

More on beats and resonance

Last class we discussed the solutions to

$$\frac{d^2y}{dt^2} + 3y = \cos \omega t$$

that satisfy the initial condition  $(y(0), y'(0)) = (0, 0)$  where  $\omega$  is a parameter. If  $\omega \neq \pm\sqrt{3}$ , the solution is

$$y(t) = \frac{1}{3 - \omega^2}(\cos \omega t - \cos \sqrt{3}t).$$

Applying a trig identity, we obtain

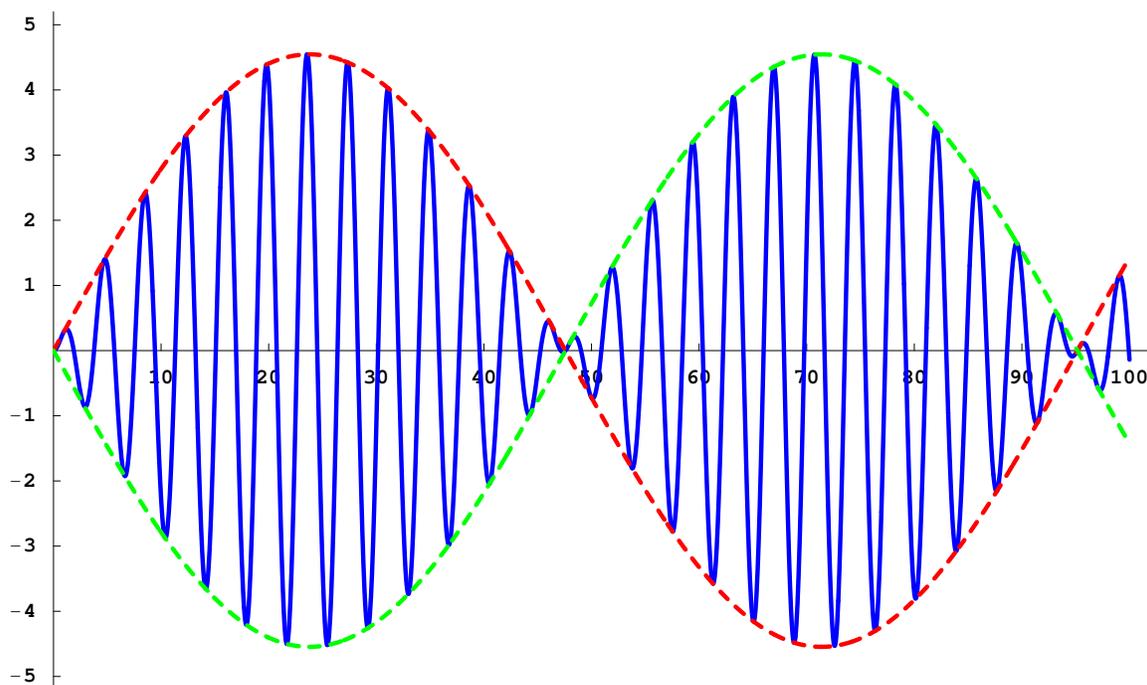
$$y(t) = \frac{-2}{3 - \omega^2} (\sin \alpha t) (\sin \beta t)$$

where

$$\alpha = \frac{\omega + \sqrt{3}}{2} \quad \text{and} \quad \beta = \frac{\omega - \sqrt{3}}{2}.$$

Here is the graph of this solution in the case where  $\omega = 1.6$ . Note that the average of  $\omega$  and  $\sqrt{3}$  in this case is approximately 1.67. The half difference is approximately  $-0.066$ . The average yields “rapid” oscillations with a period of approximately 3.76. The half-difference yields “slow” oscillations with a period of approximately 95. Also,

$$\frac{2}{3 - 1.6^2} \approx 4.55.$$



What happens if  $\omega = \sqrt{3}$ ?

**Example.**  $\frac{d^2y}{dt^2} + 3y = \cos \sqrt{3}t$ .

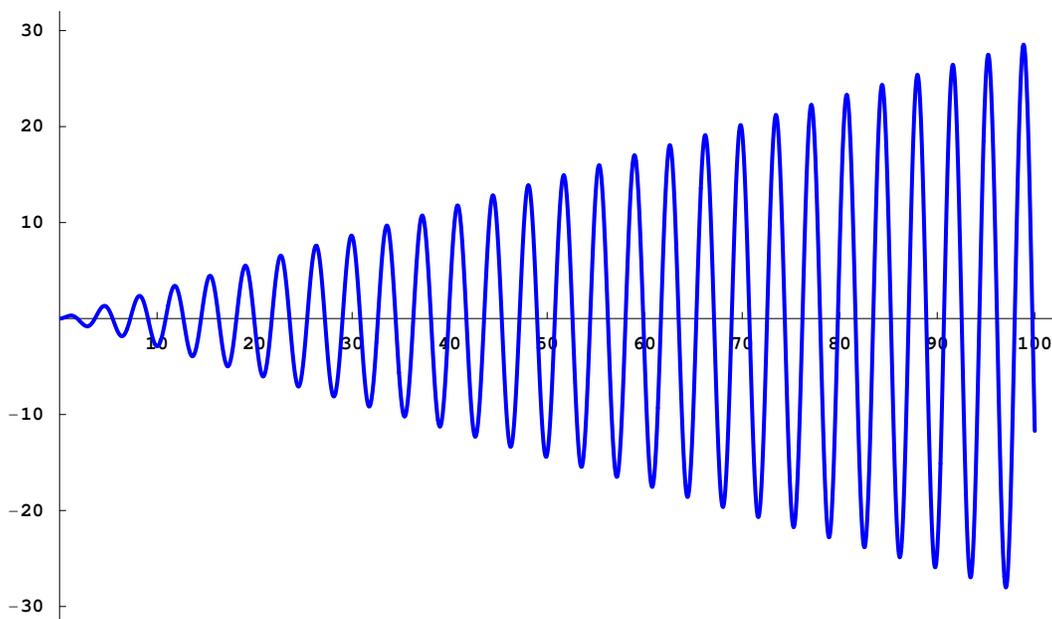
The complexified equation is  $\frac{d^2y}{dt^2} + 3y = e^{i\sqrt{3}t}$ . We guess  $y_c(t) = ate^{i\sqrt{3}t}$ , and we get

