

# 11

## Basics of Wavelets

References: I. Daubechies (Ten Lectures on Wavelets; Orthonormal Bases of Compactly Supported Wavelets)

Also: Y. Meyer, S. Mallat

Outline:

1. Need for time-frequency localization
2. Orthonormal wavelet bases: examples

3. Meyer wavelet
4. Orthonormal wavelets and multiresolution analysis

Signal:

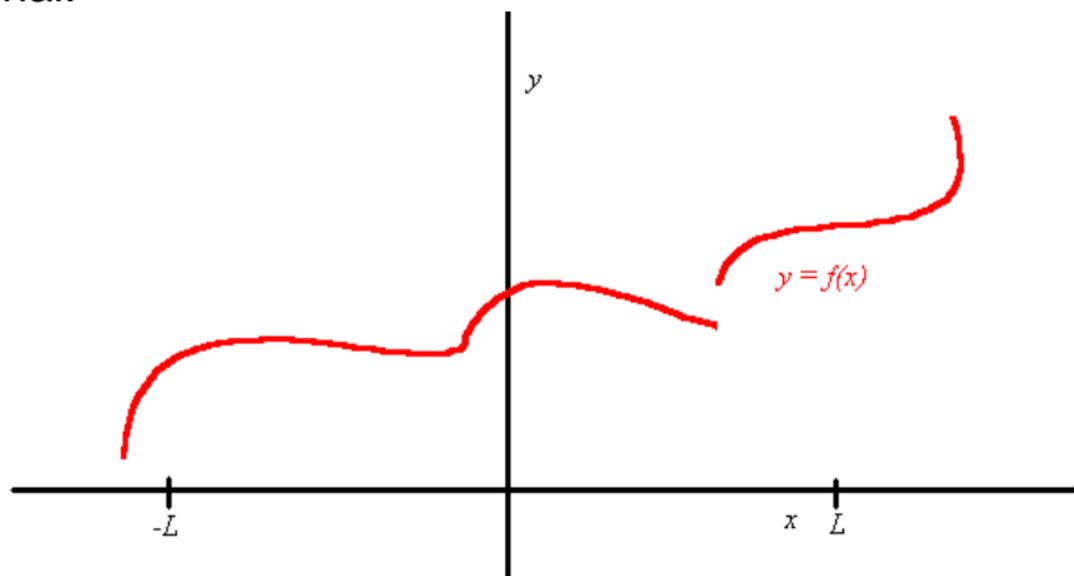


fig 1

Interested in “frequency content” of signal, locally in time.  
E.G., what is the frequency content in the interval  $[.5, .6]$ ?

Standard techniques: write in Fourier series as sum of sines and cosines: given function defined on  $[-L, L]$  as above:

$$f(x) =$$

$$\frac{1}{2} a_0 + \sum_{n=1}^{\infty} a_n \cos nx(\pi/L) + b_n \sin nx(\pi/L)$$

$(a_n, b_n$  constants)

$$a_n = \frac{1}{L} \int_{-L}^L dx f(x) \cos nx(\pi/L)$$

$$b_n = \frac{1}{L} \int_{-L}^L dx f(x) \sin nx(\pi/L)$$

(generally  $f$  is complex-valued and  $a_n, b_n$  are complex numbers).

## THEORY OF FOURIER SERIES

Consider function  $f(x)$  defined on  $[-L, L]$ .

Let  $L^2[-L, L] =$  square integrable functions

$$= \left\{ f : [-L, L] \rightarrow \mathbb{C} \mid \int_{-L}^L dx |f^2(x)| < \infty \right\}$$

where  $\mathbb{C} =$  complex numbers. Then  $L^2$  forms a Hilbert space.

Basis for Hilbert space:

$$\left\{ \frac{1}{\sqrt{L}} \cos nx(\pi/L), \frac{1}{\sqrt{L}} \sin nx(\pi/L) \right\}_{N=1}^{\infty}$$

(together with the constant function  $1/\sqrt{2L}$ ).

These vectors form an orthonormal basis for  $L^2$   
(constants  $1/\sqrt{L}$  give length 1).

## 2. Complex form of Fourier series (see previous lecture):

Equivalent representation:

Can use Euler's formula  $e^{ib} = \cos b + i \sin b$ . Can show similarly that the family

$$\left\{ \frac{1}{\sqrt{2L}} e^{inx(\pi/L)} \right\}_{n=-\infty}^{n=\infty} = \left\{ \frac{1}{\sqrt{2L}} \cos nx(\pi/L) + \frac{i}{\sqrt{2L}} \sin nx(\pi/L) \right\}_{n=-\infty}^{n=\infty}$$

is orthonormal basis for  $L^2$ .

