

p. 45 #5 (a) If $u(x,t) = -2xt - x^2$ then

$$u_t = -2x \text{ while } u_x = -2t - 2x \text{ and } u_{xx} = -2$$

$$\therefore u_t = xu_{xx}$$

By elementary calculus one finds that the max. value of $u(x,t)$ is 1 which occurs at $(x,t) = (-1,1)$. This is not on the bottom or sides of the rectangle & hence the max. principle fails.

(b) If one tries to define $v_\epsilon(x,t) = u(x,t) + \epsilon x^2$ as in the proof of the max. theorem then

$$\partial_t v_\epsilon - x \partial_x^2 v_\epsilon = -2\epsilon x.$$

which is no longer strictly negative. Thus, we cannot rule out that v (& hence u) could have a max. inside or on the top of the rectangle.

p. 45 #8. Define

$$E[u](t) = \int_0^L u^2(x,t) dx$$

Then

$$\frac{dE}{dt} = 2 \int_0^L u u_t dx$$

Int. by parts

$$= 2k \int_0^L u u_{xx} dx$$

$$= 2k \left\{ u u_x \Big|_0^L - \int_0^L (u_x(x,t))^2 dx \right\}$$

use b.c.'s
to replace
 $u_x(0,t)$
and $u_x(L,t)$

$$= 2k \left\{ -a_2 (u(L,t))^2 - a_0 (u(0,t))^2 - \int_0^L (u_x(x,t))^2 dx \right\}$$

Since a_2 and a_0 are positive, each term in braces is less than or equal to zero so

$$\frac{dE}{dt} \leq 0.$$

page 52 #16:

$$u_t - k u_{xx} + b u = 0 \quad -\infty < x < \infty$$

$$u(x, 0) = \varphi(x)$$

Define $v(x, t) = e^{bt} u(x, t)$

Then

$$\begin{aligned} \partial_t v &= b e^{bt} u(x, t) + e^{bt} \partial_t u \\ &= b e^{bt} u(x, t) + k e^{bt} \partial_x^2 u - b e^{bt} u \end{aligned}$$

where the last equality used the PDE for u

$$\text{But } k e^{bt} \partial_x^2 u = k \partial_x^2 (e^{bt} u) = k \partial_x^2 v$$

so

$$\partial_t v = k \partial_x^2 v$$

and

$$v(x, 0) = u(x, 0) = \varphi(x) \quad (\text{since } e^{bt} = 1 \text{ for } t=0)$$

Thus,

$$v(x, t) = \int_{-\infty}^{\infty} S(x-y, t) \varphi(y) dy$$

$$\text{where } S(z, t) = \frac{1}{\sqrt{4\pi kt}} e^{-z^2/4kt}$$

or

$$u(x,t) = e^{-bt} v(x,t)$$

$$= \frac{e^{-bt}}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy$$

Solutions for HW#3.

p.58 #3: Since we want the function to have zero derivative at $x=0$, we extend $\varphi(x)$ to the whole line as an even function, i.e.

$$\varphi^e(x) = \begin{cases} \varphi(x) & \text{if } x > 0 \\ \varphi(-x) & \text{if } x < 0. \end{cases}$$

Then

$$w(x,t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt} \varphi^e(y) dy$$

$$= \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^0 e^{-(x-y)^2/4kt} \varphi^e(y) dy$$

$$+ \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} e^{-(x-y)^2/4kt} \varphi^e(y) dy$$

In the second integral $\varphi^e(y) = \varphi(y)$ since $y > 0$. In the first integral let $z = -y$. Then $\varphi^e(y) = \varphi^e(-z) = \varphi(z)$.

$$\therefore w(x,t) = \frac{1}{\sqrt{4\pi kt}} \int_{\infty}^0 e^{-(x+z)^2/4kt} \varphi(z) dz + \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} e^{-(x-y)^2/4kt} \varphi(y) dy$$

$$= \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} \{ e^{-(x-y)^2/4kt} + e^{-(x+y)^2/4kt} \} \varphi(y) dy$$

Note that $\partial_x w(0,t) = 0$, by explicit computation.

