# Percolation under noise: Detecting explosive percolation using the second-largest component

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We consider the problem of distinguishing between different rates of percolation under noise. A statistical model of percolation is constructed allowing for the birth and death of edges as well as the presence of noise in the observations. This graph-valued stochastic process is composed of a latent and an observed nonstationary process, where the observed graph process is corrupted by type-I and type-II errors. This produces a hidden Markov graph model. We show that for certain choices of parameters controlling the noise, the classical (Erdős-Rényi) percolation is visually indistinguishable from a more rapid form of percolation. In this setting, we compare two different criteria for discriminating between these two percolation models, based on the interquartile range (IQR) of the first component's size, and on the maximal size of the second-largest component. We show through data simulations that this second criterion outperforms the IQR of the first component's size, in terms of discriminatory power. The maximal size of the second component therefore provides a useful statistic for distinguishing between different rates of percolation, under physically motivated conditions for the birth and death of edges, and under noise. The potential application of the proposed criteria for the detection of clinically relevant percolation in the context of applied neuroscience is also discussed.

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### I. INTRODUCTION

Understanding the emergence of organized structure in dynamic networks remains an active research area [1,2]. In the study of random networks, percolation—the sudden emergence of a giant connected component (GCC)—is of critical importance from a theoretical, applied, and statistical perspective. Percolation in the Erdős-Rényi (ER) model constitutes one of the first examples of a fully characterized mathematical phase transition [3,4]. While the ER model of percolation is an example of a (second order) continuous phase transition, recent efforts have focused on identifying the conditions under which a random network process can yield a (first order) discontinuous percolation [5].

One of the most popular attempts to model discontinuous percolation has been the Achlioptas' process and its variants [6]. The Achlioptas' product rule (PR) slows down the growth of the GCC by favoring the creation of edges between small connected components. Although this particular percolation model has been shown to be, in fact, continuous and therefore of second order [5,7,8], it nonetheless provides an interesting alternative to the ER model [9]. Achlioptas' processes have indeed generated a substantial amount of theoretical work, whereby authors have explored related strategies for producing explosive percolation in random networks [10–12]. In addition, Riordan and Warnke [5] have shown that genuine first order phase transitions can be realized by systematically adding, at every step of the process, the edge that joins the two smallest components in the entire network. In this paper, we use the phrase "explosive percolation" to refer to the sudden emergence of a large connected component in the Achlioptas process. However, we note that explosive percolation in this context is rapid, but continuous.

Interest in network percolation has been fueled by its relevance to several application domains. In clinical neuroscience, for instance, epileptic seizures have been associated with the sudden emergence of coupled activity across the brain [13–18]. The resulting functional networks—in which edges indicate strong enough coupling between brain regions [19]—are consistent with the notion of percolation. A better understanding of the type of phase transitions undergone at different stages of the seizure may aid in the development of strategies for the treatment of epilepsy [20]. In this paper, we have therefore concentrated our attention on relatively small networks, with sizes ranging between 100 and 1000 vertices. Such a number of vertices is comparable to the size of the networks often studied in neuroscience [17,21–25]. However, note that percolation on such small networks can suffer from finite-size effects [26].

The rich theory on percolation, and its application to real world data, motivates the following question: How can we distinguish between different percolation regimes in practice? Previous theoretical work has concentrated on noise-free percolation, which constitutes an idealized perspective on percolation processes. In practice, however, the sampling of real-world networks is likely to be corrupted by measurement errors. Moreover, network growth has generally been conceived as a monotonic process, whereby only edge creations are allowed. However, this assumption may be too restrictive, since in real-world networks, the number of edges may increase and decrease over time, in a stochastic manner (see example in Fig. 1). Finally, to the best of our knowledge, there does not currently exist a statistical framework for distinguishing between different types of percolation regimes in the presence of edge birth and edge death, as well as noise.

In this paper, we propose a framework to distinguish between different percolation regimes in practice. To do so, we formulate the problem of recognizing a percolation regime from noisy observations as a question of statistical inference. Under this framework, we compare the discriminatory power of two potential percolation features deduced from the evolution of the first and second component of an observed dynamic network. We test this framework in simulation by constructing a hidden Markov graph model, which encompasses both a



FIG. 1. Proportion of nodes in the largest component as a function of time for a functional network deduced from the electrocorticogram of a single patient with epilepsy during a seizure (see [27] for details). The weighted functional networks have been binarized by conducting independent hypothesis tests on the maximum absolute values of the cross correlations over 0.5-s windows with 50% overlap and after correcting for multiple comparisons. The black trace is a smoothed version of this process. Seizure onset was clinically estimated to occur at the vertical dotted line.

nonstationary latent process characterized by birth and death of edges, and an observed graph process that introduces both type-I and type-II errors. We show that edge death and noise renders the statistic deduced from the first component ineffective in distinguishing between the standard ER second order percolation and Achlioptas' explosive percolation. However, a different detection criterion—based on the size of the second component—successfully discriminates between the two percolation regimes in the presence of edge death and noise. These results provide a framework for distinguishing percolation regimes in practice.