$$a = \frac{1}{2i\sqrt{3}} = -\frac{1}{2\sqrt{3}}i.$$

Consequently, if  $\omega = \sqrt{3}$ , the general solution is

$$y(t) = k_1 \cos \sqrt{3}t + k_2 \sin \sqrt{3}t + \frac{1}{2\sqrt{3}}t \sin \sqrt{3}t.$$

Here is the graph for the case where  $k_1 = k_2 = 0$ .



This value of  $\omega$  is called the resonant value for the frequency of the forcing.

The resonance value of the forcing should be immediately apparent from the differential equation.

**Example.** What is the resonance value of  $\omega$  for the one-parameter family of equations

$$\frac{d^2y}{dt^2} + 5y = 4 \cos \omega t?$$

## Linearization

We would like to apply what we know about linear systems to nonlinear systems.

**Example.** Consider the van der Pol equation

$$\frac{d^2x}{dt^2} + (x^2 - 1)\frac{dx}{dt} + x = 0.$$

The corresponding system is

$$\begin{aligned}\frac{dx}{dt} &= y \\ \frac{dy}{dt} &= (1 - x^2)y - x.\end{aligned}$$

The only equilibrium point for this system is  $(0, 0)$ . What is the linearized system near  $(0, 0)$ ?

**Example.** Consider the (undamped) pendulum

$$\frac{d^2\theta}{dt^2} + \sin\theta = 0.$$

The corresponding system is

$$\begin{aligned}\frac{d\theta}{dt} &= v \\ \frac{dv}{dt} &= -\sin\theta.\end{aligned}$$

There are equilibria at  $(\theta, v) = (k\pi, 0)$  for all integers  $k$ .

The linearized system near  $(0, 0)$  is

$$\frac{d\mathbf{Y}}{dt} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \mathbf{Y}.$$

What is the linearized pendulum near the equilibrium point  $(\pi, 0)$ ?

Given the (nonlinear) system

$$\begin{aligned}\frac{dx}{dt} &= f(x, y) \\ \frac{dy}{dt} &= g(x, y),\end{aligned}$$

its **Jacobian** at the point  $(x_0, y_0)$  is the matrix

$$\mathbf{J}(x_0, y_0) = \begin{pmatrix} \frac{\partial f}{\partial x}(x_0, y_0) & \frac{\partial f}{\partial y}(x_0, y_0) \\ \frac{\partial g}{\partial x}(x_0, y_0) & \frac{\partial g}{\partial y}(x_0, y_0) \end{pmatrix}$$

and its linearization at  $(x_0, y_0)$  is the system

$$\frac{d\mathbf{Y}}{dt} = \mathbf{J}\mathbf{Y}.$$

For the pendulum, we have one linearization for each equilibrium point:

For the van der Pol system, we obtain the linearization:

**Linearization Theorem** Let  $\mathbf{Y}_0$  be an equilibrium point for the nonlinear autonomous system

$$\frac{d\mathbf{Y}}{dt} = \mathbf{F}(\mathbf{Y})$$

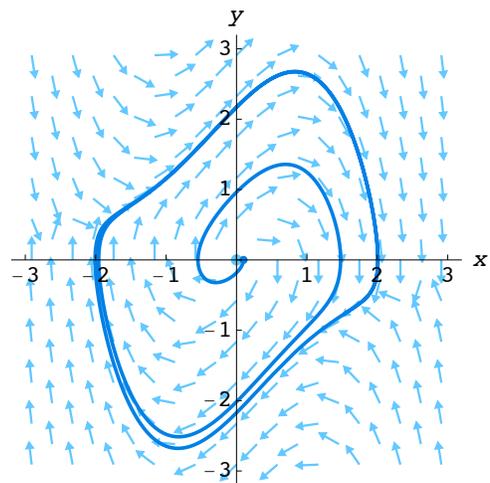
and let

$$\frac{d\mathbf{Y}}{dt} = \mathbf{J}\mathbf{Y}$$

be the corresponding linearized system. If  $\det \mathbf{J} \neq 0$  and the eigenvalues of  $\mathbf{J}$  are not purely imaginary, then the solution curves of the nonlinear system near  $\mathbf{Y}_0$  behave in the same qualitative way as the solution curves of the linear system.

**Example.** Consider the van der Pol system near the origin. The linearized system is

$$\frac{d\mathbf{Y}}{dt} = \begin{pmatrix} 0 & 1 \\ -1 & 1 \end{pmatrix} \mathbf{Y}.$$



**Example.** Consider the pendulum system. The linearized system near  $(\pi, 0)$  is

$$\frac{d\mathbf{Y}}{dt} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \mathbf{Y}.$$

Its characteristic equation is  $\lambda^2 - 1 = 0$ , and therefore the eigenvalues are  $\pm 1$ . The Linearization Theorem says that this equilibrium point is a nonlinear saddle.

The linearized system near  $(0, 0)$  is

$$\frac{d\mathbf{Y}}{dt} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \mathbf{Y},$$

and the characteristic equation is  $\lambda^2 + 1 = 0$ . The eigenvalues are  $\pm i$ . This equilibrium point is a nonlinear center, but this example is misleading. The Linearization Theorem does not apply to this equilibrium point.