P.87 #1: (a) If one clamps a violin string in the middle that's like imposing Dirichlet b.c.s at 0 & $L/2$ instead of 0 & L so the eigenvalues change from

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \quad \text{to} \quad \tilde{\lambda}_n = \left(\frac{2n\pi}{L}\right)^2$$

This means that the $T(t)$ part of the solution changes from

$$T_n(t) = A_n \cos\left(\frac{n\pi ct}{L}\right) + B_n \sin\left(\frac{n\pi ct}{L}\right)$$

to

$$\tilde{T}_n(t) = \tilde{A}_n \cos\left(\frac{2n\pi ct}{L}\right) + \tilde{B}_n \sin\left(\frac{2n\pi ct}{L}\right).$$

Note that the frequencies of $\tilde{T}_n(t)$ is exactly twice that of $T_n(t)$ & doubling the frequency raises the pitch by an octave.

(b) Increasing the tension increases $c = \sqrt{T/\rho}$ & hence the frequency, since

$$T_n(t) = A_n \cos\left(\frac{n\pi ct}{L}\right) + B_n \sin\left(\frac{n\pi ct}{L}\right).$$

p.87 #3

$$\frac{\partial u}{\partial t} = i \frac{\partial^2 u}{\partial x^2} \quad 0 < x < L$$

$$u(0,t) = u(L,t) = 0.$$

Write $u(x,t) = T(t) X(x)$. Substituting this into Schrodinger's equation gives

$$\frac{\dot{T}}{iT} = -\frac{X''}{X} = \lambda.$$

Thus, the X -eqn becomes $X'' + \lambda X = 0$; $X(0) = X(L) = 0$
& we know this has non-zero solutions only if $\lambda_n = \left(\frac{n\pi}{L}\right)^2$,
 $n=1,2,\dots$ & $X_n(x) = \sin\left(\frac{n\pi x}{L}\right)$

The T equation then becomes

$$\dot{T}_n = -i\lambda_n T_n$$

or

$$T_n(x) = e^{-i\left(\frac{n\pi}{L}\right)^2 t} A_n$$

Thus, we get an ∞ -sequence of solutions

$$u_n(x,t) = A_n e^{-i\left(\frac{n\pi}{L}\right)^2 t} \sin\left(\frac{n\pi x}{L}\right)$$

& the general solution will have the form

$$u(x,t) = \sum_{n=1}^{\infty} A_n e^{-i\left(\frac{n\pi}{L}\right)^2 t} \sin\left(\frac{n\pi x}{L}\right)$$

p.87 #4 If we set $u(x,t) = T(t)X(x)$ & insert this in the PDE we can rewrite the equation as

$$-\frac{\ddot{T} + r\dot{T}}{c^2 T} = -\frac{X''}{X} = \lambda$$

For the X equation we have (again) solution $X_n(x) = \sin\left(\frac{n\pi x}{L}\right)$ when $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $n=1,2,\dots$. Then the T equation becomes

$$\ddot{T}_n + r\dot{T}_n + \left(\frac{n\pi c}{L}\right)^2 T_n = 0$$

If we write $T_n(t) = e^{mt}$ then m satisfies

$$m^2 + rm + \left(\frac{n\pi c}{L}\right)^2 = 0$$

This can also be written in terms of sines & cosines - please accept that form too.

