right) \\ \eta := \frac{\theta(\rho - 2)}{\rho} < 1 \qquad (\eta := \theta \text{ if } \rho = +\infty) \,. \end{split}$$

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• Space Hölder continuity $\approx (1 - \eta)$. Time Hölder continuity $\approx \frac{1 - \eta}{2}$.

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- If d > 1, space-time white noise not allowed. On rectangular domains we need $\theta > \frac{d}{2}$, $\rho < \frac{2d}{d-2}$.

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- Generalization to higher spatial dimensions and other more general conditions in dissertation by Bezdek.

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• Similar conditions identified by Liu and Röckner (2010) for monotone SPDEs

$$2_{V^*} \langle A(t,v), v \rangle_V + \|B(t,v)\|_2^2 + \theta \|v\|_V^{\alpha} \le C + K \|v\|_H^2.$$

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- When η not trace-class, Ito formula arguments are not available.
- As long as $\beta > 1$, γ can grow superlinearly.

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• There exist C > 0, q > 1 such that for any $k \in \mathbb{N}$, $\varepsilon > 0$,

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• By Borel-Cantelli setting $\varepsilon = \frac{1}{k}$,

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- Eventually, $|u(\tau_k)|_{L^{\infty}}$ decreases or it takes more than 1/k time to triple. If it explodes it has more up steps than down steps.
- This is enough to prove that $\sum_{k} (\tau_{k+1} \tau_k) = +\infty$. Cannot explode in finite time.

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If β > 1, the deterministic ODE dv/dt = -|v(t)|^βsign(v(t)) has solution |v(t)| = (|v(0)|^{-(β-1)} + Ct)^{-1/β-1} ≤ max{|v(0)|, Ct^{-1/β-1}}.
 The mild solution to the SPDE

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- If $Z_k(t) := \int_{\tau_k}^{\tau_k+t} S(\tau_k+t-s)\sigma(u(s))dW(s)$ satisfies $|Z_k(t)|_{L^{\infty}} \leq \frac{1}{3}|u(t+\tau_k)|_{L^{\infty}}$, for $t \in [0, \tau_{k+1}-\tau_k]$,

$$|u(t)|_{L^{\infty}} \le \frac{3}{2} \left(|u(\tau_k)|_{L^{\infty}}^{-(\beta-1)} + Ct \right)^{-\frac{1}{\beta-1}}$$

• Moment bounds (Cerrai 2003) – For p > 1 large enough, and small numbers $\alpha, \gamma \in (0, (1 - \eta)/2)$ such that $(\alpha - \frac{\zeta}{2})p > 1$ and

$$\mathbb{E} \sup_{t \in [0,T]} \sup_{x \in D} |Z_k(t,x)|^p \le CT^{p(\alpha-\frac{\zeta}{2})-1} \int_0^T \mathbb{E} \left(\int_0^t (t-s)^{-2\alpha-\eta} |\sigma(u(s+\tau_k))|^2_{L^{\infty}} ds \right)^{\frac{p}{2}} dt.$$

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• When $2\alpha + \eta + \frac{2(\gamma-1)}{\beta-1} < 1$, this inner integral is bounded as $t \downarrow 0$. (Beta function)

Salins (BU)

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$$\mathbb{P}\left(|u(\tau_{k+1})|_{L^{\infty}} = 3|u(\tau_k)|_{L^{\infty}}, \tau_{k+1} - \tau_k < \varepsilon \Big| |u(\tau_k)|_{L^{\infty}} = c_0 3^n \right) \leq \mathbb{P}\left(\sup_{t \in [0, \varepsilon \land (\tau_{k+1} - \tau_k)]} |Z_k(t)|_{L^{\infty}} \geq \frac{1}{9} |u(\tau_k)|_{L^{\infty}} \Big| |u(\tau_k)|_{L^{\infty}} = c_0 3^n \right).$$

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- If it were to explode, then it would need more up steps than down steps. This is enough to guarantee that $\lim \tau_k = +\infty$.

Theorem (S. 2022 - same theorem)

Assume $f(u) \operatorname{sign}(u) \leq -\mu |u|^{\beta}$ for $|u| > c_0$, $|\sigma(u)| \leq C(1 + |u|^{\gamma})$, with $\gamma < 1 + \frac{(1-\eta)(\beta-1)}{2}$. Then mild solutions can never explode.

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• Consider the non-linear heat equation on a one-dimensional interval with space-time white noise

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- If $\beta \leq 3$, the Mueller results give a stronger result, so γ can be any value $\gamma < \frac{3}{2}$.
 - The semigroup causes $t^{-\frac{1}{2}}$ decay but the nonlinear term causes $t^{-\frac{1}{\beta-1}}$ decay.

Salins (BU)

Superlinear accretive forcing
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• Explodes in finite time if and only if $\int_c^{\infty} \frac{1}{f(x)} dx < +\infty$.