#### **II. PERCOLATION MODELS**

### A. Birth and death Erdős-Rényi (ER) process

We first construct a graph-valued stochastic process that exhibits the Markov property. This provides a realistic model for generating noisy percolation processes, while maintaining a sufficient level of computational tractability. We will denote a sequence of graph-valued random variables on n vertices by

$$\{G_t = (V, E_t) : t = 0, \dots, T\},$$
(1)

where V denotes the vertex set of  $G_t$ , whereas  $E_t$  represents the edge set of G at time t. Observe that the vertex set does not vary with time. At each time step, a single edge is either added or deleted. Such a sequence will be said to be *Markov* if its edge sets are controlled by a Markov chain. We impose this dependence through the use of a binary random variable, denoted  $\{Y_t : t = 0, ..., T\}$ , whose state space is  $\{0, 1\}$ . This Markov chain is characterized by the following transition probability matrix P, for some choices of the birth and death rates, denoted respectively by p and q, and taking values in [0,1]:

Following customary notation, the entries of *P* will be denoted by  $\mathbb{P}[Y_{t+1} = j | Y_t = i]$ , with rows summing to 1. The graphvalued Markov chain  $G_t$  is then obtained by associating  $Y_t$  with the addition or deletion of an edge in each edge set  $E_t$ . Thus, provided that  $p,q \neq 0$ , it follows that the state space of this graph-valued Markov chain is the space of all simple graphs on *n* vertices, since every graph is reachable with positive probability.

In general, p and q are not required to sum to 1. It will be of interest to let p > q in order to study the large-scale behavior of the  $G_t$ 's as the graph process accumulates edges. Moreover, observe that  $Y_t$  is a (time) homogeneous Markov chain, since  $\mathbb{P}[Y_{t+1} = \omega | Y_t = \omega'] = \mathbb{P}[Y_1 = \omega | Y_0 = \omega']$ , for any  $\omega, \omega' \in \{0, 1\}$  and every t.

Now, suppose that there exist  $m_t := |E_t|$  edges at time t in  $G_t$  and let  $X_t(e)$  denote the "status" of edge e at time t, such that  $X_t(e) = 1$ , if that edge is present, and  $X_t(e) = 0$ , otherwise. Note that we have here two different sources of dependence. On one hand, the edges are dependent on each other, since no more than one edge can be added or deleted at every time step. On the other hand, the edges are also dependent over time, since the status of an edge at time t + 1 depends on the status of that same edge at time t.

In the sequel, we will concentrate on a special case of this birth and death process, where we will set p = 1 - q. This leads to simplified marginal distributions for the edges. Additional details of this birth and death model are provided in Appendix.

#### B. Birth and death product rule (PR) process

We extend the standard Achlioptas' framework of PR percolation to a birth and death process, by devising death steps. This model is analogous to the aforementioned ER birth and death model, except for the choice of the probability distribution of the latent  $X_t(e)$ 's. As for the ER birth and death model, a binary random variable  $Y_t$  controls the addition or deletion of edges in each  $G_t$ . However, in the case of the PR model, the choice of the edge to be added or to be deleted is not uniform over the  $E_t$ 's. Here, this choice depends on the modular structure of the graph at time t. Therefore, as for the ER model, we obtain a nonstationary stochastic process.

Assuming that  $Y_t = 1$ , the addition of a new edge is conducted by uniformly choosing two candidate vertex pairs among all the edges in  $E_t^C$ , the complement of the edge set  $E_t$ . These two candidate edge pairs are denoted by  $e_1 := (v_{11}, v_{12})$ and  $e_2 := (v_{21}, v_{22})$ , and satisfy  $X_t(e_1) = 0$  and  $X_t(e_2) = 0$ , since  $e_1, e_2 \in E_t^C$ , as in Fig. 2. We then evaluate the size of the connected components to which  $v_{11}$ ,  $v_{12}$ ,  $v_{21}$ , and  $v_{22}$ belong. These four connected components are denoted by  $C_{11}$ ,  $C_{12}$ ,  $C_{21}$ , and  $C_{22}$ , respectively. Then, following Achlioptas *et al.* [6], we apply the following product rule: If  $|C_{11}||C_{12}| < |C_{21}||C_{22}|$ , then  $X_{t+1}(e_1) = 1$ ; otherwise,  $X_{t+1}(e_2) = 1$ .

Conversely, the death or deletion of an edge is handled in a symmetric manner. When  $Y_t = 0$ , we uniformly select



FIG. 2. Birth and death steps for the product rule (PR) in the Achlioptas' model of percolation. In this example, edge  $e_2$  was born *before* edge  $e_1$ , since  $|C_{11}||C_{12}| > |C_{21}||C_{22}|$ . Therefore, when  $e_1$  and  $e_2$  are selected during a death step,  $e_2$  is discarded *after*  $e_1$ . This death step specification ensures that births and deaths constitute genuine reverse PR operations.

two candidate edges from  $E_t$ . These vertex pairs are denoted  $e_1 := (v_{11}, v_{12})$  and  $e_2 := (v_{21}, v_{22})$  and satisfy  $X_t(e_1) = 1$  and  $X_t(e_2) = 1$ , since  $e_1, e_2 \in E_t$ . Next, we set  $X_t(e_1) = 0$  and  $X_t(e_2) = 0$ , in order to compute the size of the connected components to which  $v_{11}$ ,  $v_{12}$ ,  $v_{21}$ , and  $v_{22}$  would belong, if these edges were absent. This is done in order to ensure that the deletion of an edge exactly corresponds to the reverse operation of the addition of an edge under PR. Next, after having deleted these edges and having computed the sizes of  $C_{11}$ ,  $C_{12}$ ,  $C_{21}$ , and  $C_{22}$ , we decide which edge should re-enter  $G_t$ , in order to produce  $G_{t+1}$ . Such a decision is also based on the PR, such that if  $|C_{11}||C_{12}| < |C_{21}||C_{22}|$ , then  $X_{t+1}(e_2) = 0$ ; otherwise,  $X_{t+1}(e_1) = 0$ .

