

The McMullen Domain: Rings Around the Boundary

Robert L. Devaney
Sebastian M. Marotta

Boston University

May 28, 2005

Our goal in this paper is to consider the dynamics of families of rational maps of the form

$$F_\lambda(z) = z^n + \frac{\lambda}{z^n}$$

where $\lambda \neq 0$ is a complex parameter and n is a positive integer. The Julia sets corresponding to maps in these families have been shown to possess a number of interesting dynamical and topological properties.

In this paper we discuss some of the properties of the parameter plane for these maps. Each of these maps has $2n$ “free” critical points. However, like the well-studied quadratic family $Q_c(z) = z^2 + c$, each of these families has only one free critical orbit since all forward orbits of the critical points behave symmetrically. Hence the λ -plane is a natural parameter plane for these families.

As another similarity with the quadratic family, the point at ∞ is super-attracting fixed point for each λ when $n > 1$, and so it may be the case that the critical orbits enter the basin of this fixed point. Unlike the quadratic family, there are three distinctly different manners in which the critical orbit may escape to ∞ , and this in turn determines the topological structure of the Julia sets for the escape parameters. We denote the immediate basin of attraction of ∞ by B_λ . One possible escape scenario is that the critical values all lie in B_λ . Since 0 is a pole, there is a neighborhood of 0 that is mapped into B_λ . Now either this neighborhood is itself contained in B_λ , or else 0 lies in a disjoint preimage of B_λ which we denote by T_λ . In the latter case, we note that F_λ maps T_λ in n to 1 fashion onto B_λ while $F_\lambda|_{B_\lambda}$ is also n to 1. Hence the only preimages of B_λ are B_λ itself and T_λ . A second possible scenario occurs when the critical values lie in T_λ . And a third possibility arises when some higher iterate of the critical values lies in T_λ . The following Theorem describes the Julia sets that result from these three different situations (see [5]):

Theorem (The Escape Trichotomy). *For the family of functions*

$$F_\lambda(z) = z^n + \frac{\lambda}{z^n}$$

with $n \geq 2$ and $\lambda \in \mathbb{C}$:

- 1. If the critical values lie in B_λ , then the Julia set is a Cantor set.*
- 2. If the critical values lie in $T_\lambda \neq B_\lambda$, then the Julia set is a Cantor set of simple closed curves.*
- 3. If the critical values lie in any other preimage of T_λ , then the Julia set is a Sierpinski curve.*

A *Sierpinski curve* is a planar set that is characterized by the following five properties: it is a compact, connected, locally connected and nowhere dense set whose complementary domains are bounded by simple closed curves that are pairwise disjoint. It is known from work of Whyburn [15] that any two Sierpinski curves are homeomorphic. In fact, they are homeomorphic to the well-known Sierpinski carpet fractal. From the point of view of topology, a Sierpinski curve is a universal set in the sense that it contains a homeomorphic copy of any planar, compact, connected, one-dimensional set. The first example of a Sierpinski curve Julia set was given by Milnor and Tan Lei [11]. See also Ushiki [14].

Case 2 of the Escape Trichotomy was first observed by McMullen [8], who showed that this phenomenon occurs in each family provided that $n \neq 1, 2$ and λ is sufficiently small.

In Figure 1 we display the parameter plane for the family $F_\lambda(z) = z^3 + \lambda/z^3$. The black regions indicate parameter values for which the critical orbit does not escape to ∞ . Again in analogy with the quadratic polynomial family, for these parameters the Julia set is a connected set. The white

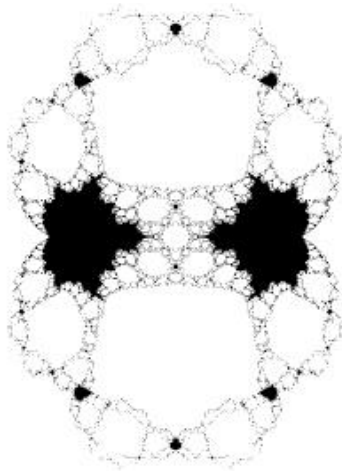


Figure 1: The parameter plane for the family $z^3 + \lambda/z^3$.

regions in this picture represent λ -values for which the critical orbit tends to ∞ . The exterior region corresponds to parameter values for which the Julia set is a Cantor set; we call this set the *Cantor set locus*. The small region in the center corresponds to parameter values for which the Julia set is a Cantor set of simple closed curves. We call this region the *McMullen domain*. The other white regions correspond to parameters for which the Julia set is a Sierpinski curve. These are called *Sierpinski holes*.

It is known that there are infinitely many disjoint Sierpinski holes for each of these families [1], [4] and that there is a parameter in each Sierpinski hole for which the orbit of the critical point lands on 0 at some iteration and therefore on ∞ at the next iteration, say at iteration $k > 2$. We then call this λ -value the center of the Sierpinski hole and k the *escape time* of the hole. All other parameters in a given Sierpinski hole have the property that the critical value has orbit that lands in B_λ at the escape time iterate. By Whyburn's result, the Julia sets corresponding to any two parameters drawn from a Sierpinski hole are homeomorphic. However, as shown in [5],

there exist Sierpinski holes corresponding to each escape time $k \geq 3$, and these have the property that if λ_1 and λ_2 lie in Sierpinski holes with different escape times, then F_{λ_1} and F_{λ_2} are not topologically conjugate on their Julia sets.

Note that, in the case $n = 3$, there appear to be two large copies of a Mandelbrot set that straddle the positive and negative real axes. These are called the *principal Mandelbrot sets* for F_λ . Using the Douady-Hubbard theory of polynomial-like maps, it is known that these sets are actually homeomorphic to the standard quadratic Mandelbrot set [2] and that, for parameters drawn from these sets, there is an invariant subset on which F_λ is conjugate to the corresponding quadratic polynomial on its Julia set. Also apparent in this image are two large Sierpinski holes along the positive and negative imaginary axis. These holes have escape time 3.

Our goal in this paper is to investigate further properties of the parameter plane for these maps and, in particular, the structure of the parameter plane in a neighborhood of the McMullen domain. It is known [3] that there is a unique McMullen domain in the parameter plane for each $n \geq 3$, and this region is an open disk that is bounded by a simple closed curve.

In Figure 2, we have displayed several magnifications of the region around the McMullen domain in the case $n = 3$. In the first image, note that there are four large Sierpinski holes symmetrically placed around the McMullen domain. These Sierpinski holes all have escape time 4. Between the two upper and the two lower Sierpinski holes there appear to be small copies of a Mandelbrot set, while between the two left and two right holes we see the period two bulb of a principal Mandelbrot set and the remainder of the “tail” of this set. Indeed, one may draw a simple closed curve that encircles the McMullen domain and passes through the centers of each of these Sierpinski holes, the centers of the main cardioids of the two smaller Mandelbrot sets,

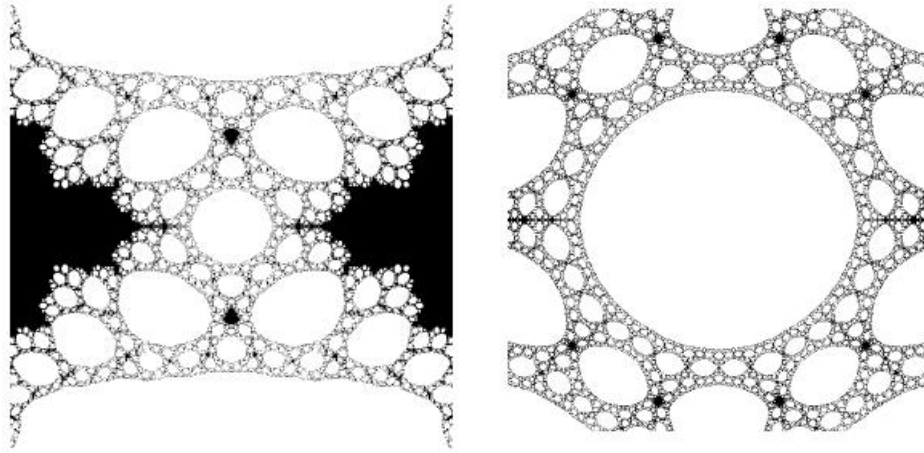


Figure 2: Magnifications of the parameter plane for the family $z^3 + \lambda/z^3$ around the McMullen domain.

and the centers of the two period two bulbs of the principal Mandelbrot sets. That is, on this simple closed curve, we find four parameter values for which F_λ^2 has a superstable fixed point and four other values for which F_λ^4 maps the critical points to ∞ , and these parameter values alternate between the superstable and the centers of Sierpinski holes as the parameter winds around the closed curve.

Inside these four Sierpinski holes appear to be another simple closed curve containing ten Sierpinski holes. Each of these holes has escape time 5. Also, each pair of these holes apparently has either a small copy of a Mandelbrot set or a portion of a principal Mandelbrot set between them. Examining the further magnification in Figure 2, we see a smaller closed curve containing 28 Sierpinski holes with escape time 6 and, inside that curve, an even smaller curve containing 82 Sierpinski holes with escape time 7. It appears that the k^{th} curve meets exactly $3^k + 1$ Sierpinski holes with escape time $k + 3$ as well as the same number of (portions of) Mandelbrot sets. We call these curves

rings around the McMullen domain.