Function  $f(x)$  can be written

$$f(x) = \sum_{n=-\infty}^{\infty} c_n \phi_n(x)$$

where

$$c_n = \langle \phi_n, f \rangle = \int_{-L}^L dx \overline{\phi_n(x)} f(x),$$

and

$$\phi_n(x) = n^{\text{th}} \text{ basis element} = \frac{1}{\sqrt{2L}} e^{inx(\pi/L)} .$$

### 3. FOURIER TRANSFORM

Fourier transform is “Fourier series” on entire line  $(-\infty, \infty)$ :

Start with function  $f(x)$  on  $(-L, L)$ :

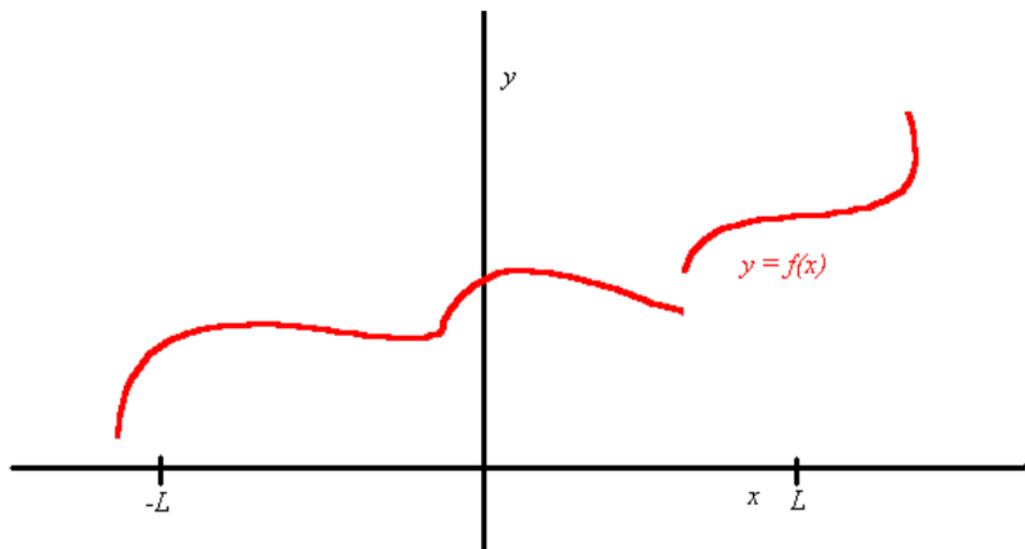


fig 2

$$f(x) = \sum_{n=-\infty}^{\infty} c_n \phi_n(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx(\pi/L)} / \sqrt{2L}$$

Let  $\xi_n = n\pi/L$ ; let  $\Delta\xi = \pi/L$ ;

let  $c(\xi_n) = c_n \sqrt{2\pi} / (\sqrt{2L} \Delta\xi)$ .

Then:

$$\begin{aligned}
f(x) &= \sum_{n=-\infty}^{\infty} c_n e^{inx(\pi/L)} / \sqrt{2L} \\
&= \sum_{n=-\infty}^{\infty} (c_n / \sqrt{2L}) e^{ix\xi_n} \\
&= \sum_{n=-\infty}^{\infty} c_n / (\sqrt{2L}\Delta\xi) e^{ix\xi_n} \Delta\xi. \\
&= \frac{1}{\sqrt{2\pi}} \sum_{n=-\infty}^{\infty} c_n \sqrt{2\pi} / (\sqrt{2L}\Delta\xi) e^{ix\xi_n} \Delta\xi. \\
&= \frac{1}{\sqrt{2\pi}} \sum_{n=-\infty}^{\infty} c(\xi_n) e^{ix\xi_n} \Delta\xi.
\end{aligned}$$

Note as  $L \rightarrow \infty$ , we have  $\Delta\xi \rightarrow 0$ , and

$$\begin{aligned}
c(\xi_n) &= c_n \sqrt{2\pi} / \left( \sqrt{2L\Delta\xi} \right) \\
&= \int_{-L}^L dx f(x) \overline{\phi_n(x)} \cdot \sqrt{2\pi} / \left( \sqrt{2L\Delta\xi} \right) \\
&= \int_{-L}^L dx f(x) \frac{\sqrt{2\pi}}{2L\Delta\xi} e^{-inx(\pi/L)} \\
&= \int_{-L}^L dx f(x) \frac{\sqrt{2\pi}}{2L(\pi/L)} e^{-inx(\pi/L)} \\
&= \frac{1}{\sqrt{2\pi}} \int_{-L}^L dx f(x) e^{-ix\xi_n} .
\end{aligned}$$

Now (informally) take the limit  $L \rightarrow \infty$ . The interval becomes

$$[-L, L] \rightarrow (-\infty, \infty).$$

We have

$$f(x) = \frac{1}{\sqrt{2\pi}} \sum_{n=-\infty}^{\infty} c(\xi_n) e^{ix\xi_n} \Delta\xi$$

fundamental thm. calculus

$$L \xrightarrow{\quad} \infty \quad \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} c(\xi) e^{ix\xi} d\xi$$

[note this is like Riemann sum of calculus, which turns into integral].