$\Rightarrow$

$$m = \frac{-r \pm \sqrt{r^2 - 4\left(\frac{n\pi c}{L}\right)^2}}{2}$$

$$= -\frac{1}{2}r \pm \frac{i}{2}\sqrt{\left(\frac{2\pi cn}{L}\right)^2 - r^2}$$

since $\left(\frac{2\pi cn}{L}\right) > r$.

$$\therefore T_n(t) = e^{-rt/2} \left(A_n e^{i\frac{t}{2}\sqrt{\left(\frac{2\pi cn}{L}\right)^2 - r^2}} + B_n e^{-i\frac{t}{2}\sqrt{\left(\frac{2\pi cn}{L}\right)^2 - r^2}} \right)$$

Multiplying by $X_n(x)$ and summing over n we find

$$u(x,t) = \sum_{n=1}^{\infty} \left\{ A_n e^{i\frac{t}{2}\sqrt{\left(\frac{2\pi cn}{L}\right)^2 - r^2}} + B_n e^{-i\frac{t}{2}\sqrt{\left(\frac{2\pi cn}{L}\right)^2 - r^2}} \right\} e^{-\frac{rt}{2}} \sin\left(\frac{n\pi x}{L}\right)$$

One then solves for A_n & B_n using the initial conditions as we did for the wave equation.

Page 87 #6: If we assume $u(x,t) = X(x)T(t)$
The equation can be written as

$$X(x) \left(t \frac{dT}{dt} \right) = T(t) \left(\frac{d^2 X}{dx^2} \right) + 2T(t)X(x)$$

or

$$\frac{t \frac{dT}{dt} - 2}{T(t)} = \frac{d^2 X / dx^2}{X(x)} = -\lambda$$

with $X(0) = X(\pi) = 0$. The X -part of the problem is the same as we considered in the lecture with $l = \pi$ so we have

$$X_n(x) = \sin(nx)$$

$$\lambda_n = n^2 \quad n = 1, 2, \dots$$

The t -dependent part of the equation becomes

$$t \frac{dT}{dt} - 2T(t) = -\lambda T(t)$$

or

$$\frac{dT}{dt} = -\frac{(\lambda - 2)}{t} T(t)$$

which has the solution $T(t) = A t^{2-\lambda}$. Using
The allowed values of λ from above we have

$$T_n(t) = A_n t^{2-n^2}$$

We're interested in $t=0$ so note that we can only consider $n=1$ (since $T_n(0)$ is undefined if $n > 1$.) But for $n=1$ we have $u_1(x,t) = A_1 t \sin x$ & we get infinitely many solutions (corresponding to different values of A_1) all of which have $u_1(x,0) = 0$.

page 50 #2:

$$u(x,t) = \frac{3}{\sqrt{4\pi kt}} \int_{-\infty}^0 e^{-\frac{(x-y)^2}{4kt}} dy + \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} e^{-\frac{(x-y)^2}{4kt}} dy$$

$$= \frac{3}{\sqrt{\pi}} \int_{-\infty}^0 e^{-\left(\frac{x}{\sqrt{4kt}} - z\right)^2} dz + \frac{1}{\sqrt{\pi}} \int_0^{\infty} e^{-\left(\frac{x}{\sqrt{4kt}} - z\right)^2} dz$$

$$\stackrel{\substack{= \\ \uparrow \\ w = \left(\frac{x}{\sqrt{4kt}} - z\right)}}{=} -\frac{3}{\sqrt{\pi}} \int_{\infty}^{\frac{x}{\sqrt{4kt}}} e^{-w^2} dw - \frac{1}{\sqrt{\pi}} \int_{\frac{x}{\sqrt{4kt}}}^{-\infty} e^{-w^2} dw$$

$$= -\frac{3}{\sqrt{\pi}} \left\{ \int_{+\infty}^0 e^{-w^2} dw + \int_0^{\frac{x}{\sqrt{4kt}}} e^{-w^2} dw \right\}$$

$$- \frac{1}{\sqrt{\pi}} \left\{ \int_{\frac{x}{\sqrt{4kt}}}^0 e^{-w^2} dw + \int_0^{-\infty} e^{-w^2} dw \right\}$$

$$= \frac{3}{2} + \frac{1}{2} - \frac{3}{2} \operatorname{Erf}\left(\frac{x}{\sqrt{4kt}}\right) + \frac{1}{2} \operatorname{Erf}\left(\frac{x}{\sqrt{4kt}}\right)$$

$$= 2 - \operatorname{Erf}\left(\frac{x}{\sqrt{4kt}}\right)$$