

# Computation of the Wodzicki – Chern – Simons form in local coordinates. Computations for $S^1$ actions on $S^2 \times S^3$ (Last update : July 1, 2010.)

We define a local system of coordinates for  $M = S^2 \times S^3$  and compute the Christoffel symbols and curvature. We use a system of coordinates to write a family of metrics over  $M$  taken from

J.P. Gauntlett, D. Martelli, J. Sparks and D. Waldram

Sasaki-Einstein metrics on  $S^2 \times S^3$ .

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[hep-th/0403002].

In this system of coordinates,  $M$  can be considered as a non-trivial  $S^1$  fibration over  $S^2 \times S^2$ . We compute the zeroth and -1 order symbols for the connection form and the zeroth and -1 order symbol for the curvature form on the loop space  $LM$ . These computations allow us to calculate the relative Wodzicki-Chern-Simons form on  $LM$ . The natural  $S^1$  action on the fibers of  $M$  determines a homology class in  $LM$ . We compute the value of the WCS form on this class and related classes for the  $n$ -fold iteration of the  $S^1$  action. Each of these integrals are different (i.e. determine nonhomologous classes in  $H_5(LM)$ ) for  $n = 0, 1, 2, \dots$ . As explained in the paper, this proves that the action is not smoothly homotopic to the trivial action, and that  $\pi_1(\text{Diff}(S^2 \times S^3))$  is infinite.

We proceed as follows: In Section 1, we compute the Christoffel symbols and curvature for these metrics, and use them to compute explicitly all needed symbols for the corresponding metrics on  $LM$ . In Section 2, we use these computations to calculate the Wodzicki-Chern-Simons form on  $LM$ . We integrate the WCS form over a 5-cycle in  $LM$  associated to a circle action on  $M$ . For certain metrics, the integral is nonzero.

**Table of contents**

1. Preliminaries -- the main procedures to compute Christoffel symbols, curvature, and symbols for the connection and curvature forms on loop space.
  - 1.1 Computation of the metric.
  - 1.2 Computations of Christoffel symbols and curvature.
  - 1.3 The symbols
2. Evaluation of the relative Wodzicki – Chern – Simons form on a cycle in  $L(S^2 \times S^3)$  associated to the fiber action
  - 2.1 The contribution from  $\text{Tr } \sigma_{-1}(\omega) \wedge \sigma_0(\Omega) \wedge \sigma_0(\Omega)$

# 1. Preliminaries

```
SetDirectory["Mathematica/PaperSteveYoshi"];
```

```
SetDirectory["research/fabian/07/notebooks"];
```

Define the metric in spherical coordinates on  $S^2 \times S^3$ .  $\phi$  and  $\theta$  are the ordinary spherical coordinates for  $S^2$ , and  $\psi$ ,  $y$  and  $\tau$  are coordinates for  $S^3$ . We have  $\phi$  in  $[0, 2\pi]$ ,  $\theta$  in  $[0, \pi]$ ; with  $\theta$  the angle with respect to the z axis.  $\psi$  in  $[0, 2\pi]$ ,  $\tau$  in  $[0, 2\pi]$ , with  $l$  and the range of  $y$  defined below. In Gauntlett et al., the range for  $\tau$  and  $y$  are chosen so the metric extends to the closure  $S^2 \times S^3$  of the coordinate ball.

```
Clear[coord, metric, inversemetric,
      Christoffel, Theta, Phi, Xi, x, X1, X2, X3, X4, A, Con, NumIt];
dimension = 5;
coord = {Phi, Theta, Psi, y, Tau};
PsiForm = {P1si, P2si, P3si, P4si};

XField = {X1, X2, X3, X4, X5};
YField = {Y1, Y2, Y3, Y4, Y5};
WField = {W1, W2, W3, W4, W5};
ZField = {Z1, Z2, Z3, Z4, Z5};
TField = {T1, T2, T3, T4, T5};
DotGamma = {Dot1, Dot2, Dot3, Dot4, Dot5};
```