Finally, from above

$$c(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-L}^L dx f(x) e^{-ix\xi}$$

$$L \xrightarrow{\quad} \infty \quad \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx f(x) e^{-ix\xi}.$$

Thus, the informal arguments give that in the limit, we can write

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} c(\xi) e^{ix\xi} d\xi,$$

where  $c(\xi)$  (called *Fourier transform* of  $f$ ) is

$$c(\xi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} dx f(x) e^{ix\xi}$$

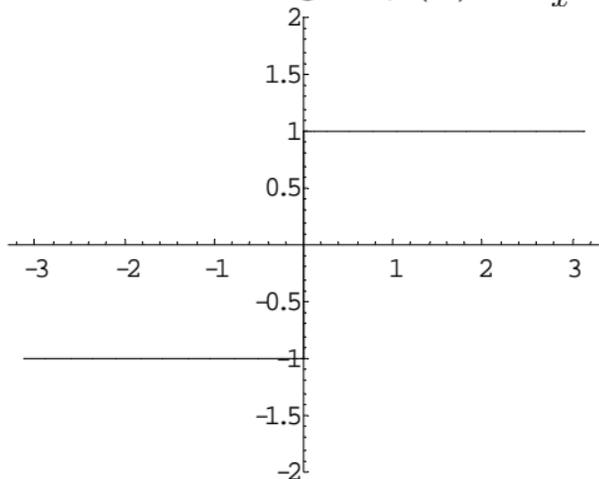
(like Fourier series with sums replaced by integrals over the real line).

**Note:** can prove that writing  $f(x)$  in the above integral form works for arbitrary  $f \in L^2(-\infty, \infty)$ .

## 4. FREQUENCY CONTENT AND GIBBS PHENOMENON

For now work with Fourier series on  $\mathbb{R}$ .  
discontinuous at  $x = 0$ , e.g. if  $f(x) = \frac{|x|}{x}$  :

If  $f(x)$



first few partial sums of Fourier series are:  
5 terms of FS:

$$\frac{4\sin x}{\pi} + \frac{4\sin 3x}{3\pi} + \frac{4\sin 5x}{5\pi}$$

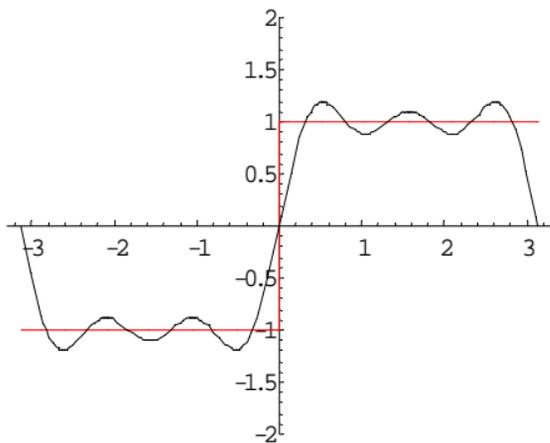
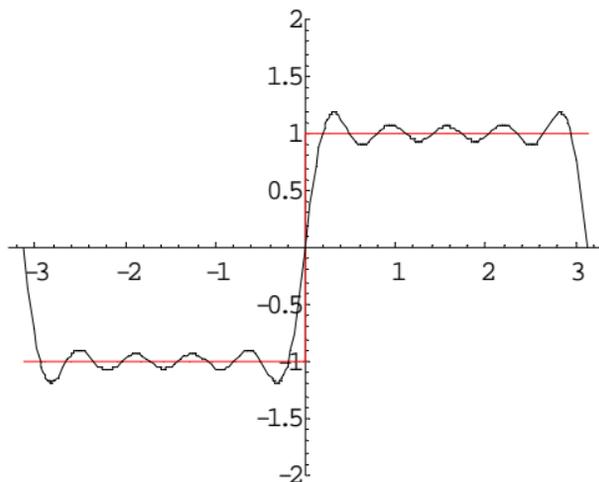