This choice of specification for the death step ensures that the *ordering* of the creation and deletion of edges is symmetrical. Given a sequence of two edges  $\{e_1, e_2\}$  successively born during two time steps of  $G_t$ , if we encounter a death step, where both  $e_1$  and  $e_2$  are selected, we would then delete these edges in the reverse order, by eliminating  $e_2$  before  $e_1$ . This order-preserving property is illustrated in Fig. 2. This constraint ensures that births and deaths are genuine reverse PR operations. In addition, observe that, as for the ER percolation process, this chain is *irreducible*, in the sense that there is positive probability of transitioning from any given edge configuration to any other in the space of the edge sets of G.

#### C. Hidden Markov graph model

Next, we assume that there exists a time-independent error process, which produces at each time point an *observed* edge status  $X_t^*(e)$ . This stochastic process is governed by two additional parameters  $\alpha$  and  $\beta$ , whose behavior can be described using a traditional "confusion matrix," such that for any  $\alpha, \beta \in [0, 1]$ , we have

The  $X_t(e)$ 's and  $X_t^*(e)$ 's are here treated as *latent* and *observed* stochastic processes, respectively, and  $\alpha$  and  $\beta$  can



FIG. 3. Directed acyclic graph (DAG) representation of the hidden Markov process combining a *latent* stochastic graph process in the first row denoted by  $G_t$ , with an *observed* stochastic graph process contaminated by noise in the second row, denoted by  $G_t^*$ . Directed arrows indicate probabilistic dependence, such that the distribution of the observed  $G_t^*$  depends on the value taken by the latent graph  $G_t$ .

therefore be interpreted as the type-I (false positive) and type-II (false negative) error probabilities. Combining the graph-valued Markov latent process with this time-independent error process, we obtain a graph-valued *hidden Markov process*, as described in Fig. 3. From this schematic representation, one can immediately see that the observed graphs denoted  $G_{t-1}^*$ ,  $G_t^*$  and  $G_{t+1}^*$  are conditionally independent, given the latent graph process  $G_t$ .

For the ER model, under the assumption that p = 1 - q, these two stochastic graph processes can be combined by taking into account the time dependence of the  $X_t^*(e)$ 's. In this case, the corresponding transition matrix linking the observed and latent processes is available in closed form. Details of these derivations are provided in Appendix.

## **III. DETECTING EXPLOSIVE PERCOLATION**

Explosive percolation is expected to produce a sharper phase transition than a typical ER percolation. When considering noisy observations, however, detecting such differences through visual inspection only, is hard. Figure 4 illustrates this problem, by comparing noisy and noise-free graph sequences for both explosive and ER models. Beyond visual inspection, the problem of discriminating between these two models of percolation can be formulated as a hypothesis-testing problem: The null hypothesis, denoted  $H_0$ , states that the observed process corresponds to an ER percolation, whereas the alternative hypothesis,  $H_1$ , is that the observed process does not correspond to this type of percolation model.

To proceed with this hypothesis-testing problem in practice, we specify a population parameter summarizing the percolation process, say  $\theta^{\text{ER}}$  and  $\theta$ , for the ER and target models, respectively. This leads to a hypothesis test of the form

$$H_0: \theta^{\text{ER}} = \theta$$
, and  $H_1: \theta^{\text{ER}} \neq \theta$ .

Several population parameters could be used for the purpose of discriminating between these two models of percolation. A natural candidate for such parameters would be a measure of the sharpness of the transition of the first component's size. The main panel of Fig. 4, however, suggests that this population parameter will not have sufficient discriminatory power, when confronted with a substantial amount of observational noise.



FIG. 4. Percentage of vertices in the giant connected component (GCC) in the main window, and the corresponding sizes of the secondlargest component (SLC) in inset, for a birth rate of p = 1 and death rate of q = 0, for the ER (black) and PR (gray) percolation models, for  $N_v = 100$ ; in which the x axis denotes the number of steps in the process divided by  $N_v$ . In bold, these results are reported for a noise-free model, whereas the thin lines represent noisy simulations with type-I and type-II error rates of  $\alpha = 0.0075$ ,  $\beta = 0$ . The details of the algebraic relationship between the type-I and type-II error rates and the graph process is described in Appendix, for the case of the ER model. Observe that the ER and PR models are nearly indistinguishable once a small amount of noise is added to these percolation processes.