## 1.1 Computation of the metric

Computation of the general metric: First, we introduce the line element from (2.1) in Gauntlett et al.

```
Clear[a, c]; coord1 = {Phi1, Theta1, Psi1, y1, Tau1};
Funcw[y_, a_, c_] := 2 * (a - y * y) / (1 - c * y);
Funcq[y_, a_, c_] := (a - 3 * y * y + 2 * y * y * y * c) / (a - y * y);
```

Gauntlett's et al. line element  $ds^2$ . Here  $a$  and  $c$  are constants.

```
DeltaSquare[Phi1_, Theta1_, Psi1_, y1_, Tau1_, a_, c_] :=
(1 - c * y) / 6 (Theta1 * Theta1 + Sin[Theta1] * Sin[Theta1] * Phi1 * Phi1) + (1 / Funcw[y, a, c]) *
(1 / Funcq[y, a, c]) * y1 * y1 + 1 / 9 Funcq[y, a, c] * (Psi1 - Cos[Theta1] * Phi1) ^ 2 +
Funcw[y, a, c] * (Tau1 + (a * c - 2 * y + y * y * c) / (6 (a - y * y)) (Psi1 - Cos[Theta1] * Phi1)) ^ 2;
```

```
dimension = 5; Clear[a, c];
GenMetric[Phi_, Theta_, Psi_, y_, Tau_, a_, c_] :=
Table[D[DeltaSquare[Phi1, Theta1, Psi1, y1, Tau1, a, c], coord1[[ii]], coord1[[jj]]] /
If[ii == jj, 2, 2], {ii, 1, dimension}, {jj, 1, dimension}]
```

```
Simplify[GenMetric[Phi, Theta, Psi, y, Tau, a, c]]
```

$$\left\{ \left\{ -\frac{5 + a c^2 - 12 c y + 6 c^2 y^2 + (-1 + a c^2) \cos[2 \theta]}{36 (-1 + c y)}, 0, \frac{(2 + a c^2 - 6 c y + 3 c^2 y^2) \cos[\theta]}{18 (-1 + c y)}, \right. \right.$$

$$\left. 0, \frac{(a c + y (-2 + c y)) \cos[\theta]}{-3 + 3 c y} \right\}, \left\{ 0, \frac{1}{6} (1 - c y), 0, 0, 0 \right\},$$

$$\left\{ \frac{(2 + a c^2 - 6 c y + 3 c^2 y^2) \cos[\theta]}{18 (-1 + c y)}, 0, \frac{2 + a c^2 - 6 c y + 3 c^2 y^2}{18 - 18 c y}, 0, \frac{a c + y (-2 + c y)}{3 - 3 c y} \right\},$$

$$\left\{ 0, 0, 0, \frac{1 - c y}{2 a - 6 y^2 + 4 c y^3}, 0 \right\},$$

$$\left\{ \frac{(a c + y (-2 + c y)) \cos[\theta]}{-3 + 3 c y}, 0, \frac{a c + y (-2 + c y)}{3 - 3 c y}, 0, \frac{2 (-a + y^2)}{-1 + c y} \right\}$$

Explicit expression for the metric  $ds^2$ . We set  $c=1$  as in Gauntlett et al.

$c = 1$ ; GenMetric[Phi, Theta, Psi, y, Tau, a, c]

$$\left\{ \left\{ \frac{1}{2} \left( \frac{(a-2y+y^2)^2 \cos[\text{Theta}]^2}{9(1-y)(a-y^2)} + \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]^2}{9(a-y^2)} + \frac{1}{3}(1-y) \sin[\text{Theta}]^2 \right), 0, \right. \right.$$

$$\left. \frac{1}{2} \left( -\frac{(a-2y+y^2)^2 \cos[\text{Theta}]}{9(1-y)(a-y^2)} - \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]}{9(a-y^2)} \right), 0, -\frac{(a-2y+y^2) \cos[\text{Theta}]}{3(1-y)} \right\},$$

$$\left\{ 0, \frac{1-y}{6}, 0, 0, 0 \right\}, \left\{ \frac{1}{2} \left( -\frac{(a-2y+y^2)^2 \cos[\text{Theta}]}{9(1-y)(a-y^2)} - \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]}{9(a-y^2)} \right), \right.$$

$$\left. 0, \frac{1}{2} \left( \frac{(a-2y+y^2)^2}{9(1-y)(a-y^2)} + \frac{2(a-3y^2+2y^3)}{9(a-y^2)} \right), 0, \frac{a-2y+y^2}{3(1-y)} \right\},$$

$$\left\{ 0, 0, 0, \frac{1-y}{2(a-3y^2+2y^3)}, 0 \right\}, \left\{ -\frac{(a-2y+y^2) \cos[\text{Theta}]}{3(1-y)}, 0, \frac{a-2y+y^2}{3(1-y)}, 0, \frac{2(a-y^2)}{1-y} \right\}$$

End of the computation of the metric.