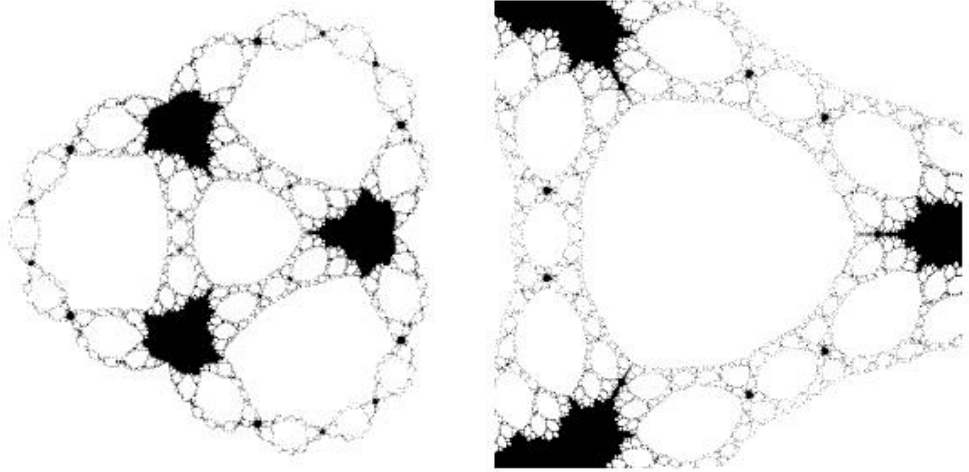


Figure 3: The parameter plane for the family $z^4 + \lambda/z^4$ and a magnification around the McMullen domain.

Actually, the formula in the general case is a little more complicated than that. In Figure 3 we display the parameter plane for the case $n = 4$ as well as a magnification of the McMullen domain. Here we see three principal Mandelbrot sets caught between three large Sierpinski holes, each of which has escape time 3. Inside these sets is a curve containing 9 Sierpinski holes, each with escape time 4, and inside another curve containing 33 holes of escape time 5. Further magnification shows that there are $2 \cdot 4^{k-1} + 1$ holes with escape time $k + 2$ in case $n = 4$.

Our main goal in this paper is to make these observations rigorous. We shall prove:

Theorem. (Rings Around the McMullen Domain.) *For each $n \geq 3$, the McMullen domain for the family $z^n + \lambda/z^n$ is surrounded by infinitely many simple closed curves \mathcal{S}^k for $k = 1, 2, \dots$ having the property that:*

1. Each curve \mathcal{S}^k surrounds the McMullen domain as well as \mathcal{S}^{k+1} , and the \mathcal{S}^k accumulate on the boundary of the McMullen domain as $k \rightarrow \infty$;
2. The curve \mathcal{S}^k meets the centers of τ_k^n Sierpinski holes, each with escape time $k + 2$ where

$$\tau_k^n = (n - 2)n^{k-1} + 1.$$

3. The curve \mathcal{S}^k also passes through τ_k^n superstable parameter values where a critical point is periodic of period k or $2k$.

In a subsequent paper we shall show that these superstable parameter values each lie at the center of a small copy of a Mandelbrot set, while the Sierpinski holes surrounding the centers are all simply connected sets.

1 Elementary Mapping Properties

In this paper we restrict attention to the family of rational maps given by

$$F_\lambda(z) = z^n + \frac{\lambda}{z^n}$$

where, in view of McMullen's result, we assume that $n \geq 3$.

In the dynamical plane, the object of principal interest is the *Julia set* of F_λ , which we denote by $J(F_\lambda)$. The Julia set is the set of points at which the family of iterates $\{F_\lambda^n\}$ fails to be a normal family in the sense of Montel. It is known that $J(F_\lambda)$ is also the closure of the set of repelling periodic points for F_λ as well as the boundary of the set of points whose orbits escape to ∞ under iteration of F_λ . See [10].

The point at ∞ is a superattracting fixed point for F_λ and we denote the immediate basin of ∞ by B_λ . It is well known that F_λ is conjugate to $z \mapsto z^n$ in a neighborhood of ∞ in B_λ [12]. There is also a pole of order n for F_λ at the origin, so there is a neighborhood of 0 that is mapped into B_λ by F_λ . If

the full preimage of B_λ that contains this neighborhood is disjoint from B_λ , then we denote the preimage of B_λ that contains 0 by T_λ . So F_λ maps both B_λ and T_λ in n -to-one fashion over B_λ . We call T_λ the *trap door* since any orbit that eventually enters the immediate basin of ∞ must “fall through” T_λ enroute to B_λ .

Besides 0 and ∞ , F_λ has $2n$ other critical points given by $\lambda^{1/2n}$. We call these points the *free critical points* for F_λ . There are, however, only two critical values, and these are given by $\pm 2\sqrt{\lambda}$. We denote a free critical point by c_λ and a critical value by v_λ . The map also has $2n$ prepoles given by $(-\lambda)^{1/2n}$. Note that all of the critical points and prepoles lie on the circle of radius $|\lambda|^{1/2n}$ centered at the origin. We call this circle the *critical circle* and denote it by C_λ .

The map F_λ has some very special properties when restricted to circles centered at the origin. The following is a straightforward computation (see [3]):

Proposition.

1. F_λ takes the critical circle $2n$ -to-one onto the line interval connecting the two critical values $\pm\sqrt{2\lambda}$;
2. F_λ takes any other circle centered at the origin to an ellipse whose foci are the critical values.

We call the image of the critical circle the *critical segment*. We call the straight line connecting the origin to ∞ and passing through one of the critical points (resp., prepoles) a *critical point ray* (resp., *prepole ray*). Similar straightforward computations show that each of the critical point rays is mapped in two-to-one fashion onto one of the two straight line segments of the form tv_λ , where $t \geq 1$ and v_λ is the image of the critical point on this

ray. So the image of a critical point ray is a straight ray connecting either v_λ or $-v_\lambda$ to ∞ . Thus the critical segment together with these two rays forms a straight line through the origin.

Similarly, each of the $2n$ prepole rays is mapped in one-to-one fashion onto the straight line given by $it\sqrt{\lambda}$, where t is now any real number. Note that the image of the prepole rays is the line that is perpendicular to the line tv_λ for $t \in \mathbb{R}$, i.e., the line that contains the critical segment and the images of all of the critical point rays.

Let U_λ be a sector bounded by two prepole rays corresponding to adjacent prepoles on C_λ , i.e., U_λ is a sector in the plane with angle $2\pi/2n$. We call U_λ a *critical point sector* since it contains at its “center” a unique critical point of F_λ . Similarly, let V_λ be the sector bounded by two critical point rays corresponding to adjacent critical points on C_λ . We call V_λ a *prepole sector*. The next result follows immediately from the above:

Proposition (Mapping Properties of F_λ).

1. F_λ maps the interior of each critical point sector in two-to-one fashion onto the open half plane that is bounded by the image of the prepole rays and contains the critical value that is the image of the unique critical point in the sector;
2. F_λ maps the interior of each prepole sector in one-to-one fashion onto the entire plane minus the two half lines $\pm tv_\lambda$ where $t \geq 1$;
3. F_λ maps the region in either the interior or the exterior of the critical circle onto the complement of the critical segment as an n -to-one covering map of $\overline{\mathbb{C}}$ (except at 0 and ∞).

We now turn to the symmetry properties of F_λ in both the dynamical and parameter planes. Let ν be the primitive $2n^{\text{th}}$ root of unity given by

$\exp(\pi i/n)$. Then, for each j , we have $F_\lambda(\nu^j z) = (-1)^j F_\lambda(z)$. Hence, if n is even, we have $F_\lambda^2(\nu^j z) = F_\lambda^2(z)$ for each j . Therefore the points z and $\nu^j z$ land on the same orbit after two iterations and so their orbits have the same eventual behavior for each j . If n is odd, the orbits of $F_\lambda(z)$ and $F_\lambda(\nu^j z)$ are either the same or else they are the negatives of each other after the first iteration. In either case it follows that the orbits of $\nu^j z$ behave symmetrically under $z \mapsto -z$ for each j . Hence the Julia set of F_λ is always symmetric under $z \mapsto \nu z$. In particular, each of the free critical points eventually maps onto the same orbit (in case n is even) or onto one of two symmetric orbits (in case n is odd). Thus these orbits all have the same behavior and so the λ -plane is a natural parameter plane for each of these families. Note also that, if n is even and the orbit of some critical point eventually lands on some other critical point at iteration $j \geq 1$, then in fact one of the critical points of F_λ must be periodic of period j . If n is odd, then there are two possibilities: either one of the critical points has period j or else it has period $2j$.

Let $H_\lambda(z)$ be one of the n involutions given by $H_\lambda(z) = \lambda^{1/n}/z$. Then we have $F_\lambda(H_\lambda(z)) = F_\lambda(z)$, so that the Julia set is also preserved by each of these involutions. Note that each H_λ maps the critical circle to itself and also fixes a pair of critical points of the form $\pm \sqrt{\lambda^{1/n}}$. H_λ also maps circles centered at the origin outside the critical circle to similar circles inside the critical circle and vice versa. It follows that two such circles, one inside and one outside the critical circle, are mapped onto the same ellipse by F_λ .

The parameter plane (see Figures 1 and 3) for F_λ also possesses several symmetries. First of all, we have

$$\overline{F_\lambda(z)} = F_{\overline{\lambda}}(\overline{z})$$

so that F_λ and $F_{\overline{\lambda}}$ are conjugate via the map $z \mapsto \overline{z}$. Therefore the parameter plane is symmetric under the map $\lambda \mapsto \overline{\lambda}$.

We also have $(n - 1)$ -fold symmetry in the parameter plane for F_λ . To see this, let ω be the primitive $(n - 1)$ st root of unity given by $\exp(2\pi i/(n - 1))$. Then, if n is even, we compute that

$$F_{\lambda\omega}(\omega^{n/2}z) = \omega^{n/2}(F_\lambda(z)).$$

As a consequence, for each $\lambda \in \mathbb{C}$, the maps F_λ and $F_{\lambda\omega}$ are conjugate under the linear map $z \mapsto \omega^{n/2}z$. In particular, since, when λ is real, the real line is preserved by F_λ , it follows that the straight line passing through 0 and $\omega^{n/2}$ is preserved by $F_{\lambda\omega}$.