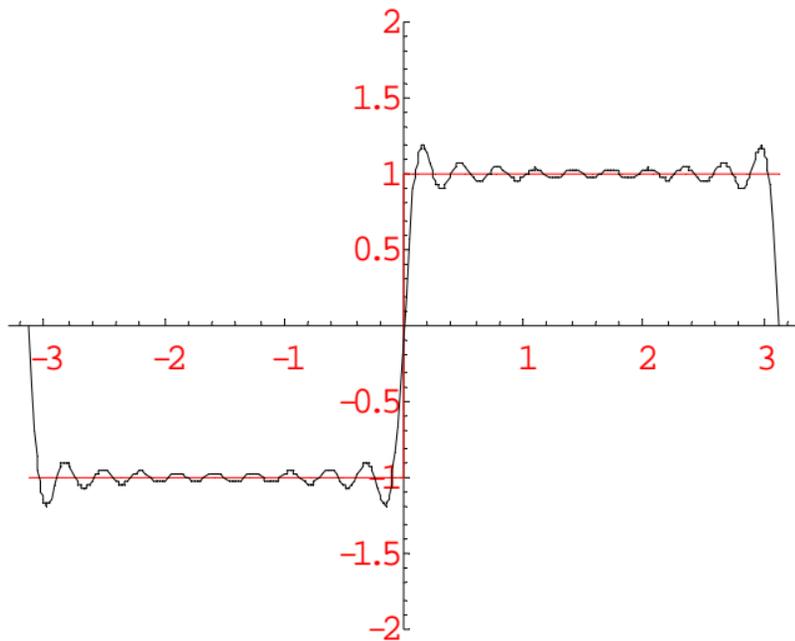
fig 3

10 terms:

$$\frac{4\sin x}{\pi} + \frac{4\sin 3x}{3\pi} + \frac{4\sin 5x}{5\pi} + \dots + \frac{4\sin 10x}{10\pi}$$



20 terms:



40 terms:

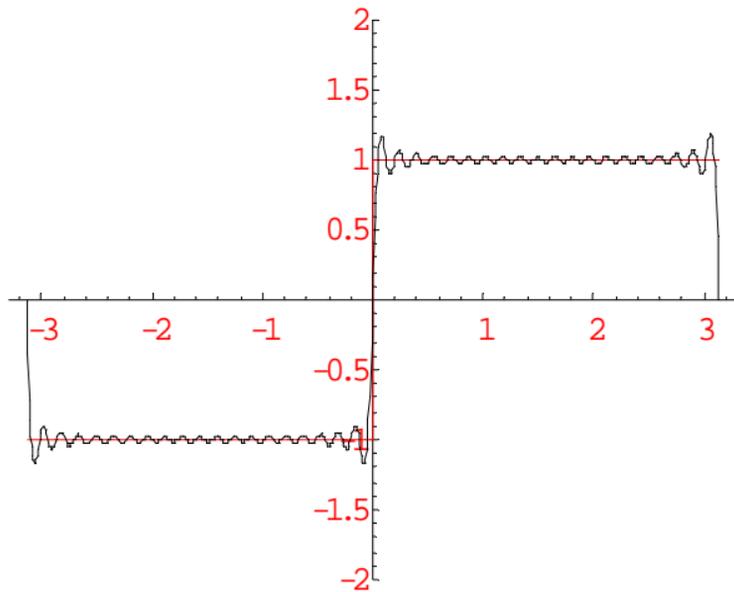


fig 4

Note there are larger errors appearing near “singularity” (discontinuity).

Specifically: “overshoot” of about 9% of the jump near singularity no matter how many terms we take!

In general, singularities (discontinuities in  $f(x)$  or derivatives) cause high frequency components so that FS

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx} / \sqrt{2\pi}$$

has large  $c_n$  for  $n$  large (bad for convergence).

But notice that singularities are at only **one point**, but cause all  $c_n$  to be large.

Wavelets can deal with problem of localization of singularities, since they are localized.

## Advantages of FS:

- “Frequency content” displayed in sizes of the coefficients  $a_k$  and  $b_k$ .
- Easy to write derivatives of  $f$  in terms of series (and use to solve differential equations)

**Fourier series are a natural for differentiation.**

Equivalently sines and cosines are eigenvectors of the derivative operator  $\frac{d^2}{dx^2}$ .

## Disadvantages:

- Usual Fourier transform or series not well-adapted for time-frequency analysis (i.e., if high frequencies are there, we have large  $a_k$  and  $b_k$  for  $k = 100$ . But what part of the function has the high frequencies?

Where  $x < 0$ ? Where  $2 < x < 3$ ?