## 1.2 Computations of Christoffel symbols and curvature

$$\text{metric} = \left\{ \left\{ \frac{1}{2} \left( \frac{(a-2y+y^2)^2 \cos[\text{Theta}]^2}{9(1-y)(a-y^2)} + \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]^2}{9(a-y^2)} + \frac{1}{3}(1-y) \sin[\text{Theta}]^2 \right), 0, \right. \right.$$

$$\left. \frac{1}{2} \left( -\frac{(a-2y+y^2)^2 \cos[\text{Theta}]}{9(1-y)(a-y^2)} - \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]}{9(a-y^2)} \right), 0, -\frac{(a-2y+y^2) \cos[\text{Theta}]}{3(1-y)} \right\},$$

$$\left\{ 0, \frac{1-y}{6}, 0, 0, 0 \right\}, \left\{ \frac{1}{2} \left( -\frac{(a-2y+y^2)^2 \cos[\text{Theta}]}{9(1-y)(a-y^2)} - \frac{2(a-3y^2+2y^3) \cos[\text{Theta}]}{9(a-y^2)} \right), \right.$$

$$\left. 0, \frac{1}{2} \left( \frac{(a-2y+y^2)^2}{9(1-y)(a-y^2)} + \frac{2(a-3y^2+2y^3)}{9(a-y^2)} \right), 0, \frac{a-2y+y^2}{3(1-y)} \right\},$$

$$\left\{ 0, 0, 0, \frac{1-y}{2(a-3y^2+2y^3)}, 0 \right\}, \left\{ -\frac{(a-2y+y^2) \cos[\text{Theta}]}{3(1-y)}, 0, \frac{a-2y+y^2}{3(1-y)}, 0, \frac{2(a-y^2)}{1-y} \right\};$$

Computation of the metric inverse.

```
inversemetric = Simplify[Inverse[metric]];
```

```
metric = Simplify[metric];
inversemetric = Simplify[inversemetric];
```

Definition of the Christoffel Symbols. Notation:  $\text{Christoffel}[i,j,k]=\Gamma^i_{j,k}$ .

```
Christoffel = Christoffel = Table[
  Simplify[(1/2) * Sum[(inversemetric[[ii, ss]]) * (D[metric[[ss, jj]], coord[[kk]] +
    D[metric[[ss, kk]], coord[[jj]]) - D[metric[[jj, kk]], coord[[ss]])],
    {ss, 1, dimension}], {ii, 1, dimension}, {jj, 1, dimension}, {kk, 1, dimension}];
```

```
listaffine := Table[{ToString[ $\Gamma$ [i, j, k]], Christoffel[[i, j, k]]},
  {i, 1, dimension}, {j, 1, dimension}, {k, 1, dimension}]
```

Definition of the Curvature. Notation:  $\text{Curvature}[i,j,k,l]=R^i_{jkl}$ .

```
Curvature = Curvature =
  Table[Simplify[D[Christoffel[[ii, jj, ll]], coord[[kk]] - D[Christoffel[[ii, jj, kk]],
    coord[[ll]] + Sum[Christoffel[[ss, jj, ll]] Christoffel[[ii, kk, ss]] -
    Christoffel[[ss, jj, kk]] Christoffel[[ii, ll, ss]], {ss, 1, dimension}],
    {ii, 1, dimension}, {jj, 1, dimension}, {kk, 1, dimension}, {ll, 1, dimension}];
```

```
ListCurvature := Table[If[UnsameQ[Curvature[[i, j, k, l]], 0],
  {ToString[R[i, j, k, l]], Curvature[[i, j, k, l]]},
  {i, 1, dimension}, {j, 1, dimension}, {k, 1, dimension}, {l, 1, dimension}]
TableForm[Partition[DeleteCases[Flatten[ListCurvature], Null], 2], TableSpacing -> {2, 2}]
```