When n is odd, the situation is a little different. We now have

$$F_{\lambda\omega}(\omega^{n/2}z) = -\omega^{n/2}(F_\lambda(z)).$$

Since $F_\lambda(-z) = -F_\lambda(z)$ when n is odd, we therefore have that $F_{\lambda\omega}^2$ is conjugate to F_λ^2 via the map $z \mapsto \omega^{n/2}z$. This means that the dynamics of F_λ and $F_{\lambda\omega}$ are “essentially” the same, though subtly different. For example, if F_λ has a fixed point, then under the conjugacy, this fixed point and its negative are mapped to a 2-cycle for $F_{\lambda\omega}$. Since the real line is invariant when λ is real, it follows that the straight lines passing through the origin and $\pm\omega^{n/2}$ are interchanged by $F_{\lambda\omega}$ and hence invariant under $F_{\lambda\omega}^2$.

To summarize the symmetry properties of F_λ , we have:

Proposition (Symmetries in the dynamical and parameter plane). *The dynamical plane for F_λ is symmetric under the map $z \mapsto \nu z$ where $\nu = \exp(\pi i/n)$. The parameter plane is symmetric under both $z \mapsto \bar{z}$ and $z \mapsto \omega z$ where $\omega = \exp(2\pi i/(n - 1))$.*

The following result shows that the McMullen domain and all of the Sierpinski holes are located inside the unit circle in parameter space.

Proposition (Location of the Cantor set locus.) *Suppose $|\lambda| \geq 1$. Then v_λ lies in B_λ so that λ lies in the Cantor set locus.*

Proof: Suppose $|z| \geq 2|\lambda|^{1/2} \geq 2$. Then, since $|z| \geq |\lambda|^{1/2}$, we have

$$|F_\lambda(z)| \geq |z|^n - \frac{|\lambda|}{|z|^n} \geq |z|^n - |\lambda|^{1-\frac{n}{2}} \geq |z|^n - 1 \geq |z|^{n-1} > |z|$$

since $n > 2$. Hence $|F_\lambda^j(z)| \rightarrow \infty$ so the region $|z| \geq 2|\lambda|^{1/2}$ lies in B_λ . In particular, $v_\lambda \in B_\lambda$. □

For each n , let $\lambda^* = \lambda_n^*$ be the unique real solution to the equation

$$|v_\lambda| = 2|\sqrt{\lambda}| = |\lambda|^{1/2n} = |c_\lambda|.$$

Using this equation, we compute easily that

$$\lambda^* = \left(\frac{1}{4}\right)^{\frac{n}{n-1}}.$$

The circle of radius λ^* plays an important role in the parameter plane, for if λ lies on this circle, it follows that both of the critical values lie on the critical circle for F_λ . If λ lies inside this circle, then F_λ maps the critical circle strictly inside itself. We call the circle of radius λ^* in parameter plane the *dividing circle*. We denote by $\mathcal{O} = \mathcal{O}_n$ the open set of parameters inside the dividing circle. We will be primarily concerned in later sections with values of the parameter that lie in \mathcal{O} . In particular, we shall show that all of the rings around the McMullen domain \mathcal{S}^k with $k > 1$ lie in this region while the ring \mathcal{S}^1 is the dividing circle itself.

2 Some Special Cases

In this section we discuss the dynamics of several special cases of F_λ that will help define the rings around the McMullen domain later.

First suppose that λ lies on the dividing circle, i.e., $|\lambda| = \lambda^*$. In this case, all of the critical points, critical values, and prepoles of F_λ lie on the same

circle (the critical circle) in dynamical plane, namely the circle

$$|z| = \left(\frac{1}{2}\right)^{\frac{1}{n-1}}.$$

As λ winds once around the dividing circle in the counterclockwise direction beginning on the real axis, the critical points and prepoles of F_λ wind $1/2n$ of a turn around the critical circle, while the critical values wind one-half of a turn around the critical circle, all in the counterclockwise direction. Hence there are exactly $n - 1$ special parameter values on the dividing circle for which a critical point of the corresponding map equals a critical value, so for these special λ -values we have a superattracting fixed or period two point for F_λ . Equivalently, one computes that these $n - 1$ parameters are given by

$$\lambda = \left(\frac{1}{4}\right)^{\frac{n}{n-1}}.$$

There are $n - 1$ other parameters on this circle for which the critical value is a prepole, and these are given by

$$\lambda = \left(\frac{-1}{4}\right)^{\frac{n}{n-1}}.$$

This proves the case $k = 1$ of the Rings Around the McMullen Domain Theorem.

Theorem. *The ring \mathcal{S}^1 is the dividing circle in parameter plane. It contains $n - 1$ superstable parameters and the same number of centers of Sierpinski holes.*

See Figure 4.

We next restrict attention to values of λ lying in \mathbb{R}^+ . The graph of F_λ shows that, in this case, F_λ maps \mathbb{R}^+ to itself and that there is a unique critical point lying in \mathbb{R}^+ . We denote this critical point by $c_0 = c_0(\lambda)$. See Figure 5.

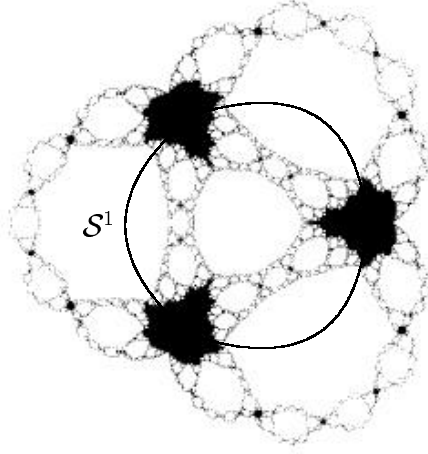


Figure 4: The curve \mathcal{S}^1 in the parameter plane for $n = 4$.

It is known [2] that there is a Mandelbrot set whose central spine lies along the interval $[\lambda_-, \lambda_+]$ contained in \mathbb{R}^+ . Moreover, if $\lambda > \lambda_+$, then λ lies in the Cantor set locus, whereas if $0 < \lambda < \lambda_-$, then λ lies in the McMullen domain. The graph of $F_\lambda | \mathbb{R}^+$ shows that F_λ undergoes a saddle-node bifurcation at λ_+ and that the critical point c_λ maps onto the repelling fixed point in $\partial B_\lambda \cap \mathbb{R}^+$ after two iterations when $\lambda = \lambda_-$. Since each F_λ is conjugate on the real line to a quadratic polynomial of the form $Q_c(x) = x^2 + c$, standard facts from quadratic dynamics yield the following:

Proposition (Superstable parameters for $\lambda \in \mathbb{R}^+$.) *There is a decreasing sequence of parameters in \mathbb{R}^+ $\lambda_1 > \lambda_2 \dots$ converging to λ_- such that, for $\lambda = \lambda_k$, the critical point c_0 is periodic with period k and the critical orbit in \mathbb{R}^+ has the special form when $k \geq 2$:*

$$0 < v_\lambda = F_\lambda(c_0) < c_0 = F_\lambda^k(c_0) < F_\lambda^{k-1}(c_0) < \dots < F_\lambda^3(c_0) < F_\lambda^2(c_0).$$

In particular, λ_k is a superstable parameter value of period k and the orbit of $F_{\lambda_k}^2(c_0)$ is monotonically decreasing for $k - 1$ iterations along \mathbb{R}^+ .

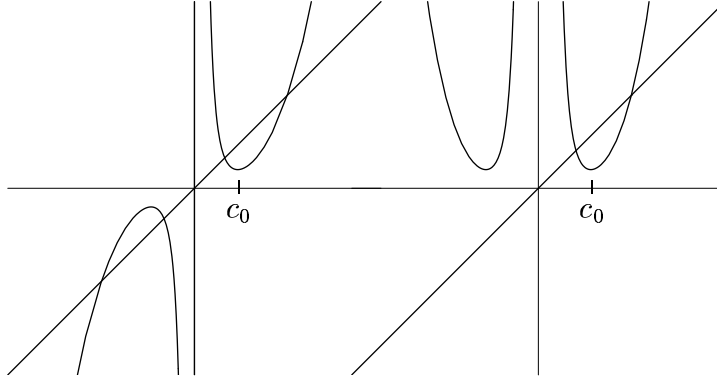


Figure 5: The graphs of $x^3 + 0.01/x^3$ and $x^4 + 0.01/x^4$.

Portions of the graphs of F_{λ_k} for $k = 4$ and $k = 8$ when $n = 4$ are displayed in Figure 6. Note that the parameter λ_1 necessarily lies on the dividing circle \mathcal{S}^1 . We shall show below that each λ_k lies on \mathcal{S}^k .

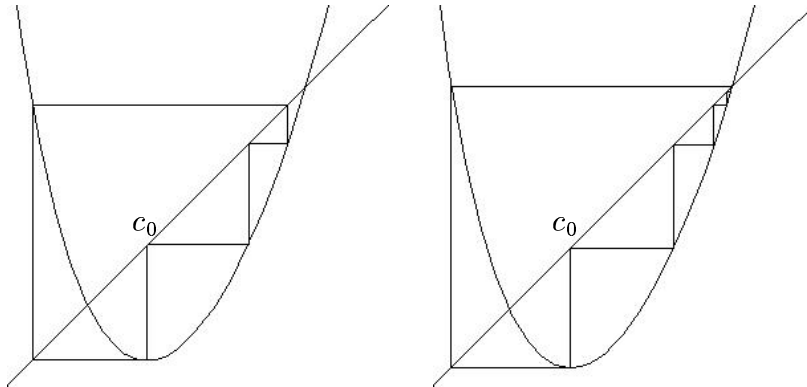


Figure 6: The graphs of F_λ for $\lambda = \lambda_4$ and $\lambda = \lambda_8$ when $n = 4$.

Because of the $(n - 1)$ -fold symmetry in the parameter plane, we have a similar sequence of superstable parameter values along the ray $\lambda = \omega \cdot \mathbb{R}^+$ in parameter plane. To be more precise, first suppose that n is even. Suppose

that $\lambda = a\omega$ with $a > 0$ and, as before, $\omega = \exp(2\pi i/(n-1))$. Then, using the results in Section 1, we have that, if $t > 0$,

$$F_\lambda(\omega^{\frac{n}{2}}t) = \omega^{\frac{n}{2}}F_a(t)$$

so that F_λ on the line $\omega^{n/2} \cdot \mathbb{R}^+$ is conjugate to F_a on \mathbb{R}^+ .