Therefore, as a second candidate population parameter, we consider the following natural extension: The size of the second-largest component (SLC). This provides a more sensitive marker of the sharp phase transition exhibited by explosive percolation models [28]. Several authors have considered the size of the second largest component as a useful marker. Margolina et al. [28] have investigated the ratio of the sizes of the first and second component in cubic and triangular lattices. They have shown that the number of nodes in the second largest component reaches a maximum at the percolation threshold. More recently, and also in the context of cubic and hypercubic lattices, da Silva et al. [29] have studied the scale invariance of the ratio of the sizes of the first and second components. The size of the second largest component has also been used in neuroimaging, in order to identify the percolation threshold [30], as well as in an effort to detect and prevent epileptic seizures [31]. However, to the best of our knowledge, it has not been previously used in a statistical context for the purpose of discriminating between different types of percolation regimes. In what follows, we will show that the use of the size of the SLC as a statistical marker to distinguish the two percolation regimes exhibits greater discriminatory power than a statistic solely based on the first component.

The differences between the candidate percolation models are therefore quantified using two criteria: (i) the interquartile range (IQR) of the distribution of the size of the GCC, and (ii) the maximal size of the SLC over the entire time period. These two criteria are formally defined as follows. Given the graph process,  $G_t = (V, E_t)$ , and denoting the vertex subset of the largest component in  $G_t$  by  $S_{1,t}$ , we define the cumulative edge function as the *cardinality* of  $S_{1,t}$ , normalized by the maximal number of edges in the graph, such that

$$F(t) := \frac{|S_{1,t}|}{\binom{n}{2}}.$$

Although this function is not a cumulative distribution function (CDF), one can nonetheless uniquely define quantiles using the standard definition of quantiles for the CDFs of discrete random variables, such that for any  $x \in [0,1]$ , we have

$$Q(x) := \min_{t=1,\dots,T} \{t : F(t) \ge x\},\$$

where T is the maximal number of time steps in the graph process. In this paper, we are especially interested in the classical interquartile range,

$$IQR := Q(0.75) - Q(0.25).$$
(2)

This parameter quantifies the steepness of the phase transition: the larger the IQR criterion, the longer the transition to a fully connected graph.

As a second criterion to distinguish the two percolation regimes, we consider the maximal size attained by the SLC over the entire time period of the dynamic network observation. If one defines the vertex set of the SLC at time *t* by  $S_{2,t}$ , this second criterion can be expressed as

$$\theta_{\text{SLC}} = \max\{|S_{2,t}| : t = 1, \dots, T\}.$$

This quantity is known to constitute a good marker of the steepness of the phase transition, since it reflects the extent of separation of the graph process into large connected subgraphs [28,31]. Indeed, a direct consequence of the Achlioptas' construction rule is that by inhibiting the growth of a single large component, we necessarily increase the production of several subcomponents.

Statistical inference on these two criteria is then drawn using a Monte Carlo hypothesis test. Letting the parameter  $\theta := \theta_{IQR}$ , and selecting the candidate percolation to be drawn from a PR process, we consider the following null and alternative hypotheses,

$$H_0: \theta^{\text{ER}} = \theta^{\text{PR}}, \text{ and } H_1: \theta^{\text{ER}} > \theta^{\text{PR}},$$

respectively. The direction of this test is justified by the fact that we expect explosive percolation to occur rapidly, and thus to exhibit a smaller amount of variability in the size of its GCC, when transitioning to a fully connected graph. Our second criterion, by contrast, is tested in the opposite direction, since we naturally anticipate the PR process to be characterized by a *larger* maximal SLC. Thus, for  $\theta := \theta_{SLC}$ , the alternative hypothesis becomes  $H_1:\theta^{ER} < \theta^{PR}$ .

In the results reported in this paper, the distributions of the ER and PR graph processes are known. It therefore suffices to simulate from these densities in order to construct the distribution of the two test statistics at hand. This procedure is illustrated in Fig. 5. We are especially interested in the discriminatory powers of these statistics, and we will therefore compare their respective merits, using the true positive and false positive rates, within a receiver operating characteristic (ROC) framework. The ROC curve illustrates the performance of a binary classifier by plotting the false positive rate against



FIG. 5. Densities of the IQR of the GCC in the main windows, and densities of the maximal size of the SLC in the insets, are reported for networks of sizes 100 and 1000 vertices in panels (a) and (b), respectively, for the ER (black) and PR (gray) percolation models, and based on 1000 realizations from the distributions of these two (noise-free) models, specifying a birth rate of p = 1 and death rate of q = 0, in which the x axis denotes the number of steps in the process divided by  $N_v$ . The IQR criterion has been described in Eq. (2). Note that the difference in scales of the density values of the x axes, which represent the number of steps scaled by network sizes,  $t/N_v$ .

the true positive rate, as the discriminating threshold varies. For presentational convenience, the distributions of interest were smoothed using a normal density kernel, before computing the ROC curves and corresponding *areas under the curves* (AUCs). The computation of the AUCs allows us to summarize the differences between the models over the entire time period (see Figs. 6-10).

#### **IV. RESULTS**

### A. Birth and death processes

The ER and PR percolation models were simulated on graphs with n = 100 and n = 1000 vertices. We first explored the effect of varying the birth and death rates (i.e., p) on the



FIG. 6. Percentage of vertices in the giant connected component (GCC) in the main window, and the corresponding sizes of the secondlargest component (SLC) in inset, for a birth rate of p = 1 and death rate of q = 0, for the ER (black) and the PR (gray) percolation models, for  $N_v = 1000$ , in which the x axis denotes the number of steps in the process divided by  $N_v$ . In bold, these results are reported for a noise-free model, whereas the thin lines represent noisy simulations with type-I and type-II error rates of  $\alpha = 0.0075$ ,  $\beta = 0$ . Observe that the ER and PR models are nearly indistinguishable once noise is added to these percolation processes.

behavior of the two statistical criteria of interest, under the ER and PR models. The results of these Monte Carlo simulations are reported in Fig. 11, where each curve is the mean of 1000 independent synthetic data sets. In these simulations, the death rate is set to q = 1 - p, and therefore the value of p controls both the birth and death rates.