## Possible solution:

Sliding Fourier transform -

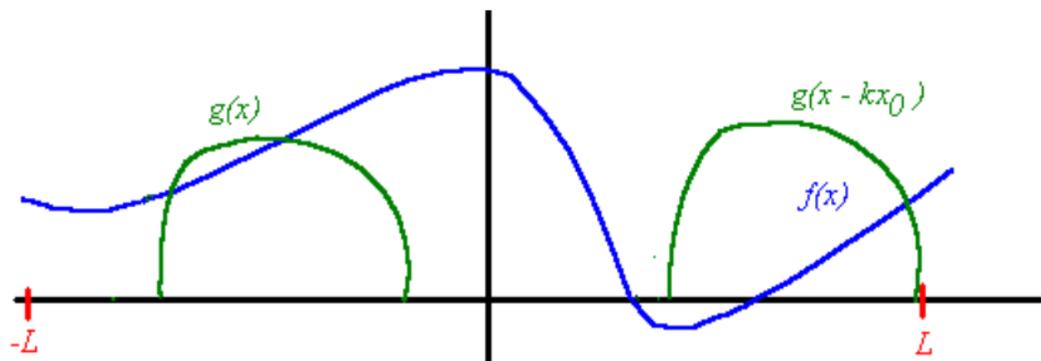


fig 5

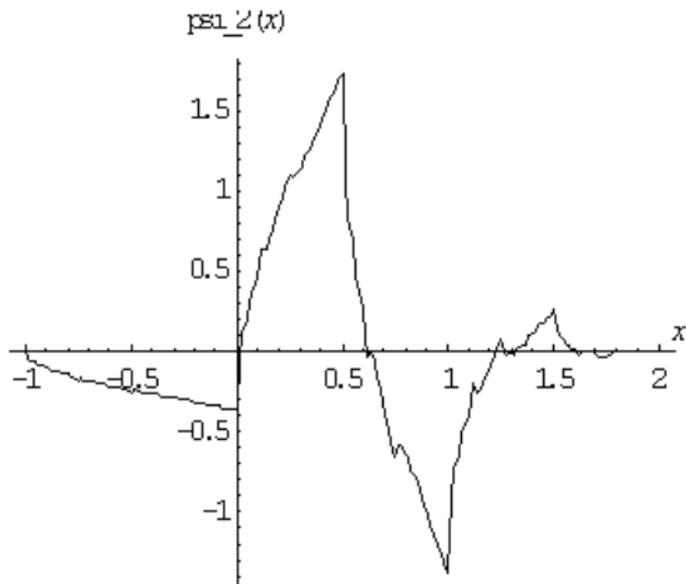
Thus first multiply  $f(x)$  by “window”  $g(x - kx_0)$ , and look at Fourier series or take Fourier transform: look at

$$\int_{-L}^L dx f(x) g_{jk}(x) = \int_{-L}^L dx f(x) g(x - kx_0) e^{ij\frac{\pi}{L}x} \equiv c_{jk}$$

Note however: functions  $g_{jk}(x) = g(x - kx_0) e^{ij\frac{\pi}{L}x}$  not orthonormal like sines and cosines; do not form a nice basis – need something better.

## 5. Wavelet transform

Try: Wavelet transform - fix appropriate function  $h(x)$ .



Then form all translations by integers, and all 'scalings' by powers of 2:

$$h_{jk}(x) = 2^{j/2} h(2^j x - k)$$

( $2^{j/2}$  = normalization constant)

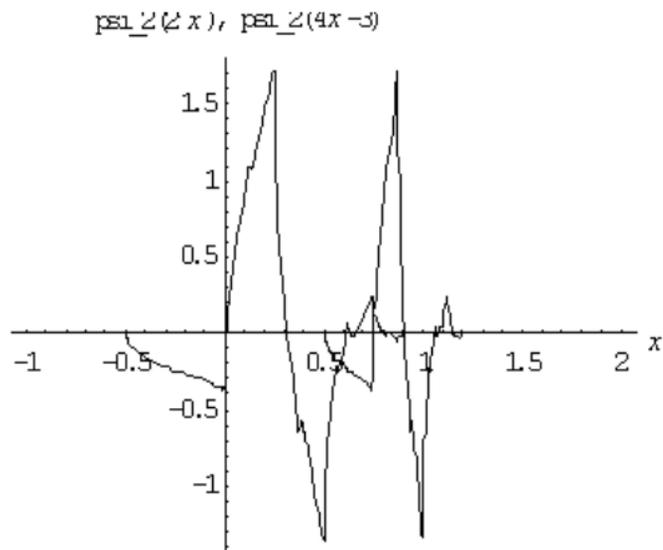


fig. 6:  $h(2x)$  and  $h(4x - 3)$

Let

$$c_{jk} = \int dx f(x) h_{jk}(x).$$

If  $h$  chosen properly, then can get back  $f$  from the  $c_{jk}$ :

$$f(x) = \sum_{j,k} c_{jk} h_{jk}(x)$$

These new functions and coefficients are easier to manage. Sometimes much better –

## Advantages over windowed Fourier transform:

- Coefficients  $c_{jk}$  are all real
- For high frequencies ( $j$  large), functions  $h_{jk}(t)$  have good localization (get thinner as  $j \rightarrow \infty$ ; above diagram). Thus short lived (i.e. of small duration in  $x$ ) high frequency components can be seen from wavelet analysis, but not from windowed Fourier transform.