Now F_λ has critical points at

$$\begin{aligned} c_0 &= (a\omega)^{\frac{1}{2n}} \\ c_1 &= \nu(a\omega)^{\frac{1}{2n}} \\ c_{n+1} &= \nu^{n+1}(a\omega)^{\frac{1}{2n}} = -\nu(a\omega)^{\frac{1}{2n}} = -c_1. \end{aligned}$$

Note that the critical point c_{n+1} lies on the line $\omega^{n/2} \cdot \mathbb{R}^+$. This follows since

$$\begin{aligned} -\nu(a\omega)^{\frac{1}{2n}} &= -(a)^{\frac{1}{2n}} \left(\exp\left(\frac{\pi i}{n}\right) \exp\left(\frac{\pi i}{n(n-1)}\right) \right) \\ &= -(a)^{\frac{1}{2n}} \exp\left(\frac{\pi i}{n-1}\right) \\ &= -a^{\frac{1}{2n}} \omega^{\frac{1}{2}} = a^{\frac{1}{2n}} \omega^{\frac{n}{2}}. \end{aligned}$$

Therefore the above Proposition goes over to the case where $\lambda = a\omega$ with $a = \lambda_k \in \mathbb{R}^+$ provided we now use the critical point c_{n+1} lying on the line $\omega^{n/2} \cdot \mathbb{R}^+$. We note that the symmetric critical point c_1 lies on the line $\omega^{1/2} \cdot \mathbb{R}^+$ and maps onto the critical value on the line $\omega^{n/2} \cdot \mathbb{R}^+$ after one iteration.

The case where n is odd is similar modulo the $z \mapsto -z$ symmetry described earlier. The difference is that the superattracting cycles now have period $2k$ and alternate back and forth between $\omega \cdot \mathbb{R}^+$ and $-\omega \cdot \mathbb{R}^+$. We have:

Proposition (Superstable parameters for $\lambda \in \omega \cdot \mathbb{R}^+$). *Let $\lambda_1 > \lambda_2 \dots$ be the decreasing sequence in \mathbb{R}^+ in the previous Proposition. Suppose n is even. For $\lambda = \lambda_k \omega$, the critical point c_{n+1} is periodic with period k and the critical*

orbit along the line $\omega^{n/2} \cdot \mathbb{R}^+$ has the special form when $k \geq 2$

$$F_\lambda(c_{n+1}) < c_{n+1} = F_\lambda^k(c_{n+1}) < F_\lambda^{k-1}(c_{n+1}) < \dots < F_\lambda^3(c_{n+1}) < F_\lambda^2(c_{n+1}).$$

In particular, $\lambda = \lambda_k \omega$ is a superstable parameter value of period k and the orbit of $F_\lambda^2(c_{n+1})$ is monotonically decreasing for $k - 1$ iterations along $\omega^{n/2} \cdot \mathbb{R}^+$. When n is odd, replace F_λ with F_λ^2 . The cycle corresponding to $\lambda = \lambda_k \omega$ now has period $2k$.

3 Rings in Dynamical Plane

In this section we prove the existence of infinitely many rings γ_λ^k for $k = 0, 1, \dots$ in the dynamical plane. Each ring γ_λ^k is a smooth, simple closed curve that is mapped n^k -to-1 onto the critical circle by F_λ^k . We shall use these rings in the next section to construct the rings \mathcal{S}^k in the parameter plane.

We begin by defining γ_λ^0 to be the critical circle. Recall that, if $\lambda \in \mathcal{O}$, then F_λ maps γ_λ^0 strictly inside itself. Since all of the critical points of F_λ lie on γ_λ^0 , it follows that F_λ takes the exterior of γ_λ^0 as an n -to-1 covering onto the plane minus the critical segment and hence over the entire exterior of γ_λ^0 . Thus there is a preimage γ_λ^1 lying outside of γ_λ^0 and mapped n -to-1 onto γ_λ^0 by F_λ . Since F_λ is a covering map, it follows that γ_λ^1 must be a single simple closed curve. Then F_λ maps the exterior of γ_λ^1 as an n -to-1 covering onto the exterior of γ_λ^0 , so there is a preimage of γ_λ^1 lying in this region and mapped n -to-1 to γ_λ^1 . Call this simple closed curve γ_λ^2 . Continuing inductively, we find a collection of simple closed curves γ_λ^k for $k \geq 1$ having the properties that:

1. γ_λ^{k+1} lies in the exterior of γ_λ^k ;
2. F_λ takes γ_λ^{k+1} as an n -to-1 covering onto γ_λ^k ;

3. so F_λ takes γ_λ^{k+1} as an n^{k+1} -to-1 covering of the critical circle;
4. the γ_λ^{k+1} converge outward to the boundary of B_λ as $k \rightarrow \infty$.

We now construct a parameterization of each of the γ_λ^k . In order for this parametrization to be well-defined, we need to restrict attention to parameters in the region $\mathcal{O}' = \mathcal{O} - (-\lambda^*, 0]$, so that $-\pi < \text{Arg } \lambda < \pi$. We therefore assume that λ lies in \mathcal{O}' for the remainder of this paper.

For $\lambda \in \mathcal{O}'$, there is a unique critical point of F_λ lying in the region $|\text{Arg } z| < \pi/2n$. Call this critical point $c_0 = c_0(\lambda)$. Note that $c_0 \in \mathbb{R}^+$ if $\lambda \in \mathbb{R}^+$. We index the remaining critical points by c_j with the argument of c_j increasing as j increases.

To parametrize the critical circle γ_λ^0 , we set $\gamma_\lambda^0(0) = c_0(\lambda)$. By the Mapping Properties Proposition, for each $\theta \in \mathbb{R}$, we then let $\gamma_\lambda^0(\theta)$ be the natural continuation of this parametrization of the circle in the counterclockwise direction. So $\gamma_\lambda^0(\theta)$ is 2π -periodic in θ and depends analytically on λ for $\lambda \in \mathcal{O}'$.

To parametrize $\gamma_\lambda^1(\theta)$, consider the portion of the critical point sector containing $c_0(\lambda)$ that lies outside the critical circle. There is a unique point in this region mapped to c_0 by F_λ ; call this point $\gamma_\lambda^1(0)$. Then define $\gamma_\lambda^1(\theta)$ by requiring that

$$F_\lambda(\gamma_\lambda^1(\theta)) = \gamma_\lambda^0(\theta)$$

and that $\gamma_\lambda^1(\theta)$ varies continuously with θ . Note that $\gamma_\lambda^1(\theta)$ is $2n\pi$ periodic since F_λ is n -to-1 on γ_λ^1 . We then proceed inductively to define $\gamma_\lambda^k(\theta)$ by first specifying that, in the outside portion of the critical point sector containing c_0 , $\gamma_\lambda^k(0)$ is the unique point that is mapped by F_λ to $\gamma_\lambda^{k-1}(0)$ and then using F_λ to complete this parameterization. As above, for each k , $\gamma_\lambda^k(\theta)$ is $2n^k\pi$ periodic in θ and depends analytically on λ .

To prove the existence of the rings in the parameter plane, we need to be more specific about the location of the rings in the dynamical plane. Let V_+

be the portion of the prepole sector lying on and outside the critical circle and also between the two critical point rays through c_0 and c_1 . That is,

$$V_+ = \left\{ z \mid |z| \geq |\lambda|^{1/2n}, \frac{\text{Arg } \lambda}{2n} \leq \text{Arg } z \leq \frac{\text{Arg } \lambda}{2n} + \frac{\pi}{n} \right\}.$$

Let $V_- = \nu^{-1} \cdot V_+$. So V_- is the portion of the prepole sector bounded by the critical lines through c_0 and c_{-1} and lying on or outside the critical circle. Let $V_\lambda = V_+ \cup V_-$. See Figure 7.

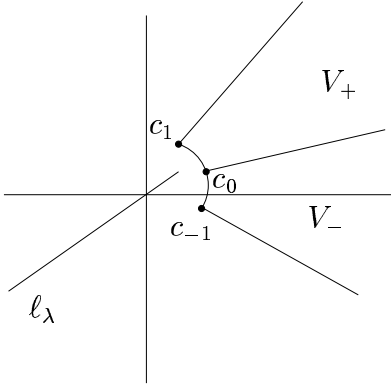


Figure 7: The region $V_\lambda = V_+ \cup V_-$.

Since $|\text{Arg } \lambda| < \pi$ and $n \geq 3$, we have for $z \in V_\lambda$

$$|\text{Arg } z| \leq \left| \frac{\text{Arg } \lambda}{2n} \right| + \frac{\pi}{n} < \frac{3\pi}{2n} \leq \frac{\pi}{2}.$$

So for each $\lambda \in \mathcal{O}'$, the region V_λ is contained in the half plane $\text{Re } z > 0$.

Now F_λ maps the portion of boundary of V_+ lying along the critical circle one-to-one to the critical segment since the endpoints of this arc are adjacent critical points along C_λ that are mapped to distinct critical values. Also, F_λ maps the portion of the critical point line containing c_0 lying on the boundary of V_+ one-to-one onto the ray $tv_\lambda = 2t\sqrt{\lambda}$ with $t \geq 1$ and $\text{Arg } \sqrt{\lambda} > 0$, while F_λ maps the other boundary ray containing c_1 to the negative of this ray.

Hence the boundary of V_+ is mapped onto the entire straight line passing through $\pm v_\lambda$ and the origin. Therefore F_λ maps V_+ univalently onto one of the half planes bounded by this line. Similarly, F_λ maps V_- univalently onto the opposite half plane.