The main effect of a change in p is to delay percolation, and to diminish the steepness of the phase transition. Observe that as p decreases, percolation tends to occur at a later time step in both the ER and PR models (see Fig. 11). In particular, for the lowest birth rate that we investigated (p = 0.7), the ER model did not produce a fully connected graph within the number of iterations considered, as can be seen from Fig. 11(a). When pwas set to values equal to or less than 0.5, no phase transition could be observed, and these results are not reported.

The size of the second component was similarly affected by changes in *p*. Decreasing the birth rate delayed the time at which the size of the SLC attained its highest value. Moreover, lower values of *p* also yielded SLCs with smaller maximal sizes, under both the ER and PR models. Interestingly, we note that the time points at which the SLC reaches a maximal size tended to coincide in both models. Thus, it would be difficult to distinguish between these two percolation models on the sole basis of the *timing* of the occurrence of the maximal size of the SLCs. By contrast, the relative maximal *size* of the SLCs in the ER and PR models differ by approximately one order of magnitude, thereby providing a natural criterion for discriminating between these two types of percolation, as can be observed by comparing Figs. 11(a) and 11(b).

We formally quantified these differences in discriminatory powers by studying the ROC curves of these two criteria under different choices of p (see Fig. 8). The maximal size of the SLC substantially outperforms the relative size of the first component, for all values of p. The stark difference between



FIG. 7. Percentage of vertices in giant connected component (GCC), and in second largest component (SLC), in (a) and (b), respectively; with respect to scaled time  $(t/N_v)$ , for  $N_v = 1000$  in which the *x* axis denotes the number of steps in the process divided by  $N_v$ . These results are reported for both the ER and PR models, in black and gray, respectively; and for different choices of the birth rate *p*. Each curve represents the mean of 1000 independent simulations.

these discriminatory criteria can be understood by considering the amount of overlap of the distributions of these two criteria in Fig. 5. Whereas the distributions of the IQR of the size of the GCC under the two models exhibit a large amount of overlap, the distributions of the maximal sizes of the SLC, by contrast, share very little common support. These differences in support account for the substantive gains in discriminatory power by the maximal size of the second component, as reported in Fig. 10.

In addition, we note that the IQR of the size of the first component was more sensitive to choices of p than



FIG. 8. Receiver operating characteristic (ROC) curves for tests based on the IQR and SLC statistics for networks of size  $N_v = 100$  and  $N_v = 1000$ , in (a) and (b), respectively, for different choices of birth rate p.

the maximal size of the SLC. As p diminishes, it becomes increasingly more difficult to discriminate between the ER and PR percolation models, using the IQR of the size of the first component. This suggests that this criterion is more sensitive to a nonzero death rate than the maximal size of the SLC, which provides further support for the use of this latter criterion, in practice.

### **B.** Percolation under noise

Second, we considered the effect of introducing noise in these models. The results reported in Fig. 12 were produced using our proposed hidden Markov graph model, and were averaged over 1000 simulations. We were especially interested



FIG. 9. Percentage of vertices in giant connected component (GCC), and in second largest component (SLC), in (a) and (b), respectively, with respect to scaled time  $(t/N_v)$ , for  $N_v = 1000$ ; in which the *x* axis denotes the number of steps in the process divided by  $N_v$ . These results are reported for both the ER and PR models, in black and gray, respectively, and for different choices of the error rate  $\alpha$ . Each curve represents the mean of 1000 independent simulations.

in the effect of type-I and type-II errors on our ability to discriminate between classical and explosive percolation, using the two criteria under scrutiny. Both the type-I and type-II error rates were made to vary between 0 and 0.01.

From Fig. 12, one can observe that the two types of errors had markedly different effects on the AUCs of the two discriminatory criteria. Introducing type-I errors led to a substantial diminution of the AUCs for both the IQR of the size of the first component in Fig. 12(a), and the maximal size of the SLC in Fig. 12(b). In particular, note that the two



FIG. 10. Receiver operating characteristic (ROC) curves for tests based on the IQR and SLC statistics for networks of size  $N_v = 100$  and  $N_v = 1000$ , in (a) and (b), respectively, for different choices of error rate  $\alpha$ .

criteria reached equivalent levels of discriminatory power for  $\alpha = 0.01$ . Thus, although the maximal size of the SLC remains a more useful criterion for distinguishing between the ER and PR models than the IQR of the size of the first component, these two criteria exhibited comparable performance, under a moderate amount of type-I error.

The impact of increasing the type-II error rate on the behavior of these two criteria was negligible. Introducing false negatives in the ER and PR models slightly increased the AUCs of both the IQR of the first component's size, and the maximal size of the second component. Thus, large type-II error rates may be marginally advantageous for discriminating between these two models of percolation, under the scenarios studied.