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 - Global solution if $\int_c^{\infty} \frac{1}{f(x)} dx = +\infty$.
- Dalang, Khoshnevisan, and Zhang (2019) showed that σ can be superlinear too. Studied the space-time white noise on a bounded one-dimensional spatial domain
 - Assume $|f(u)| \le C(1+|u|\log|u|), \ \sigma \in o(|u|(\log|u|)^{\frac{1}{4}}).$
 - Then solutions never explode.

Theorem (S. 2022)

Assume f and σ are locally Lipschitz continuous functions. Assume that there exists a positive, increasing function $h: [0, +\infty) \to [0, +\infty)$ such that

$$\int_{0}^{\infty} \frac{1}{h(u)} du = \infty \text{ and}$$
$$|f(u)| \le h(|u|)$$
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$$|\sigma(u)| \le |u|^{1-\gamma} (h(|u|))^{\gamma} \text{ for all } |u| > 1.$$

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• Define the cutoff versions of f and σ

$$f_n(x) = \begin{cases} f(x) & \text{if } x \in [-3^n, 3^n] \\ f(3^n) & \text{if } x > 3^n \\ f(-3^n) & \text{if } x < -3^n \end{cases}$$

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- Solution EXPLODES in finite time if $\sup_n \tau_n < +\infty$ and solution is global in time if $\sup_n \tau_n = +\infty$.

Salins (BU)

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• For any $\zeta \in (0, 1 - \eta)$, $p > \max\left\{\frac{2}{1 - \eta - \zeta}, \frac{d}{\zeta}\right\}$, there exists $C = C(\zeta, p)$ such that for any adapted random field that is almost surely bounded

$$\mathbb{P}\left(\sup_{s\in[0,t]}\sup_{x\in D}|\Phi(t,x)| \le M\right) = 1.$$
$$\mathbb{E}\sup_{t\in[0,\varepsilon]}\left|\int_{0}^{t}S(t-s)\Phi(s)dw(s)\right|_{L^{\infty}}^{p} \le CM^{p}\varepsilon^{\frac{p(1-\eta-\zeta)}{2}}.$$
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• One interpretation is that time Hölder continuity is slightly worse than $\frac{1-\eta}{2}$ – same number that describes how superlinear σ can be.

Idea of the proof

• Consider the ODE $\frac{dv}{dt} = f(v(t))$ where $|f(v)| \le h(|v|)$ and $\int_0^\infty \frac{1}{h(v)} dv = +\infty$ and h increasing.

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• This is enough to prove that solutions cannot explode in finite time.

$$\sum_{n} \frac{2^{n}}{h(2^{n+1})} = +\infty \text{ iff } \int_{1}^{\infty} \frac{2^{x}}{h(2^{x})} dx = +\infty \text{ iff } \int_{0}^{\infty} \frac{1}{h(x)} dx = +\infty.$$

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• So the times required to double have the property that $\sum (T_{n+1} - T_n) = +\infty$. Cannot explode in finite time.

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- For any fixed n and t > 0,

$$u(\tau_n + t) = S(t)u(\tau_n) + \int_{\tau_n}^{\tau_n + t} S(\tau_n + t - s)f(u(s))ds$$
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• For $t \in [0, a_{n+1} \land (\tau_{n+1} - \tau_n)]$ $|S(t)u(\tau_n)|_{L^{\infty}} \le |u(\tau_n)|_{L^{\infty}} \le 3^n.$

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$$\left| \int_{\tau_n}^{\tau_n + t} S(\tau_n + t - s) f(u(s)) ds \right|_{L^{\infty}} \le a_{n+1} h(3^{n+1}) \le 3^n.$$

Salins (BU)

• The only way that $(\tau_{n+1} - \tau_n)$ can be less than a_{n+1} is if

$$\sup_{t \in [0, a_{n+1} \land (\tau_{n+1} - \tau_n)]} \left| \int_{\tau_n}^{\tau_n + t} S(\tau_n + t - s) \sigma(u(s)) dw(s) \right|_{L^{\infty}} > 3^n.$$

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• Chebyshev

$$\leq 3^{-np} \mathbb{E} \sup_{t \in [0, a_{n+1} \land (\tau_{n+1} - \tau_n)]} \left| \int_{\tau_n}^{\tau_n + t} S(\tau_n + t - s) \sigma(u(s)) dw(s) \right|_{L^{\infty}}^p$$

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• By the moment bounds and $\sigma(u(s)) \leq 3^{(n+1)(1-\gamma)} (h(3^{n+1})^{\gamma})$,

$$\leq 3^{-np} a_{n+1}^{\frac{p(1-\eta-\zeta)}{2}} 3^{p(1-\gamma)(n+1)} (h(3^{n+1}))^{\gamma p} \leq C a_{n+1}^{\frac{p(1-\eta-\zeta-2\gamma)}{2}}$$

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$$\mathbb{P}(\tau_{n+1} - \tau_n \le a_{n+1}) \le C a_{n+1}^{\frac{p(1-\eta-\zeta-2\gamma)}{2}} \le C n^{-\frac{p(1-\eta-\zeta-2\gamma)}{2}}$$
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• Therefore
$$\mathbb{P}(\sup_n \tau_n = +\infty) = 1.$$

Thank you