Note  $h_{jk}$  has width of order  $2^{-j}$ , and is centered about  $k2^{-j}$  (see diagram earlier).

## DISCRETE WAVELET EXPANSIONS:

Take a basic function  $h(x)$  (the basic wavelet);

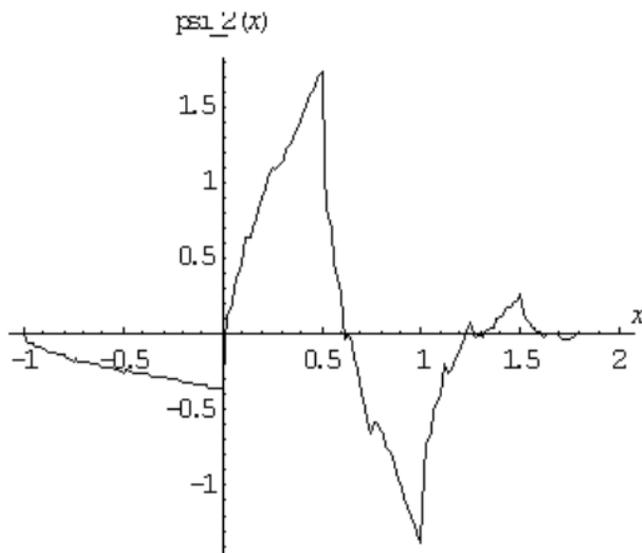


fig 7

let

$$h_{jk}(x) = 2^{j/2} h(2^j x - k).$$

Form discrete wavelet coefficients:

$$c_{jk} = \int dx f(x) h_{jk}(x) \equiv \langle f, h_{jk} \rangle.$$

## Questions:

- Do  $c_{jk}$  characterize  $f$ ?
- Can we expand  $f$  in an expansion of the  $h_{jk}$ ?
- What properties must  $h$  have for this to happen?
- How can we reconstruct  $f$  in a numerically stable way from knowing  $c_{jk}$ ?

*We will show: It is possible to find a function  $h$  such that the functions  $h_{jk}$  form such a perfect basis for functions on  $\mathbb{R}$ .*

That is,  $h_{jk}$  are orthonormal:

$$\langle h_{jk}, h_{j'k'} \rangle \equiv \int h_{jk}(x) h_{j'k'}(x) dx = 0$$

unless  $j = j'$  and  $k = k'$ .

And any function  $f(x)$  can be represented by the  $h_{jk}$ :

$$f(x) = \sum_{j,k} c_{jk} h_{jk}(x).$$

So: like Fourier series, but  $h_{jk}$  are better (e.g., non-zero only on a small sub-interval, i.e., compactly supported)



## 6. A SIMPLE EXAMPLE: HAAR WAVELETS

Motivation: suppose have basic function

$$\phi(x) = \begin{cases} 1 & \text{if } 0 \leq x \leq 1 \\ 0 & \text{otherwise} \end{cases} = \text{basic "pixel"}.$$

We wish to build all other functions out of pixel and translates  $\phi(x - k)$

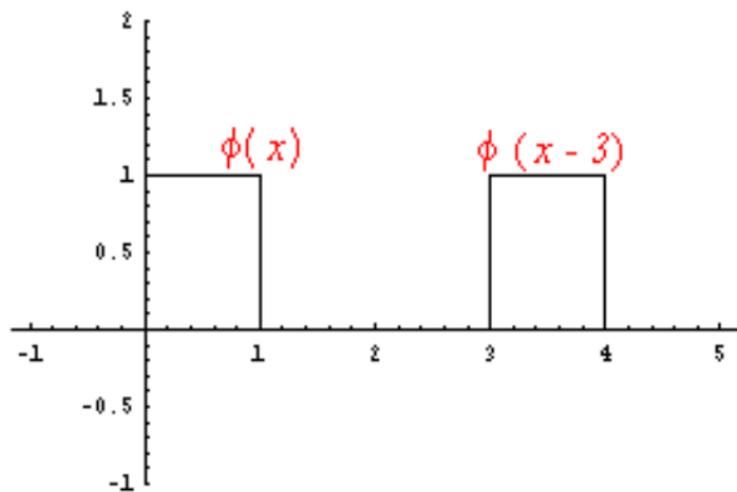


fig 8:  $\phi$  and its translates

Linear combinations of the  $\phi(x - k)$ :