Computation of the covariant derivative for the curvature. Notation:  $\text{CovCurvature}[i,j,k,m,p]=R^m_{ijk;p}$

```

CovCurvature = CovCurvature = Table[D[Curvature[[m, i, j, k]], coord[[p]]] +
  Sum[Curvature[[t, i, j, k]] * Christoffel[[m, p, t]], {t, 1, dimension}] -
  Sum[Curvature[[m, t, j, k]] * Christoffel[[t, p, i]], {t, 1, dimension}] -
  Sum[Curvature[[m, i, t, k]] * Christoffel[[t, p, j]], {t, 1, dimension}] -
  Sum[Curvature[[m, i, j, t]] * Christoffel[[t, p, k]], {t, 1, dimension}],
  {i, 1, dimension}, {j, 1, dimension}, {k, 1, dimension},
  {m, 1, dimension}, {p, 1, dimension}];

```

Computation of the Ricci curvature  $\text{Ricci}_{mk} = R^l_{mk}$ . This step is not needed for the computations related to the WCS form; we just did this to confirm the metric is Einstein with Einstein constant 4.

```

Ricci = Table[Sum[Curvature[[kk, ii, kk, jj]], {kk, 1, dimension}],
  {ii, 1, dimension}, {jj, 1, dimension}];

```

```

Ricci2 = Table[
  Sum[D[Christoffel[[kk, j, i]], coord[[kk]]] - D[Christoffel[[kk, kk, i]], coord[[j]]] +
  Sum[Christoffel[[kk, kk, s]] Christoffel[[s, j, i]] -
  Christoffel[[s, i, kk]] Christoffel[[kk, j, s]], {s, 1, dimension}],
  {kk, 1, dimension}, {i, 1, dimension}, {j, 1, dimension}];

```

```

Simplify[Ricci2 - 4 * metric] // TableForm

```

```

0 0 0 0 0
0 0 0 0 0
0 0 0 0 0
0 0 0 0 0
0 0 0 0 0

```

### 1.3 The symbols

Computation of the symbols for the connection.

```

SigmaMinusConn[XF_] :=
  -1/2 * Table[Sum[(-2 Curvature[[jj, kk, mm, ll]] - Curvature[[jj, mm, kk, ll]] +
  Curvature[[jj, ll, mm, kk]]) * XF[[mm]] * DotGamma[[ll]], {ll, 1, dimension},
  {mm, 1, dimension}], {jj, 1, dimension}, {kk, 1, dimension}];

```

```

SigmaZeroConn[XF_] :=
  Table[Sum[Christoffel[[jj, mm, kk]] * XF[[mm]], {mm, 1, dimension}],
    {jj, 1, dimension}, {kk, 1, dimension}];
SigmaMinusConn[XF_] := Simplify[-1/2 * Table[
  Sum[(Curvature[[jj, ll, mm, kk]]) * XF[[mm]] * DotGamma[[ll]], {ll, 1, dimension},
    {mm, 1, dimension}], {jj, 1, dimension}, {kk, 1, dimension}]];

```

Computation of the symbols for the curvature.

```

SigmaMinusCurOne[XF_, YF_] := 1/2 * Table[
  Sum[(2 * CovCurvature[[kk, ll, nn, jj, mm]] + CovCurvature[[ll, kk, nn, jj, mm]] +
    CovCurvature[[nn, ll, kk, jj, mm]]) * XF[[mm]] * YF[[ll]] * DotGamma[[nn]],
    {nn, 1, dimension}, {mm, 1, dimension}, {ll, 1, dimension}] +
  Sum[YF[[ll]] * (2 * Curvature[[jj, kk, ll, mm]] +
    Curvature[[jj, ll, kk, mm]] + Curvature[[jj, mm, ll, kk]]) *
    (Sum[D[XF[[mm]], coord[[nn]]] * DotGamma[[nn]], {nn, 1, dimension}] +
    Sum[Christoffel[[mm, ee, ff]] * DotGamma[[ee]] * XF[[ff]],
      {ee, 1, dimension}, {ff, 1, dimension}]), {mm, 1, dimension},
    {ll, 1, dimension}], {jj, 1, dimension}, {kk, 1, dimension}];