Let ℓ_λ be the straight line given by $2t\sqrt{\lambda}$ where $t \in (-\infty, 1]$. So ℓ_λ is the straight line that starts at $2\sqrt{\lambda}$ at $t = 1$ and passes through the origin and $-2\sqrt{\lambda}$ enroute to ∞ as $t \rightarrow \infty$. Note that the boundary of V_λ is mapped two-to-one onto ℓ_λ by F_λ . Hence F_λ maps the interior of V_λ univalently onto $\mathbb{C} - \ell_\lambda$. Now, for each $\lambda \in \mathcal{O}'$, the critical segment lies outside V_λ since neither V_+ nor V_- meets the interior of the critical circle. Also, the portion of ℓ_λ extending from $-2\sqrt{\lambda}$ to ∞ lies in the left half plane, so the entire line ℓ_λ does not intersect V_λ . So we have:

Proposition. *For each $\lambda \in \mathcal{O}'$, F_λ maps the interior of V_λ univalently onto $\mathbb{C} - \ell_\lambda$ and so the image of V_λ contains V_λ .*

Recall that the k^{th} ring in the dynamical plane is parametrized by $\gamma_\lambda^k(\theta)$ and is periodic with period $2n^k\pi$.

Proposition. *For each $k \geq 1$, the portion of the ring $\gamma_\lambda^k(\theta)$ with $|\theta| \leq n^{k-1}\pi$ lies in the region*

$$-\frac{3\pi}{2n} < \text{Arg } z < \frac{3\pi}{2n}.$$

Proof: We deal first with the case $0 \leq \theta \leq n^{k-1}\pi$; the other case is handled by applying the $\rightleftharpoons \nu^{-1}z$ symmetry, as we describe below.

We claim that the portion of the ring $\gamma_\lambda^k(\theta)$ with $0 \leq \theta \leq n^{k-1}\pi$ actually lies in the smaller region

$$-\frac{\pi}{2n} < \text{Arg } z < \frac{3\pi}{2n}.$$

To see this, we first consider the simplest case where $\lambda \in \mathbb{R}^+$. In this case, V_+ is bounded by \mathbb{R}^+ and $\nu \cdot \mathbb{R}^+$ and F_λ maps V_+ univalently onto $\text{Im } z \geq 0$.

Recall that $\gamma_\lambda^0(\theta)$ lies in the region $\text{Im } z \geq 0$ if $\theta \in [0, \pi]$. Hence there is a continuous preimage of $\gamma_\lambda^0(\theta)$ lying in V_+ . This preimage is, by definition, $\gamma_\lambda^1(\theta)$ for $\theta \in [0, \pi]$. So $\gamma_\lambda^1(\theta)$ lies in the region $0 \leq \text{Arg } z \leq \pi/n$ and thus the result is true when $k = 1$.

Next note that $\gamma_\lambda^1(\pi)$ lies on the line $\nu \cdot \mathbb{R}^+$ and is given by $\nu\gamma_\lambda^1(0)$. So we can use the symmetry in the dynamical plane to extend the definition of $\gamma_\lambda^1(\theta)$ to a continuous curve defined for $\theta \in [0, n\pi]$ as follows: if $\theta \in [j\pi, (j+1)\pi]$, let $\gamma_\lambda^1(\theta) = \nu^j \gamma_\lambda^1(\theta - j\pi)$ for $j = 1, \dots, n-1$. So $\gamma_\lambda^1(\theta)$ lies in $\text{Im } z \geq 0$ for $\theta \in [0, n\pi]$. Then the sector V_+ is again mapped over $\gamma_\lambda^1(\theta)$ for these θ -values, so we have a continuous preimage $\gamma_\lambda^2(\theta)$ lying in V_+ , mapped onto $\gamma_\lambda^1(\theta)$, and defined for $\theta \in [0, n\pi]$.

Then we extend the definition of $\gamma_\lambda^2(\theta)$ to $[0, n^2\pi]$ as above using the symmetry in the dynamical plane. So we have that $\gamma_\lambda^3(\theta)$ lies in V_+ for all $\theta \in [0, n^2\pi]$. Continuing in this fashion proves the stronger result that $\gamma_\lambda^k(\theta)$ in fact lies in V_+ for $\theta \in [0, n^{k-1}\pi]$ for all k as long as $\lambda \in \mathbb{R}^+$.

Now suppose that $0 < \text{Arg } \lambda < \pi$. We no longer have the fact that V_+ is mapped over $\gamma_\lambda^0(\theta)$ for $0 \leq \theta \leq \pi$. Indeed, the point $\gamma_\lambda^1(0)$ now lies in V_- . This follows from the fact that the critical point ray through c_0 is mapped to a line whose argument is strictly larger than that of c_0 , so the preimage of c_0 must lie below this critical point line. By the previous Proposition, we have that F_λ maps the interior of the entire region V_λ univalently onto $\mathbb{C} - \ell_\lambda$. Let ℓ'_λ denote the portion of ℓ_λ lying in the lower half plane. Then

$$\pi < \frac{\text{Arg } \lambda}{2} + \pi = \text{Arg } \ell'_\lambda < \frac{3\pi}{2}.$$

Since, for $\theta \in [0, \pi]$, we have

$$0 < \text{Arg } c_0 \leq \text{Arg } \gamma_\lambda^0(\theta) \leq \text{Arg } c_0 + \pi < \frac{\text{Arg } \lambda}{2} + \pi = \text{Arg } \ell'_\lambda,$$

it follows that the entire line ℓ_λ never meets $\gamma_\lambda^0(\theta)$ for these θ -values. Hence there is a continuous preimage of $\gamma_\lambda^0(\theta)$ in $V_+ \cup V_-$ for each $\theta \in [0, \pi]$. This

defines $\gamma_\lambda^1(\theta)$ over this interval. Note that $\gamma_\lambda^1(\pi) = \nu\gamma_\lambda^1(0)$ must lie in V_+ . In fact, we can say more:

$$-\frac{\pi}{2n} < \frac{\text{Arg } \lambda}{2n} - \frac{\pi}{2n} \leq \text{Arg } \gamma_\lambda^1(\theta)$$

for $0 \leq \theta \leq \pi$. This follows since F_λ maps the prepole line in V_- to a line perpendicular to ℓ_λ in $-\pi/2 < \text{Arg } z < 0$. This line does not intersect the curve $\gamma_\lambda^0(\theta)$ for $\theta \in [0, \pi]$. So $\gamma_\lambda^1(\theta)$ does not meet the prepole line in V_- . We therefore have

$$-\frac{\pi}{2n} < \text{Arg } \gamma_\lambda^1(\theta) < \frac{3\pi}{2n}$$

for $\theta \in [0, \pi]$, so this proves the case $k = 1$ when $0 < \text{Arg } \lambda < \pi$.

Now we extend the definition of $\gamma_\lambda^1(\theta)$ to $\theta \in [0, n\pi]$ as in the previous case using symmetry. Then we have, for $0 \leq \theta \leq n\pi$,

$$-\frac{\pi}{2n} < \text{Arg } \gamma_\lambda^1(\theta) \leq \text{Arg } c_0 + \pi.$$

But $\text{Arg } c_0 + \pi < \text{Arg } \lambda/2 + \pi = \text{Arg } \ell'_\lambda$. So again ℓ_λ does not meet the extension of $\gamma_\lambda^1(\theta)$. So we have that $\gamma_\lambda^2(\theta)$ lies in the interior of $V_+ \cup V_-$ for $0 \leq \theta \leq n\pi$ and so $\text{Arg } \gamma_\lambda^2(\theta) < 3\pi/2n$. As above we in fact also have $-\pi/2n \leq \text{Arg } \gamma_\lambda^2(\theta)$, so this proves the case $k = 2$. Continuing inductively proves the result for all k -values when $0 < \text{Arg } \lambda < \pi$ and $0 \leq \theta \leq n^{k-1}\pi$.

The case of negative values of θ is handled by symmetry as follows. We again assume that $0 < \text{Arg } \lambda < \pi$. For each k we have, since $\gamma_\lambda^k(\theta)$ is $2n^k\pi$ periodic,

$$\begin{aligned} F_\lambda(\nu^{-1}\gamma_\lambda^k(\theta)) &= -F_\lambda(\gamma_\lambda^k(\theta)) \\ &= -\gamma_\lambda^{k-1}(\theta) \\ &= \gamma_\lambda^{k-1}(\theta - n^{k-1}\pi) \\ &= F_\lambda(\gamma_\lambda^k(\theta - n^{k-1}\pi)). \end{aligned}$$

Therefore

$$\nu^{-1}\gamma_\lambda^k(\theta) = \gamma_\lambda^k(\theta - n^{k-1}\pi)$$

follows since $\gamma_\lambda^k(\theta)$ is continuous in θ . Therefore we have that, when $\theta \in [-n^{k-1}\pi, 0]$, $\gamma_\lambda^k(\theta)$ lies in the region

$$-\frac{3\pi}{2n} < \text{Arg } z < \frac{\pi}{2n}.$$

So altogether the curve $\gamma_\lambda^k(\theta)$ lies in the region $|\text{Arg } z| < 3\pi/2n$ for all $|\theta| \leq n^{k-1}\pi$. This concludes the proof when $0 \leq \text{Arg } \lambda < \pi$.

If $-\pi < \text{Arg } \lambda < 0$, we invoke the $z \mapsto \bar{z}$ symmetry in the parameter plane. Since F_λ is conjugate to $F_{\bar{\lambda}}$ via $z \mapsto \bar{z}$, it follows that the curves $\gamma_\lambda^k(\theta)$ are mapped to $\gamma_{\bar{\lambda}}^k(-\theta)$ by the conjugacy. Hence these curves lie in the same region when $-\pi < \text{Arg } \lambda < 0$. This concludes the proof.

4 Rings in Parameter Plane

Before turning to the proof of the existence of rings in the parameter plane, we need to examine more carefully the parametrizations of the rings in the dynamical plane in two of the special cases discussed earlier, namely when $\lambda \in \mathbb{R}^+$ and $\lambda \in \omega \cdot \mathbb{R}^+$.