$$f(x) = 2\phi(x) + 3\phi(x - 1) - 2\phi(x - 2) + 4\phi(x - 3)$$

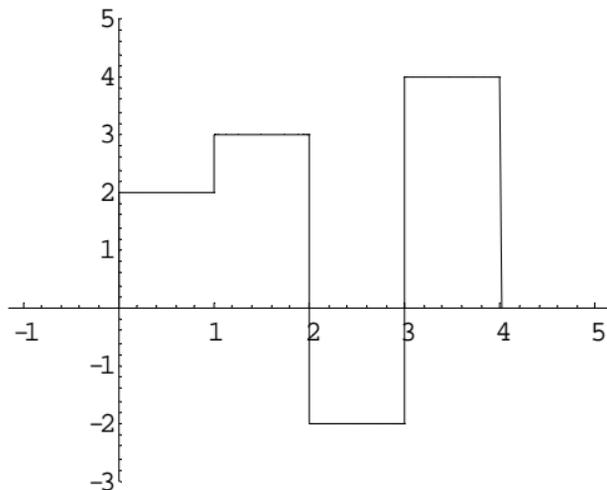


fig 9: linear combination of  $\phi(x - k)$

[Note that any function which is constant on the integers  
can be written in such a form:]

Given function  $f(x)$ , approximate  $f(x)$  by a linear combination of  $\phi(x - k)$ :

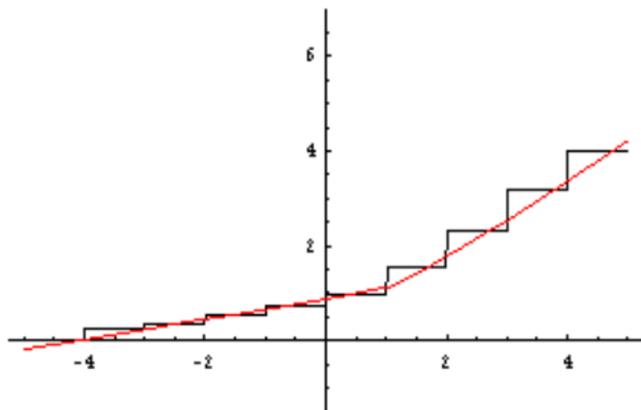


fig 10: approximation of  $f(x)$  using the pixel  $\phi(x)$  and its translates.

Define  $V_0$  = all square integrable functions of the form

$$g(x) = \sum_k a_k \phi(x - k)$$

= all square integrable functions which are constant on  
integer

intervals

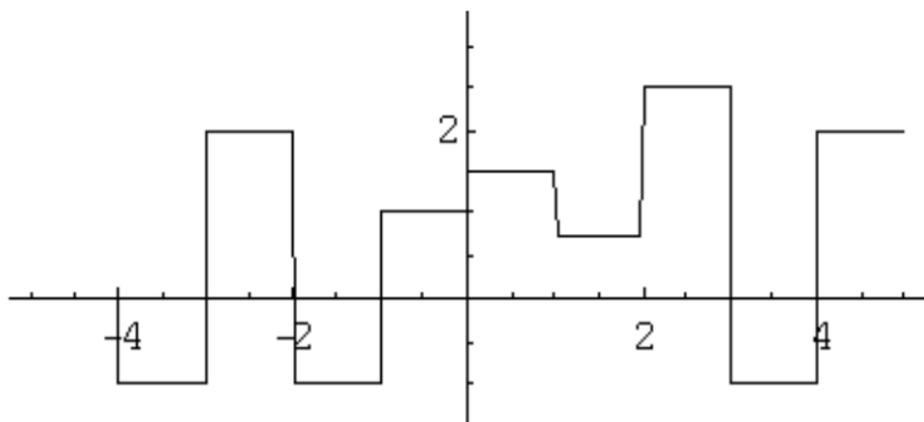


fig 11: a function in  $V_0$

To get better approximations, shrink the pixel :