```

```

SigmaZeroCur[XF_, YF_] :=
  -Table[Sum[Curvature[[jj, kk, mm, ll]] * XF[[mm]] * YF[[ll]], {mm, 1, dimension},
    {ll, 1, dimension}], {jj, 1, dimension}, {kk, 1, dimension}];
SigmaMinusCurOne[XF_, YF_] := 1/2 * Table[
  Sum[(2 * CovCurvature[[kk, ll, nn, jj, mm]] + CovCurvature[[ll, kk, nn, jj, mm]] +
    CovCurvature[[nn, ll, kk, jj, mm]]) * XF[[mm]] * YF[[ll]] * DotGamma[[nn]],
    {nn, 1, dimension}, {mm, 1, dimension}, {ll, 1, dimension}] +
  Sum[YF[[ll]] * (2 * Curvature[[jj, kk, ll, mm]] +
    Curvature[[jj, ll, kk, mm]] + Curvature[[jj, mm, ll, kk]]) *
    (Sum[D[XF[[mm]], coord[[nn]]] * DotGamma[[nn]], {nn, 1, dimension}] +
    Sum[Christoffel[[mm, ee, ff]] * DotGamma[[ee]] * XF[[ff]],
    {ee, 1, dimension}, {ff, 1, dimension}]), {mm, 1, dimension},
    {ll, 1, dimension}], {jj, 1, dimension}, {kk, 1, dimension}];
SigmaMinusCur[XF_, YF_] := SigmaMinusCurOne[XF, YF] - SigmaMinusCurOne[YF, XF];

```

## 2. Evaluation of the relative Wodzicki-Chern-Simons form on a cycle in LM associated to the fiber action.

As explained in the paper, the  $S^1$  action on the fiber determines a cycle in LM. We evaluate the WCS form on this cycle by pullback to an integral over  $M$  and show that this integral is non-zero. This implies that the  $S^1$  action is nontrivial in the sense defined in the paper.

We compute the WCS form at a typical loop  $\gamma(\tau)$  of the  $S^1$  action. In the coordinates from Gauntlett et al. the loop obtained is:

$$S^1 \rightarrow L(S^2 \times S^3)$$

$$t \rightarrow (\phi, \theta, \psi, y, \tau + \text{Palpha} * t),$$

where Palpha measures the number of iterates of the action.

We compute the derivative of  $\gamma(t)$  with respect to  $t$ , and express the result in terms of the of the standard basis for  $T_{\gamma(t)} S^2 \times S^3$ . The outcome is stored in the variable (Dot1, Dot2, Dot3, Dot4, Dot5)=(0,0,0,0,Palpha).

This is an action via isometries, since the metric is independent of  $\tau$ . In the WCS integrand, the pullback is an integral over  $S^1$  of " $\omega \wedge \Omega \wedge \Omega$ " acting on  $a_*(\text{partial}_\phi)(t)$ ,  $a_*(\text{partial}_\psi)(t)$ , ... at  $a(t)$ , which we call  $\gamma(t)$ . We store in the vector fields X, Y, Z, T, and W the pushforwards  $a_*(\text{partial}_\phi)(t)$ ,  $a_*(\text{partial}_\psi)(t)$ , .... Since our action is an isometry, this integrand is independent of  $t$ , so instead of integrating over  $S^1$ , we can just multiply the integrand by  $2\pi$ . Below we first compute the integrand and then the pullback integral.



```

InitializeAction = Module[{Dotgam},
  Dot1 = 0;
  Dot2 = 0;
  Dot3 = 0;
  Dot4 = 0;
  Dot5 = PAlpha;
  (* The vector fields *)
  X1 = 1; X2 = 0; X3 = 0; X4 = 0; X5 = 0;
  Y1 = 0; Y2 = 1; Y3 = 0; Y4 = 0; Y5 = 0;
  Z1 = 0; Z2 = 0; Z3 = 1; Z4 = 0; Z5 = 0;
  T1 = 0; T2 = 0; T3 = 0; T4 = 1; T5 = 0;
  W1 = 0; W2 = 0; W3 = 0; W4 = 0; W5 = 1;
]