First suppose that $\lambda \in \mathbb{R}^+$. For the special parameters λ_k among the superstable parameters in \mathbb{R}^+ , we have seen that $F_{\lambda_k}(c_0)$ always lies in \mathbb{R}^+ and satisfies

$$0 < F_{\lambda_k}(c_0) < c_0 = F_{\lambda_k}^k(c_0) < F_{\lambda_k}^{k-1}(c_0) < \dots < F_{\lambda_k}^2(c_0).$$

Hence $F_{\lambda_k}^2(c_0)$ lies on $\gamma_{\lambda_k}^{k-2} \cap \mathbb{R}^+$ and $F_{\lambda_k}^j(c_0)$ lies on $\gamma_{\lambda_k}^{k-j} \cap \mathbb{R}^+$ for $j = 2, \dots, k$.

In particular, since the definition of the parametrization requires that $F_\lambda(\gamma_\lambda^j(0)) = \gamma_\lambda^{j-1}(0)$, it follows that, for the special parameter value λ_k , we

have

$$\begin{aligned}
\gamma_{\lambda_k}^0(0) &= c_0 \\
\gamma_{\lambda_k}^{k-2}(0) &= F_{\lambda_k}^2(c_0) \\
\gamma_{\lambda_k}^{k-3}(0) &= F_{\lambda_k}^3(c_0) \\
&\vdots \\
\gamma_{\lambda_k}^1(0) &= F_{\lambda_k}^{k-1}(c_0)
\end{aligned}$$

Next we turn attention to the special parameter values $\lambda_k \omega$ lying along the line $\omega \cdot \mathbb{R}^+$ in the parameter plane. Here the situation is somewhat more complicated. For simplicity of notation, we fix a value of k and set $\mu = \lambda_k \omega$.

As we showed earlier, the line $\omega^{n/2} \cdot \mathbb{R}^+$ contains the critical point c_{n+1} and is either invariant under F_μ (if n is even) or interchanged with the symmetric line $-\omega^{n/2} \cdot \mathbb{R}^+$ by F_μ (if n is odd). In either case the symmetric line $-\omega^{n/2} \cdot \mathbb{R}^+$ is mapped to this line by F_μ and contains the critical point $c_1 = -c_{n+1}$. Also, the critical point line through c_0 is mapped to $-\omega^{n/2} \cdot \mathbb{R}^+$ by F_μ and then to $\omega^{n/2} \cdot \mathbb{R}^+$ by F_μ^2 .

We have, by definition, $\gamma_\mu^0(0) = c_0$. Since $c_1 = \nu c_0$ where, as usual, $\nu = \exp(\pi i/n)$, we also have

$$\begin{aligned}
c_1 &= \gamma_\mu^0\left(\frac{\pi}{n}\right) \\
c_{n+1} &= \gamma_\mu^0\left(\frac{\pi}{n} + \pi\right).
\end{aligned}$$

Consider the portion of the critical point sector containing c_0 and lying on or outside C_λ . $\gamma_\mu^1(0)$ is the unique point in this region that is mapped to c_0 by F_μ . Since F_μ takes the critical point line through c_0 to the critical point line through c_1 , it follows that $\gamma_\mu^1(0)$ lies below this line and that $\gamma_\mu^1(\pi/n)$, the preimage of c_1 , lies on the critical point line through c_0 . By symmetry, $\gamma_\mu^1((\pi/n) + \pi)$ then lies on the critical point line through c_1 and, since γ_μ^1 is

$2n\pi$ -periodic, the point

$$\gamma_\mu^1\left(\frac{\pi}{n} + \pi + n\pi\right)$$

lies on the line $\omega^{n/2} \cdot \mathbb{R}^+$ containing c_{n+1} .

Continuing, we have that $\gamma_\mu^2((\pi/n) + \pi)$ lies on the critical point line through c_0 and is mapped by F_μ to $\gamma_\mu^1((\pi/n) + \pi)$. The point

$$\gamma_\mu^2\left(\frac{\pi}{n} + \pi + n\pi\right)$$

then lies on the critical point line through c_1 and is mapped to

$$\gamma_\mu^1\left(\frac{\pi}{n} + \pi + n\pi\right)$$

on $\omega^{n/2} \cdot \mathbb{R}^+$.

Continuing inductively, we see that the critical point line through c_0 contains the points

$$\begin{aligned} c_0 &= \gamma_\mu^0(0) \\ &\gamma_\mu^1\left(\frac{\pi}{n}\right) \\ &\gamma_\mu^2\left(\frac{\pi}{n} + \pi\right) \\ &\vdots \\ &\gamma_\mu^j\left(\frac{\pi}{n} + \pi + n\pi + \dots + n^{j-2}\pi\right) = \gamma_\mu^j\left(\frac{\pi}{n}(1 + n + \dots + n^{j-1})\right). \end{aligned}$$

and the critical point line through c_1 contains the points

$$\begin{aligned} c_1 &= \gamma_\mu^0\left(\frac{\pi}{n}\right) \\ &\gamma_\mu^1\left(\frac{\pi}{n} + \pi\right) \\ &\gamma_\mu^2\left(\frac{\pi}{n} + \pi + n\pi\right) \\ &\vdots \\ &\gamma_\mu^j\left(\frac{\pi}{n} + \pi + n\pi + \dots + n^{j-1}\pi\right) = \gamma_\mu^j\left(\frac{\pi}{n}(1 + n + \dots + n^j)\right). \end{aligned}$$

Equivalently, $\gamma_\mu^j(\theta)$ lies on the critical point line through c_1 for

$$\theta = \frac{\pi}{n} \left(\frac{n^{j+1} - 1}{n - 1} \right).$$

Now consider the corresponding points on the critical point line through c_{-1} . Since the parametrization corresponding to points on this line and γ_μ^j is obtained by subtracting $n^{j-1}\pi$ from the corresponding critical point line through c_0 , we find the following points on this critical point line:

$$\begin{aligned} c_{-1} &= \gamma_\mu^0 \left(-\frac{\pi}{n} \right) \\ &\gamma_\mu^1 \left(\frac{\pi}{n} - \pi \right) \\ &\gamma_\mu^2 \left(\frac{\pi}{n} + \pi - n\pi \right) \\ &\vdots \\ &\gamma_\mu^j \left(\frac{\pi}{n} + \pi + n\pi + \dots + n^{j-2}\pi - n^{j-1}\pi \right). \end{aligned}$$

Equivalently, $\gamma_\mu^j(\theta)$ lies on the critical point line through c_{-1} for

$$\theta = \frac{\pi}{n} (1 + n + n^2 + \dots + n^{j-1} - n^j) = \frac{\pi}{n} \left(\frac{n^j - 1}{n - 1} \right) - n^{j-1}\pi.$$

For later use, this value of θ is called $\theta_{n,j}$. See Figure 8.

We now turn to the proof the existence of the rings \mathcal{S}^k in parameter plane for $k > 1$. For technical reasons, we consider only the case when $n \geq 5$ in this section; the special cases $n = 3, 4$ are handled later.

Recall that, from the results of the previous section, we have that, when $k \geq 1$, the portion of the curve $\gamma_\lambda^k(\theta)$ for $|\theta| \leq n^{k-1}\pi$ lies in the region

$$-\frac{3\pi}{2n} < \text{Arg } z < \frac{3\pi}{2n}.$$

We call this region W_n and note that W_n lies in the right half plane. Let H_λ denote the involution that fixes c_0 , i.e.,

$$H_\lambda(z) = \frac{\lambda^{1/n}}{z}.$$

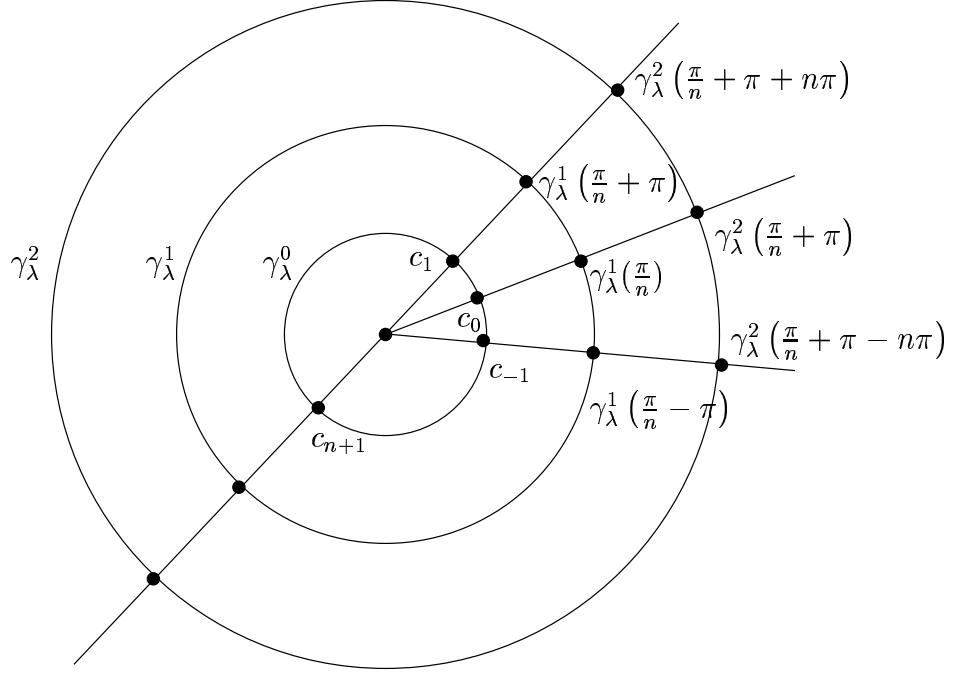


Figure 8: Parametrization of $\gamma_\lambda(\theta)$ when $\lambda = \lambda_k\omega$.