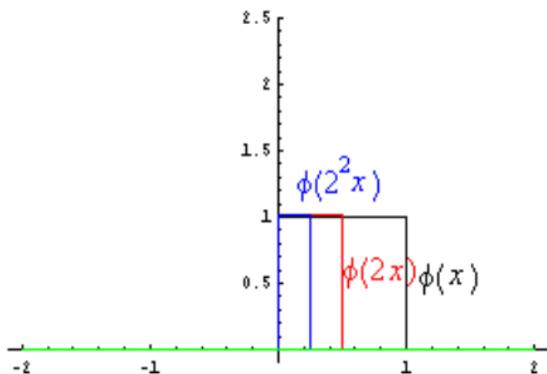


fig 12:  $\phi(x)$ ,  $\phi(2x)$ , and  $\phi(2^2 x)$

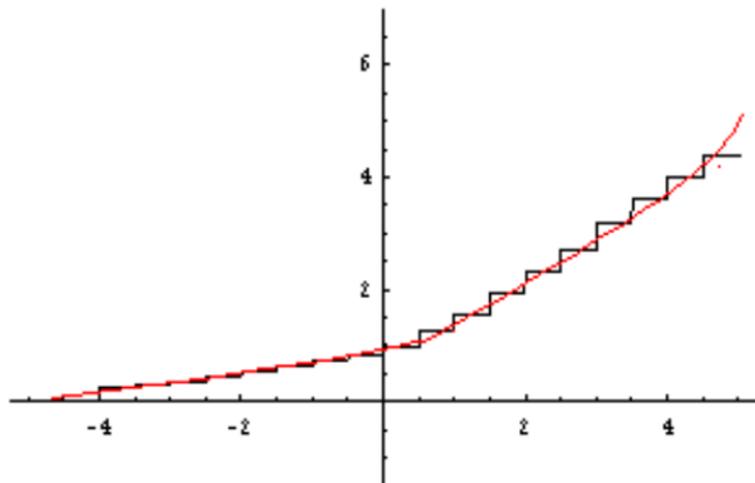


fig 13: approximation of  $f(x)$  by translates of  $\phi(2x)$ .

Define

$V_1$  = all square integrable functions of the form

$$g(x) = \sum_k a_k \phi(2x - k)$$

= all square integrable functions which are constant on all half-integers

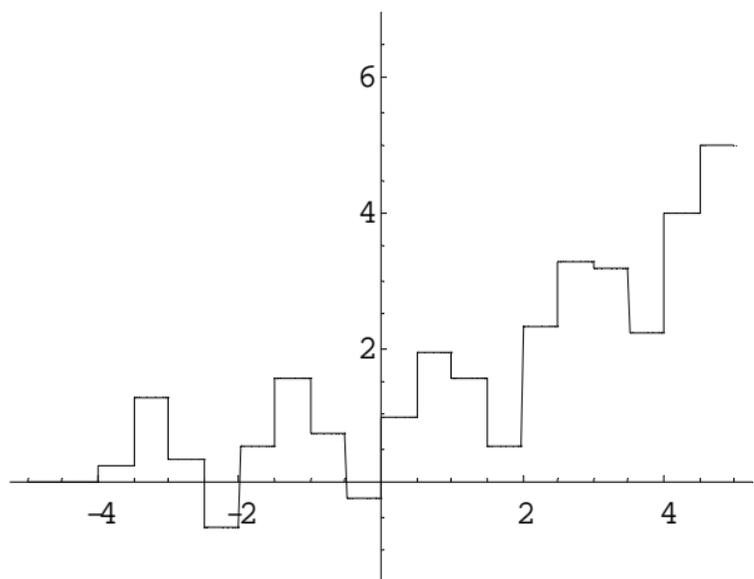


fig 14: Function in  $V_1$

Define  $V_2 =$  sq. int. functions

$$g(x) = \sum_k a_k \phi(2^2 x - k)$$

= sq. int. fns which are constant on quarter integer intervals

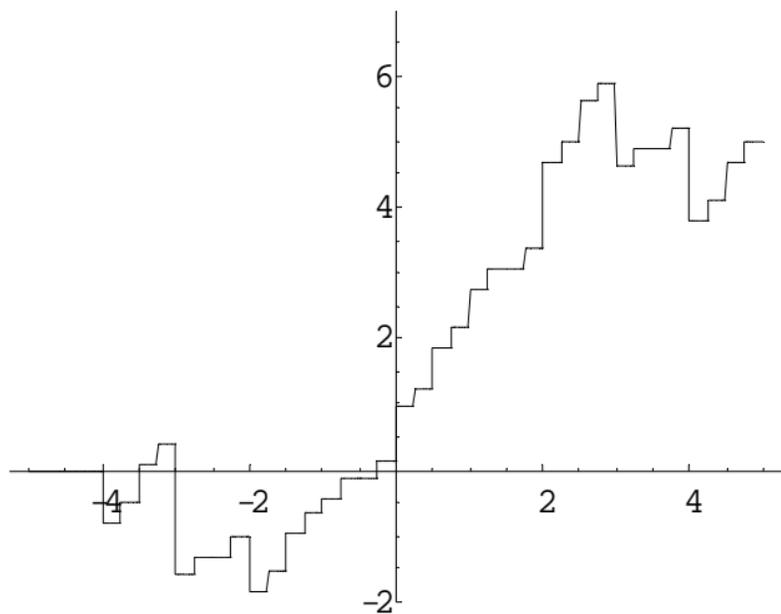


fig 15: function in  $V_2$

Generally define  $V_j$  = all square integrable functions of the form

$$g(x) = \sum_k a_k \phi(2^j x - k)$$

= all square integrable functions which are constant on  $2^{-j}$  length intervals

[note if  $j$  is negative the intervals are of length greater than 1].