```

The next two matrices contain all the permutations of the fields X,Y,Z,W and T and their corresponding sign. They will be used to compute the evaluation of each permutation of the basis vector fields that contribute to the CSW form.

```

Permu = ListPermutations = Permutations[{XField, YField, ZField, WField, TField}];
SignPermu = {-1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1,
  -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1,
  1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1,
  1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1,
  -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1, -1, 1, 1, -1};

```

## 2.1 The contribution from $\text{Tr } \sigma_{-1}(\omega) \wedge \sigma_0(\Omega) \wedge \sigma_0(\Omega)$

```
Clear[ZeroMinusOneTerm1, ZeroMinusOneTerm2, ZeroMinusOneTerm3]
```

Definition of the 5-form  $\text{Tr } \sigma_{-1}(\omega) \wedge \sigma_0(\Omega) \wedge \sigma_0(\Omega)$  using the expressions above for the symbols. The 1/120 factor is added after the evaluation of the pullback.

```

ZeroMinusOneTerm1[TF_, WF_, XF_, YF_, ZF_] := Sum[SigmaMinusConn[TF][[mm, jj]] *
  SigmaZeroCur[WF, XF][[jj, ii]] * SigmaZeroCur[YF, ZF][[ii, mm]],
  {ii, 1, dimension}, {jj, 1, dimension}, {mm, 1, dimension}]

```

Compute the evaluation of every term of  $\text{Tr } \sigma_{-1}(\omega) \wedge \sigma_0(\Omega) \wedge \sigma_0(\Omega)$  on all possible permutations of T, X, Y, W, Z. The result is stored in the matrix Final.

```
Final = Apply[ZeroMinusOneTerm1, Permu, {1}];
```

```
DotGam = << GammaDot5.dat;
XFF = << S1xS2XField.dat;
YFF = << S1xS2YField.dat;
WFF = << S1xS2ZField.dat;
```

The variable Final is a vector with 120 entries. We add all the entries with the corresponding sign for the permutation.

```
Fin = Sum[Final[[j]] * SignPermu[[j]], {j, 1, Length[SignPermu]}];
```

Simplification of the expression for general a. Notice that the contribution vanishes when a = 1.

```
TotalComputation = Simplify[Fin]
```

$$\frac{32 (-1 + a)^2 P\text{Alpha } \gamma \text{ Sin}[\text{Theta}]}{3 (-1 + \gamma)^5}$$

Multiplication by the 1/120 factor.

```
TotalComputation = TotalComputation * 1 / 120
```

$$\frac{4 (-1 + a)^2 P\text{Alpha } \gamma \text{ Sin}[\text{Theta}]}{45 (-1 + \gamma)^5}$$

We have to integrate the variable  $TotalComputation$  over  $S^2 \times S^3$ . We begin with the integration with respect to the variable  $\theta$ .

```
TotalComputationOne = Integrate[TotalComputation, {Theta, 0, Pi}]
```

$$\frac{8 (-1 + a)^2 P\text{Alpha } y}{45 (-1 + y)^5}$$

Now, we integrate with respect to  $y$ . We do this in two steps. First we compute the indefinite integral.

```
FunctionSigma = TotalComputationOne ;
FunctionSigma = Simplify[Integrate[FunctionSigma, y]]
```

$$\frac{2 (-1 + a)^2 P\text{Alpha } (1 - 4 y)}{135 (-1 + y)^4}$$

For the second step, we evaluate the integral at the bounds  $y_1 =$

$\frac{1}{2}(1 - \lambda - \text{Sqrt}[1 - \lambda^2/3])$  and  $y_2 = y_1 + \lambda$ . We also express the value of  $a = 3y_1^2 - 2y_1^3$  in terms of  $\lambda$  as in (3.3), (3.4) and (3.5) in Gauntlett et al.