Lemma. *If $n \geq 5$ and $\lambda \in \mathcal{O}'$, then $H_\lambda(W_n)$ lies in the half plane $\operatorname{Re} z > 0$.*

Proof: Since

$$\operatorname{Arg} H_\lambda(z) = \frac{\operatorname{Arg} \lambda}{n} - \operatorname{Arg} z,$$

we have, if $z \in W_n$ and $n \geq 5$,

$$-\frac{\pi}{2} \leq -\frac{5\pi}{2n} \leq -\frac{3\pi}{2n} + \frac{\operatorname{Arg} \lambda}{n} < \operatorname{Arg} H_\lambda(z) < \frac{3\pi}{2n} + \frac{\operatorname{Arg} \lambda}{n} \leq \frac{5\pi}{2n} \leq \frac{\pi}{2}.$$

□

We remark that this result is false when $n = 3, 4$; that is the reason why these are special cases.

Now consider the curves

$$\xi_\lambda^k(\theta) = H_\lambda(\gamma_\lambda^k(\theta)).$$

Since the involution H_λ interchanges the inside and outside of C_λ , each of the curves ξ_λ^k is a simple closed curve lying inside the critical circle. We have

$$F_\lambda(\xi_\lambda^k(\theta)) = \gamma_\lambda^{k-1}(\theta)$$

since $F_\lambda(H_\lambda(z)) = F_\lambda(z)$. By the Lemma, we also have that $\xi_\lambda^k(\theta)$ lies in $\text{Re } z > 0$ for $|\theta| \leq n^{k-1}\pi$, at least if $n \geq 5$.

Theorem. *For each $k \geq 1$ and any θ satisfying $|\theta| \leq n^{k-1}\pi$, there exists a unique parameter $\lambda = \lambda_{\theta,k}$ such that*

$$v_\lambda = 2\sqrt{\lambda} = \xi_\lambda^k(\theta).$$

Proof: The function $G(\lambda) = v_\lambda = 2\sqrt{\lambda}$ takes the subset \mathcal{O}' of the parameter plane univalently onto an open subset of $\text{Re } z > 0$. For each $\lambda \in \mathcal{O}'$, $G(\lambda)$ lies inside C_λ , but for λ on the dividing circle (which is the circular boundary of \mathcal{O}'), $G(\lambda)$ lies on the critical circle. Hence G maps \mathcal{O}' univalently onto the interior of a half disk in the right half plane that contains the region inside C_λ in $\text{Re } z > 0$ for each $\lambda \in \mathcal{O}'$. Call this half disk D .

Also, for fixed θ , the function $\lambda \mapsto \xi_\lambda^k(\theta)$ is analytic on \mathcal{O}' and takes this set strictly inside the portion of the critical circle bounded by the rays $|\text{Arg } z| = 3\pi/2n$. Hence, for each θ , the set of points $\xi_\lambda^k(\theta)$ lies inside a compact sector in D . That is, this set of points can possibly accumulate on the boundary of D only at the origin. Hence we may consider the composition $Q(\lambda) = G^{-1}(\xi_\lambda^k(\theta))$. As a function of λ , Q is analytic and maps the simply connected region \mathcal{O}' inside itself. By the Schwarz Lemma, Q has a unique fixed point in this set or on its boundary. But the fixed point cannot lie at $\lambda = 0$ since 0 is surrounded by the McMullen domain so that the curves ξ_λ^k are bounded away from the origin. Hence there must be a unique fixed point in the interior of D . This fixed point is $\lambda_{\theta,k}$.

□

Note that the fixed points $\lambda_{\theta,k}$ vary continuously with θ , so $\theta \mapsto \lambda_{\theta,k}$ is a curve in the parameter plane.

The following Proposition identifies the specific values of $\lambda_{\theta,k}$ corresponding to the special cases considered earlier.

Proposition. *When $\theta = 0$ and $k \geq 1$, the parameter values $\lambda_{0,k}$ are given by the parameters $\lambda_{k+1} \in \mathbb{R}^+$. When $\theta = \theta_{n,k}$, $\lambda(\theta, k)$ is given by $\omega \lambda_{k+1}$ on the symmetry line $\omega \cdot \mathbb{R}^+$.*

Proof: When $\lambda \in \mathbb{R}^+$, the points $\gamma_\lambda^j(0)$ also lie in \mathbb{R}^+ for each j . Since, as shown earlier, the parameter λ_{k+1} has the property that $v_{\lambda_{k+1}} \in \xi_{\lambda_{k+1}}^k$, $F_{\lambda_{k+1}}^2(c_0) \in \gamma_{\lambda_{k+1}}^{k-1} \cap \mathbb{R}^+$ and the forward orbit of this point decreases along \mathbb{R}^+ until meeting c_0 , it follows from the uniqueness of the parameter $\lambda_{0,j}$ that we must have $\lambda_{0,k} = \lambda_{k+1}$ for each $k \geq 1$.

When $\lambda = \lambda_{k+1}\omega$ and $\theta = \theta_{n,k}$, we know that the point $\gamma_\lambda^k(\theta_{n,k})$ lies on the critical point line through c_{-1} . Hence $H_\lambda(\gamma_\lambda^k(\theta_{n,k}))$ lies on the critical point line through c_1 and is given by $\xi_\lambda^k(\theta_{n,k})$. This point is then mapped by F_λ to the point on $\omega^{n/2} \cdot \mathbb{R}^+$ whose orbit meets c_{n+1} after $k - 1$ iterations of F_λ or F_λ^2 , depending upon whether n is even or odd. Hence $\lambda_{\theta_{n,k},k} = \lambda_{k+1}\omega$ as claimed.

□

Now the parameters in the previous Proposition are the unique parameters on the corresponding lines in parameter space for which the orbit of the second iterate of the appropriate critical point monotonically decreases along the corresponding line(s) for $k - 1$ iterations before returning to itself and becoming periodic. So the curve $\theta \mapsto \lambda_{\theta,k}$ meets each of these two symmetry lines only once. Hence the portion of this curve defined for $0 \leq \theta \leq \theta_{n,k}$

either lies outside the sector

$$0 \leq \text{Arg } \lambda \leq \frac{2\pi}{n-1}$$

for all values of θ or else this entire curve lies inside the sector. But the former cannot occur since this would imply that some $\lambda_{\theta,k}$ would lie in \mathbb{R}^- , contradicting the fact that each $\lambda_{\theta,k}$ lies in \mathcal{O}' . Hence the portion of the curve $\lambda_{\theta,k}$ defined for $0 \leq \theta \leq \theta_{n,k}$ is a continuous arc connecting $\theta = 0$ and $\theta = 2\pi/(n-1)$. It then follows by the $(n-1)$ -fold symmetry that, for each $k \geq 1$, $\lambda_{\theta,k}$ is a simple closed curve in parameter space which is periodic of period

$$\begin{aligned} (n-1)\theta_{n,k} &= (n-1) \left(\frac{\pi}{n} \left(\frac{n^k - 1}{n-1} \right) - n^{k-1}\pi \right) \\ &= \frac{\pi}{n} (-n^{k+1} + 2n^k - 1). \end{aligned}$$

We therefore define the ring \mathcal{S}^{k+1} to be the simple closed curve $\theta \mapsto \lambda_{\theta,k}$. That is, \mathcal{S}^{k+1} consists of parameter values for which the critical orbit has the following behavior:

1. both critical values lie inside the critical circle;
2. $F_\lambda^2(c_\lambda)$ lies on γ_λ^{k-1} ;
3. subsequent iterates decrease through the γ_λ^j until, at the k^{th} iterate, the critical orbit lands back on the critical circle.

We have shown:

Theorem. *When $n \geq 5$, the ring \mathcal{S}^{k+1} in parameter space is a simple closed curve that is parameterized by $\theta \mapsto \lambda_{\theta,k}$ and is periodic of period*

$$\frac{\pi}{n} (n^{k+1} - 2n^k + 1) = \frac{\pi}{n} ((n-2)n^k + 1).$$

In particular, since the critical points (resp., prepoles) of F_λ are located on $\gamma_\lambda^0(\theta)$ at $\theta = \pi j/n$ (resp., $(2j+1)\pi/2n$) for $0 \leq j < 2n$, we have the following count of superstable parameters and centers of Sierpinski holes along \mathcal{S}^{k+1} :

Corollary. *There are precisely $(n-2)n^k + 1$ parameters along \mathcal{S}^{k+1} that are superstable parameters. There are the same number of parameters that are centers of Sierpinski holes. These parameters alternate between these two types as the parameter winds around \mathcal{S}^{k+1} .*

5 The Special Cases $n = 3$ and $n = 4$

In this section we turn attention to the special cases of the Rings Around the McMullen Domain Theorem that cannot be handled by our general methods; these are the cases when $n = 3$ and $n = 4$. The problems that arose in these cases occurred because the function $G(\lambda) = F_\lambda(c_\lambda) = 2\sqrt{\lambda}$ did not cover enough of the curves γ_λ^k to provide us with the means to find the fixed point λ_θ . We remedy this in these cases by using the function $G(\lambda) = F_\lambda^2(c_\lambda)$ instead.

We first deal with the case $n = 4$. That is, suppose

$$F_\lambda(z) = z^4 + \frac{\lambda}{z^4}$$

so F_λ has eight free critical points, two critical values, and eight prepoles.

Let $\lambda^* = \lambda_4^*$ be the radius of the dividing circle in the parameter plane. Since λ^* is the unique solution of the equation $2|\sqrt{\lambda}| = |\lambda|^{1/8}$, we compute that $\lambda^* = |2^{-8/3}|$. Let \mathcal{O} be the open disk in parameter plane inside the dividing circle, so that $|\lambda| < \lambda^*$ in \mathcal{O} . As before, if $\lambda \in \mathcal{O}$, then the critical circle is mapped strictly inside itself by F_λ .

The following result gives a bound on the location of the boundary of B_λ that will prove useful later.

Proposition (Location of ∂B_λ .) *Suppose $|\lambda| \leq \lambda^*$. Then the boundary of B_λ is contained in the open annulus bounded by the circles of radius 0.9 and 1.2 centered at the origin. The inner (resp., outer) boundary of this annulus is mapped strictly inside (resp., outside) itself by F_λ . Moreover, if $|z| \geq 1.2$, then $z \in B_\lambda$.*

Proof: A straightforward computation shows that if $|z| = 0.9$, then

$$\begin{aligned} |F_\lambda(z)| &\leq (0.9)^4 + \frac{|\lambda|}{(0.9)^4} \\ &\leq (0.9)^4 + \left(\frac{1}{2}\right)^{8/3} (0.9)^{-4} \\ &< 0.897, \end{aligned}$$

so F_λ maps the circle of radius 0.9 inside a smaller circle about the origin. Hence B_λ lies strictly outside the circle of radius 0.9.