FIG. 11. Percentage of vertices in giant connected component (GCC), and in second largest component (SLC), in panels (a) and (b), respectively; with respect to scaled time  $(t/N_v)$ , with  $N_v = 100$ ; in which the *x*-axis denotes the number of steps in the process divided by  $N_v$ . These results are reported both the ER and PR models, in black and gray, respectively; for different choices of birth rates, *p*. Each curve represents the mean of 1,000 independent simulations.

# V. DISCUSSION AND CONCLUSIONS

In this paper, we have extended existing models of percolation by allowing for edge deletion steps and noisy observations. These modeling extensions have been articulated within a hidden Markov graph process, which builds links with the existing literature on the statistical properties of this family of models [32–34]. Moreover, we have compared different summary statistics for distinguishing between the ER and PR percolation models. Overall, for different birth and death rates, and for a range of noise levels, the maximal size of the SLC was found to have greater discriminatory power than the IQR of the size of the GCC.

Several methodological challenges remain before such models can be directly used for percolation detection on real-world data. Throughout this paper, we have considered the IQR of the size of the first component, using a particular



FIG. 12. Area under the curve (AUC) surface plot for the receiver operating characteristic (ROC) curves of the IQR of the GCC, and of the maximal size of the SLC, for different choices of the error rates  $\alpha$  and  $\beta$ , and with p = 1 and q = 0, under the simulation settings of Fig. 13, and with  $N_v = 100$ .

choice of quantiles for this discriminatory criterion. In practice, an optimal choice of quantiles for quantifying the steepness of such phase transitions may be motivated by different factors, including (i) the range of the observations, and (ii) the need for early detection. We discuss these two practical aspects, in turn.

First, note that when considering real-world applications, we rarely observe fully connected networks. In the data reported in Fig. 1, for instance, the size of the GCC encompasses at most 90% of the edges in the saturated network. The choice of the quantile interval of interest for the first component will be therefore automatically constrained by the range of the observations in the data at hand. Therefore, as in sequential detection analysis, the statistical objective is to detect the outcome, on the basis of as little data as possible. Such constraints would naturally lead to a relatively narrow quantile range.

Second, in the context of clinical neuroscience and with particular emphasis on the prevention of a seizure, the



FIG. 13. Percentage of vertices in giant connected component (GCC), and in second largest component (SLC), in (a) and (b), respectively, with respect to scaled time  $(t/N_v)$ , for  $N_v = 100$ , in which the *x* axis denotes the number of steps in the process divided by  $N_v$ . These results are reported for both the ER and PR models, in black and gray, respectively, and for different choices of the error rate  $\alpha$ . Each curve represents the mean of 1000 independent simulations.

detection of a percolation regime may be linked with patients' health and survival. In such cases, early detection will usually be favored, as this is likely to be associated with desirable clinical outcomes. Explosive percolation, such as the Achlioptas' PR process studied in this paper, is consistent with the sudden manifestation of a seizure as a highly synchronized event. Classifying models of percolation may then be utilized to deepen our understanding of seizures, and to gain a greater understanding of the mechanisms underlying epilepsy.

Further work in this area could be focused on estimating a percolation model from a given sequence of observed networks. In this sense, this work also contributes to the growing literature on time-indexed graph processes [35]. In such cases, the birth and death rates will need to be estimated, as well as the type-I and type-II error probabilities. These different parameters may not be fully identifiable from the data, and further assumptions are likely to be necessary, in order to discriminate between the two percolation models considered in this paper. Such estimation, however, may be amenable to a Bayesian formulation, as commonly implemented for hidden Markov models [36]. Note also that this work may be extended to some of the recently proposed generalized versions of the Achlioptas model [37–39].

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## APPENDIX: DETAILS OF BIRTH AND DEATH ER PROCESS

Here, we describe the closed-form formulas of the probabilities of edge inclusion and edge deletion in the observed graph processes under the ER model. These analytic results are obtained by assuming that the birth and death probabilities are straightforwardly related, such that p = 1 - q. Such derivations may be useful for other authors, who may want to replicate these results, or extend the applications of the noisy model of percolation.

In this birth and death graph process, each edge is treated separately by integrating out the dependence of all other edges in the graph, and considering the marginal distribution of every  $X_t(e)$ . As before, we will here refer to  $X_t(e)$  as the *latent* edge status, and  $m_t := |E_t|$  will indicate the number of edges in the graph at time t. Given the Markov random variable  $Y_t$ , the conditional transition matrix for every  $X_{t+1}(e)$ , given some value of  $Y_t$ , takes the following form:

	$X_{t+1}(e) = 0$	$X_{t+1}(e) = 1$
$X_t(e)=0$	$\frac{\binom{n}{2}-m_t-I\{Y_t=1\}}{\binom{n}{2}-m_t}$	$\frac{I\{Y_t=1\}}{\binom{n}{2}-m_t}$
$X_t(e) = 1$	$\frac{I\{Y_t=0\}}{m_t}$	$\frac{m_t - I\{Y_t = 0\}}{m_t}$

where  $\binom{n}{2}$  denotes the number of edges in a saturated graph of size *n*, and where  $I\{f(x)\}$  is the indicator function, which takes a value of 1 if f(x) is true and 0, otherwise.