$$TotalComputationFunction[y_] := \frac{2 (-1 + a)^2 P\text{Alpha } (1 - 4 y)}{135 (-1 + y)^4}$$

```
Rooty1 = 1 / 2 (1 - 1 - Sqrt[1 - 1^2 / 3]);
Rooty2 = Rooty1 + 1;
```

```
FunctionPlot =
Simplify[TotalComputationFunction[Rooty2] - TotalComputationFunction[Rooty1] /.
a -> 3 * Rooty1 ^ 2 - 2 * Rooty2 ^ 3] /. 1 -> λ
```

$$\frac{1}{45 \left(3 - 2\lambda^2 + \sqrt{9 - 3\lambda^2}\right)^4} 6 \text{PAlpha} \lambda^3 \left(144 \lambda^5 \sqrt{9 - 3\lambda^2} + 252 \lambda^4 \left(-5 + \sqrt{9 - 3\lambda^2}\right) - 108 \left(3 + \sqrt{9 - 3\lambda^2}\right) - 108 \lambda^3 \left(3 + \sqrt{9 - 3\lambda^2}\right) + 16 \lambda^6 \left(3 + \sqrt{9 - 3\lambda^2}\right) + 9 \lambda^2 \left(57 + 17 \sqrt{9 - 3\lambda^2}\right)\right)$$

As in (3.1), (3.8) in Gauntlett et al., we set  $\lambda = 3q/(2p)$  for fixed relatively prime integers  $p, q$  satisfying  $4p^2 - 3q^2 = n^2$  for a positive integer  $n$ .

```
FunctionPlot = Simplify[FunctionPlot /. {sqrt[9 - 3 λ^2] -> (3 n) / (2 p), λ -> (3 q) / (2 p)}]
```

$$\frac{1}{5 p^2 (n p + 2 p^2 - 3 q^2)^4} 6 \text{PAlpha} q^3 \left(-2 p \left(16 p^6 - 57 p^4 q^2 + 54 p^3 q^3 + 315 p^2 q^4 - 27 q^6\right) + n \left(-16 p^6 + 51 p^4 q^2 - 54 p^3 q^3 + 189 p^2 q^4 + 162 p q^5 + 27 q^6\right)\right)$$

To compute the total contribution we trivially integrate with respect to  $0 < \alpha < 2\pi$ , where  $l = q/(3q^2 - 2p^2 + p(4p^2 - 3q^2)^{1/2})$  as in (3.6) in Gauntlett et al.; with respect to  $\phi$ , which gives a factor of  $2\pi$ ; with respect to  $\psi$ , which gives a factor of  $\pi$ , and with respect to the variables of the unit cotangent bundle of the loop, which gives a factor of  $(2)(2\pi) = 4\pi$ . This shows that the evaluation is a rational number times some powers of  $\pi$ . The following function allows to compute the total contribution plugging only values for  $p$  and  $q$ .

```
In[70]:= TotalContribution[p_, q_] := Module[{n}, n = Sqrt[4 p^2 - 3 q^2];
```

$$\frac{1}{5 p^2 (n p + 2 p^2 - 3 q^2)^4} 6 \text{PAlpha} q^3 \left(-2 p \left(16 p^6 - 57 p^4 q^2 + 54 p^3 q^3 + 315 p^2 q^4 - 27 q^6\right) + n \left(-16 p^6 + 51 p^4 q^2 - 54 p^3 q^3 + 189 p^2 q^4 + 162 p q^5 + 27 q^6\right)\right) * 2^4 * \text{Pi}^4 * (q / (3 q^2 - 2 p^2 + p * n))$$

We compute for  $(p, q) = (7, 3)$  and for  $(p, q) = (19, 5)$ , corresponding to the values  $(A, B) = (3, 1)$  and  $(A, B) = (5, 1)$ , respectively, as explained in (3.12) in Gauntlett et al. The results are nonzero.

In[74]:= **TotalContribution** [7, 3]

$$\text{Out[74]= } - \frac{1849 \text{ PAlpha } \pi^4}{22050}$$

In[73]:= **TotalContribution** [19, 5]

$$\text{Out[73]= } - \frac{312481 \text{ PAlpha } \pi^4}{36843660}$$

Conclusion: We compute that the WCS integral over  $S^2 \times S^3$  for the metric described in Gauntlett et al. is nonzero.