On the other hand, if $|z| \geq 1.2$ and $|\lambda| \leq \lambda^* < 1/4$, we have

$$|F_\lambda(z)| \geq |z|^4 - \frac{|\lambda|}{|z|^4} \geq |z|^4 - \frac{1}{4(1.2)^4} > |z|^4 - 0.13 > |z|^2 > |z|.$$

It follows that F_λ maps each circle of the form $|z| = r \geq 1.2$ strictly outside itself and hence the entire region $|z| \geq 1.2$ is contained in B_λ . Therefore the boundary of B_λ must lie somewhere strictly inside the annulus bounded by the circles of radii 0.9 and 1.2.

□

It is known that the McMullen domain consists of a single, open, simply connected region surrounding the origin in the parameter plane [3]. We denote this domain by \mathcal{M} . Recall that F_λ^2 maps all eight critical points onto the same point. The location of this second image of the critical points as λ

varies can also be used to parametrize the rings \mathcal{S}^k around \mathcal{M} in parameter plane. Therefore we let G be the function that gives the location of this point, that is,

$$G(\lambda) = F_\lambda^2(\lambda^{1/8}) = 16\lambda^2 + \frac{1}{16\lambda}.$$

Note that G is defined in parameter plane, depends analytically on λ , and has a simple pole when $\lambda = 0$. So even though F_λ has no free critical points when $\lambda = 0$, the map G is still well-defined and analytic at this parameter value.

The following result gives an estimate of the size and location of \mathcal{M} in parameter space.

Proposition (Location of \mathcal{M} .) *The McMullen domain \mathcal{M} is contained in the open disk of radius $1/8$ centered at the origin in parameter space. The closed disk of radius $1/20$ centered at the origin lies inside \mathcal{M} .*

Proof: First suppose that $|\lambda| = 1/8$. Then we have

$$|G(\lambda)| \leq 16|\lambda|^2 + \frac{1}{16|\lambda|} = \frac{3}{4} < 0.9.$$

Since $1/8 < (1/2)^{8/3} = \lambda^*$, it follows from the previous Proposition that the second image of the critical points do not lie in B_λ when $|\lambda| = 1/8$. Therefore \mathcal{M} lies inside this circle.

On the other hand, if $|\lambda| \leq 1/20$, we have that

$$\begin{aligned} |G(\lambda)| &= \left| \frac{1}{16\lambda} + 16\lambda^2 \right| \\ &\geq \frac{1}{16|\lambda|} - 16|\lambda|^2 \\ &\geq \frac{20}{16} - \frac{16}{400} = 1.21. \end{aligned}$$

By the previous Proposition, $G(\lambda)$ lies in B_λ , but since $|\lambda| < \lambda^*$, the critical values themselves do not lie in B_λ . Therefore the critical values lie in T_λ and so $\lambda \in \mathcal{M}$.

□

Let \mathcal{W} be the interior of the disk of radius $1/8$ in parameter plane. Then we have $\mathcal{M} \subset \mathcal{W} \subset \mathcal{O}$.

Proposition. *The function G is one-to-one on \mathcal{W} .*

Proof: Suppose otherwise. Then there exist $\lambda, \mu \in \mathcal{W}$ with $\lambda \neq \mu$ but

$$16\lambda^2 + \frac{1}{16\lambda} = 16\mu^2 + \frac{1}{16\mu}$$

so that

$$16^2(\lambda^2 - \mu^2) = \frac{\lambda - \mu}{\lambda\mu}.$$

Since $\lambda \neq \mu$, we must then have

$$16^2|\lambda + \mu| = \frac{1}{|\lambda\mu|}.$$

But $16^2|\lambda + \mu| < 64$ whereas $1/|\lambda\mu| > 64$. This yields a contradiction.

□

Indeed, a computation shows that the function G has three critical points and each lies on the circle of radius $1/8$ centered at the origin. Therefore \mathcal{W} is the largest open circular disk surrounding the origin on which G is univalent.

Now let \mathcal{W}' denote the open simply connected region obtained by removing both the closed disk of radius $1/20$ about $\lambda = 0$ and the negative real axis from \mathcal{W} . We have:

Proposition. *The function G maps \mathcal{W}' univalently over a region that contains $\{z \mid 3/4 < |z| < 1.21\} - \mathbb{R}^-$.*

Proof: As in the proof of the Location of \mathcal{M} Proposition, if $|\lambda| = 1/8$, then $|G(\lambda)| \leq 3/4$. On the other hand, if $|\lambda| = 1/20$, then we have

$$|G(\lambda)| \geq \frac{1}{16|\lambda|} - 16|\lambda|^2 \geq 1.21.$$

Since G is one-to-one on all of \mathcal{W} and maps the interval $[-1/8, -1/20]$ into \mathbb{R}^- , the result follows. □

Proposition. *For each $\lambda \in \mathcal{W}'$ and each $k \geq 1$, the curves $\gamma_\lambda^k(\theta)$ all lie in the annulus $3/4 < |z| < 1.21$.*

Proof: As we showed earlier, the circle of radius 1.21 lies in B_λ for each $\lambda \in \mathcal{W}'$, so all of the γ_λ^k lie inside this circle. On the other hand, if $|z| = 3/4$, then

$$|F_\lambda(z)| \leq \left(\frac{3}{4}\right)^4 + |\lambda| \left(\frac{4}{3}\right)^4 < |\lambda|^{1/8}.$$

The second inequality follows by comparing the graphs of the linear expression on the left to that of $|\lambda|^{1/8}$ over the interval $1/20 \leq |\lambda| \leq 1/8$. Therefore, for each $\lambda \in \mathcal{W}'$, F_λ maps the circle of radius $3/4$ centered at the origin strictly inside C_λ . Hence all of the curves γ_λ^k with $k \geq 1$ must lie outside this circle. □

The proof of the existence of the rings \mathcal{S}^k now follows exactly as in the previous section, where we now use the function G defined on \mathcal{W}' . The image of G covers the region $|\text{Arg } z| \leq 3\pi/8$ in the annulus in the previous Proposition, and we showed in Section 3 that this region contains $\gamma_\lambda^k(\theta)$ for each $k \geq 1$. Thus we find the fixed point λ_θ that parameterizes \mathcal{S}^k for each $k \geq 3$. The case $k = 1$ is the dividing circle case described in Section 2. The case $k = 2$ in which the critical points land back on C_λ after two iterations (and hence do not meet any of the γ_λ^k with $k \geq 1$) is a similar straightforward computation.

When $n = 3$, the proof is essentially the same. The main difference in this case is that

$$G(\lambda) = 8\lambda^{3/2} + \frac{1}{8\sqrt{\lambda}}$$

no longer maps the disk about the origin univalently outside a circle in \mathbb{C} . Rather, in this case, G maps such a disk minus \mathbb{R}^- univalently onto the exterior of a circle in the right half plane. However, this is precisely where the curves γ_λ^k are constrained to lie when $n = 3$, so the proof goes through as above. We leave the details in this case to the reader.

References

- [1] Blanchard, P., Devaney, R. L., Look, D. M., Seal, P., and Shapiro, Y. Sierpinski Curve Julia Sets and Singular Perturbations of Complex Polynomials. To appear in *Ergodic Theory and Dynamical Systems*.
- [2] Devaney, R. L. Baby Mandelbrot Sets Adorned with Halos. To appear.
- [3] Devaney, R. L. Structure of the McMullen Domain in the Parameter Space of Rational Maps. To appear in *Fundamenta Mathematicae*.
- [4] Devaney, R. L. and Look, D. M. A Criterion for Sierpinski Curve Julia Sets. To appear in *Topology Proceedings*.
- [5] Devaney, R. L., Look, D. M., and Uminsky, D. The Escape Trichotomy for Singularly Perturbed Rational Maps. To appear in *Indiana U. Math. J.*
- [6] Douady, A. and Hubbard, J. Etude Dynamique des Polynômes Complexes. *Publ. Math. D'Orsay* (1984).
- [7] Douady, A. and Hubbard, J. On the Dynamics of Polynomial-like Mappings. *Ann. Sci. ENS Paris* **18** (1985), 287-343.

- [8] McMullen, C. Automorphisms of Rational Maps. *Holomorphic Functions and Moduli*. Vol. 1. Math. Sci. Res. Inst. Publ. **10**. Springer, New York, 1988.
- [9] McMullen, C. The Classification of Conformal Dynamical Systems. *Current Developments in Mathematics*. Internat. Press, Cambridge, MA, (1995) 323-360.
- [10] Milnor, J. *Dynamics in One Complex Variable*. Vieweg, 1999.
- [11] Milnor, J. and Tan Lei. A “Sierpinski Carpet” as Julia Set. Appendix F in Geometry and Dynamics of Quadratic Rational Maps. *Experiment. Math.* **2** (1993), 37-83.
- [12] Petersen, C. and Ryd, G. Convergence of Rational Rays in Parameter Spaces. In *The Mandelbrot set: Theme and Variations*, London Mathematical Society, Lecture Note Series 274, Cambridge University Press, 161-172, 2000.
- [13] Roesch, P. On Captures for the Family $f_\lambda(z) = z^2 + \lambda/z^2$. To appear.
- [14] Ushiki, S. Julia Sets with Polyhedral Symmetries. In *Fractals in the Fundamental and Applied Sciences*. Eds. Peitgen, H.-O., Henriques, J. M., and Penedo, L. F. Elsevier Publishers (1991), 425-442.
- [15] Whyburn, G. T. Topological Characterization of the Sierpinski Curve. *Fund. Math.* **45** (1958), 320-324.