In this paper, we have concentrated on a special case of this birth and death process, where we have set p = 1 - q. This choice of p and q leads to the following characterization of

the  $Y_t$  process:

$$\mathbb{P}[Y_{t+1} = 1] = \mathbb{P}[Y_{t+1} = 1 | Y_t = 0]$$
  
=  $\mathbb{P}[Y_{t+1} = 1 | Y_t = 1] = p,$  (A1)

and similarly,  $\mathbb{P}[Y_{t+1} = 0] = 1 - p$ . Under this simplifying assumption, the preceding conditional transition matrix becomes

	$X_{t+1}(e) = 0$	$X_{t+1}(e) = 1$
$X_t(e) = 0$	$\frac{\binom{n}{2} - m_t - p}{\binom{n}{2} - m_t}$	$\frac{p}{\binom{n}{2}-m_t}$
$X_t(e) = 1$	$\frac{1-p}{m_t}$	$\frac{m_t - 1 + p}{m_t}$

Each entry is obtained by taking the expectation with respect to  $Y_t$ . That is,  $\mathbb{P}[X_{t+1}(e) = \omega | X_t(e) = \omega'] = \mathbb{E}[\mathbb{P}[X_{t+1}(e) = \omega | X_t(e) = \omega', Y_t]]$ , for every  $\omega, \omega' \in \{0, 1\}$ , and where the marginal distribution of  $Y_t$  is known from Eq. (A1).

One can then combine the noise process described in Sec. II C with the birth and death stochastic process in order to

- L. D. Valdez, P. A. Macri, and L. A. Braunstein, Temporal percolation of the susceptible network in an epidemic spreading, PloS One 7, e44188 (2012).
- [2] J. E. de Freitas, L. dos Santos Lucena, and S. Roux, Percolation as a dynamical phenomenon, Physica A: Stat. Mech. Appl. 266, 81 (1999).
- [3] P. Erdős and A. Rényi, The evolution of random graphs, Magyar Tud. Akad. Mat. Kutató Int. Közl. 5, 17 (1960).
- [4] N. Alon and J. H. Spencer, *The Probabilistic Method* (Springer, London, 2004).
- [5] O. Riordan and L. Warnke, Explosive percolation is continuous, Science 333, 322 (2011).
- [6] D. Achlioptas, R. M. D'Souza, and J. Spencer, Explosive percolation in random networks, Science **323**, 1453 (2009).
- [7] R. A. da Costa, S. N. Dorogovtsev, A. V. Goltsev, and J. F. F. Mendes, Explosive Percolation Transition is Actually Continuous, Phys. Rev. Lett. **105**, 255701 (2010).
- [8] O. Riordan and L. Warnke, Achlioptas process phase transitions are continuous, Ann. Appl. Probab. 22, 1450 (2012).
- [9] R. M. D'Souza and J. Nagler, Anomalous critical and supercritical phenomena in explosive percolation, Nat. Phys. 11, 531 (2015).
- [10] T. Bohman and A. Frieze, Avoiding a giant component, Random Struct. Algorithms 19, 75 (2001).
- [11] A. Beveridge, T. Bohman, A. Frieze, and O. Pikhurko, Product rule wins a competitive game, Proc. Am. Math. Soc. 135, 3061 (2007).
- [12] M. Krivelevich, E. Lubetzky, and B. Sudakov, Hamiltonicity thresholds in Achlioptas processes, Random Struct. Algorithms 37, 1 (2010).
- [13] M. Guye, J. Regis, M. Tamura, F. Wendling, A. McGonigal, P. Chauvel, and F. Bartolomei, The role of corticothalamic coupling in human temporal lobe epilepsy, Brain 129, 1917 (2006).

link the latent and observed parts of the Markov hidden model. This gives the following table:

$$\begin{aligned} \frac{X_{t+1}^*(e) = 0}{X_t(e) = 0} & X_{t+1}^*(e) = 1\\ \hline X_t(e) = 0 & (1-\alpha) \left(\frac{\binom{n}{2} - m_t - p}{\binom{n}{2} - m_t}\right) & \alpha \left(\frac{p}{\binom{n}{2} - m_t}\right)\\ \hline X_t(e) = 1 & \beta \left(\frac{1-p}{m_t}\right) & (1-\beta) \left(\frac{m_t - 1+p}{m_t}\right) \end{aligned}$$

Since this transition matrix links the latent and observed stochastic processes, one can immediately derive the marginal probabilities of the  $X_t^*(e)$ 's, such that

$$\mathbb{P}[X_{t+1}^*(e) = 0] = (1 - \alpha) \left( \frac{\binom{n}{2} - m_t - p}{\binom{n}{2} - m_t} \right) + \beta \left( \frac{1 - p}{m_t} \right),$$

and similarly for  $\mathbb{P}[X_t^*(e) = 1]$ . Observe that the resulting  $X_t$  process is *nonstationary*. Moreover, nonstationarity also holds when considering the case p = 1 - q. Indeed, since the probability of adding a new edge at time t + 1 is dependent on the number of existing edges,  $m_t := |E_t|$ , at time t; it follows that the resulting joint distribution of any subset of the  $X_t$ 's depends on the choice of t.

- [14] S. C. Ponten, F. Bartolomei, and C. J. Stam, Small-world networks and epilepsy: Graph theoretical analysis of intracerebrally recorded mesial temporal lobe seizures, Clin. Neurophysiol. 118, 918 (2007).
- [15] K. Schindler, H. Leung, C. E. Elger, and K. Lehnertz, Assessing seizure dynamics by analysing the correlation structure of multichannel intracranial EEG, Brain 130, e65 (2007).
- [16] K. Schindler, C. E. Elger, and K. Lehnertz, Increasing synchronization may promote seizure termination: Evidence from status epilepticus, Clin. Neurophysiol. 118, 1955, (2007).
- [17] M. A. Kramer, U. R. Eden, E. D. Kolaczyk, R. Zepeda, E. N. Eskandar, and S. S. Cash, Coalescence and fragmentation of cortical networks during focal seizures, J. Neurosci. 30, 10076 (2010).
- [18] K. Schindler, F. Amor, H. Gast, M. Muller, A. Stibal, L. Mariani, and C. Rummel, Peri-ictal correlation dynamics of high-frequency (80–200Hz) intracranial EEG, Epilepsy Res. 89, 72 (2010).
- [19] M. Rubinov and O. Sporns, Complex network measures of brain connectivity: Uses and interpretations, Neuroimage 52, 1059 (2010).
- [20] M. A. Kramer and S. S. Cash, Epilepsy as a disorder of cortical network organization, Neuroscientist 18, 360 (2012).
- [21] E. Bullmore and O. Sporns, Complex brain networks: Graph theoretical analysis of structural and functional systems, Nat. Rev. Neurosci. 10, 1 (2009).
- [22] D. S. Bassett, E. T. Bullmore, A. Meyer-Lindenberg, J. A. Apud, D. R. Weinberger, and R. Coppola, Cognitive fitness of costefficient brain functional networks, Proc. Natl. Acad. Sci. USA 106, 11747 (2009).
- [23] C. E. Ginestet and A. Simmons, Statistical parametric network analysis of functional connectivity dynamics during a working memory task, NeuroImage 55, 688 (2011).

- [24] R. Salvador, J. Suckling, M. R. Coleman, J. D. Pickard, D. Menon, and E. Bullmore, Neurophysiological architecture of functional magnetic resonance images of human brain, Cereb. Cortex 15, 1332 (2005).
- [25] C. J. Honey, O. Sporns, L. Cammoun, X. Gigandet, J. P. Thiran, R. Meuli, and P. Hagmann, Predicting human resting-state functional connectivity from structural connectivity, Proc. Natl. Acad. Sci. USA 106, 2035 (2009).
- [26] Z. Wu, C. Lagorio, L. A. Braunstein, R. Cohen, S. Havlin, and H. E. Stanley, Numerical evaluation of the upper critical dimension of percolation in scale-free networks, Phys. Rev. E 75, 066110 (2007).
- [27] M. A. Kramer, U. T. Eden, S. S. Cash, and E. D. Kolaczyk, Network inference with confidence from multivariate time series, Phys. Rev. E 79, 061916 (2009).
- [28] A. Margolina, H. J. Herrmann, and D. Stauffer, Size of largest and second largest cluster in random percolation, Phys. Lett. A 93, 73 (1982).
- [29] C. R. da Silva, M. L. Lyra, and G. M. Viswanathan, Largest and second largest cluster statistics at the percolation threshold of hypercubic lattices, Phys. Rev. E 66, 056107 (2002).
- [30] A. Haimovici, E. Tagliazucchi, P. Balenzuela, and D. R. Chialvo, Brain Organization into Resting State Networks Emerges at Criticality on a Model of the Human Connectome, Phys. Rev. Lett. 110, 178101 (2013).
- [31] A. Ciurea, I. Mindruta, M. D. Maliia, A. Ciurea, J. Ciurea, A. Barborica, C. Donos, M. F. Casanova, and I. Opris, Modular

signatures and neural avalanches in epileptic brain networks. In *Recent Advances on the Modular Organization of the Cortex* (Springer, New York, 2015), pp. 271–295.

- [32] P. Flocchini, B. Mans, and N. Santoro, On the exploration of time-varying networks, Theor. Comput. Sci. 469, 53 (2012).
- [33] A. Casteigts, P. Flocchini, W. Quattrociocchi, and N. Santoro, Time-varying graphs and dynamic networks, Int. J. Parallel, Emerg. Distrib. Syst. 27, 387 (2012).
- [34] H. H. K. Lentz, T. Selhorst, and I. M. Sokolov, Unfolding accessibility provides a macroscopic approach to temporal networks, Phys. Rev. Lett. 110, 118701 (2013).
- [35] P. Grindrod and D. J. Higham, Evolving graphs: Dynamical models, inverse problems and propagation, Proc. R. Soc., Ser. A 0456, 1 (2009).
- [36] Y. Ephraim and N. Merhav, Hidden markov processes, IEEE Trans. Inf. Theory 48, 1518 (2002).
- [37] P. Giazitzidis and P. Argyrakis, Generalized achlioptas process for the delay of criticality in the percolation process, Phys. Rev. E 88, 024801 (2013).
- [38] R. A. da Costa, S. N. Dorogovtsev, A. V. Goltsev, and J. F. F. Mendes, Critical exponents of the explosive percolation transition, Phys. Rev. E 89, 042148 (2014).
- [39] R. A. da Costa, S. N. Dorogovtsev, A. V. Goltsev, and J. F. F. Mendes, Inverting the Achlioptas rule for explosive percolation, Phys. Rev. E 91, 042130 (2015).